

CP asymmetry in flavour-specific B decays

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$B - \bar{B}$ mixing and a_{fs}

Schrödinger equation:

$$i \frac{d}{dt} \begin{pmatrix} |B(t)\rangle \\ |\bar{B}(t)\rangle \end{pmatrix} = \left(M - i \frac{\Gamma}{2} \right) \begin{pmatrix} |B(t)\rangle \\ |\bar{B}(t)\rangle \end{pmatrix}$$

3 physical quantities in $B - \bar{B}$ mixing:

$$|M_{12}|, \quad |\Gamma_{12}|, \quad \phi = \arg \left(-\frac{M_{12}}{\Gamma_{12}} \right)$$

Two mass eigenstates:

$$\text{Lighter eigenstate: } |B_L\rangle = p|B^0\rangle + q|\bar{B}^0\rangle.$$

$$\text{Heavier eigenstate: } |B_H\rangle = p|B^0\rangle - q|\bar{B}^0\rangle \quad \text{with } |p|^2 + |q|^2 = 1.$$

with masses $M_{L,H}$ and widths $\Gamma_{L,H}$.

Here B represents either B_d or B_s .

To determine $|M_{12}|$, $|\Gamma_{12}|$ and ϕ measure

$$\Delta m = M_H - M_L \simeq 2|M_{12}|, \quad \Delta\Gamma = \Gamma_L - \Gamma_H \simeq 2|\Gamma_{12}| \cos \phi$$

and

$$a_{\text{fs}} = \text{Im} \frac{\Gamma_{12}}{M_{12}} = \left| \frac{\Gamma_{12}}{M_{12}} \right| \sin \phi.$$

a_{fs} is the CP asymmetry in decays $B \rightarrow f$ which are flavour-specific, i.e.

$$\bar{B} \not\rightarrow f \text{ and } B \not\rightarrow \bar{f}.$$

Examples: $B_d \rightarrow X \ell^+ \nu_\ell$ or $B_s \rightarrow D_s^- \pi^+$.

$$a_{\text{fs}} = \frac{\Gamma(\bar{B}(t) \rightarrow f) - \Gamma(B(t) \rightarrow \bar{f})}{\Gamma(\bar{B}(t) \rightarrow f) + \Gamma(B(t) \rightarrow \bar{f})}$$

The time dependence of the decay rates $\Gamma(\bar{B}(t) \rightarrow f)$ and $\Gamma(B(t) \rightarrow \bar{f})$ drops out.

a_{fs} measures CP violation in mixing.

Standard Model expectation:

$$\frac{\Delta\Gamma}{\Delta m} \simeq \left| \frac{\Gamma_{12}}{M_{12}} \right| = \mathcal{O} \left(\frac{m_b^2}{M_W^2} \right) = 4 \cdot 10^{-3}$$
$$a_{\text{fs}} = \left| \frac{\Gamma_{12}}{M_{12}} \right| \sin \phi = \begin{cases} -5 \cdot 10^{-4} & \text{for } B_d \\ 2 \cdot 10^{-5} & \text{for } B_s, \end{cases}$$

because

$$\phi = \mathcal{O} \left(\frac{m_c^2}{m_b^2} \right) = \begin{cases} -0.1 = -5^\circ & \text{for } B_d \\ 3 \cdot 10^{-3} = 0.2^\circ & \text{for } B_s. \end{cases}$$

New physics can affect M_{12} lifting the GIM suppression of ϕ .

⇒ large enhancement of a_{fs} possible

⇒ Even crude upper bounds constrain models of new physics.

Laplace, Ligeti, Nir, Perez, hep-ph/0202010

No tagging is necessary: With $\Gamma[f, t] = \Gamma(B(t) \rightarrow f) + \Gamma(\bar{B}(t) \rightarrow f)$,

$$a_{\text{fs}}^{\text{untagged}} = \frac{\Gamma[f, t] - \Gamma[\bar{f}, t]}{\Gamma[f, t] + \Gamma[\bar{f}, t]} = \frac{a_{\text{fs}}}{2} \left[1 - \frac{\cos(\Delta m t)}{\cosh(\Delta\Gamma t/2)} \right]$$

Relation to $|q/p|$:

$$\left| \frac{q}{p} \right|^2 = 1 - a_{\text{fs}}$$

\Rightarrow The time evolution of any decay $B \rightarrow f$ can be used to determine a_{fs}

2001 world average of a_{fs} from $B_d \rightarrow X\ell^+\nu_\ell$ decays:

$$a_{\text{fs}} = (2 \pm 14) \cdot 10^{-3}$$

Nir

Babar, hep-ex/0311037, decays with one fully reconstructed B_d :

$$a_{\text{fs}}(B_d) = 2 \left(1 - \left| \frac{q}{p} \right| \right) \in [-104 \cdot 10^{-3}, -2 \cdot 10^{-3}] \quad @90\%CL$$

Any decay mode $B \rightarrow f$:

$$\lambda_f = \frac{q \langle f | \bar{B} \rangle}{p \langle f | B \rangle}$$

E.g. $\lambda_f \simeq -\exp(-2i\beta) + \mathcal{O}(a_{\text{fs}})$ for $B_d \rightarrow J/\psi K_S$.

Untagged decay rate (for $\Delta\Gamma = 0$):

$$\Gamma[f, t] \propto e^{-\Gamma t} \left[1 + \frac{a_{\text{fs}}}{2} \left(1 - \frac{1 - |\lambda_f|^2}{1 + |\lambda_f|^2} \cos(\Delta m t) + 2 \frac{\text{Im } \lambda_f}{1 + |\lambda_f|^2} \sin(\Delta m t) \right) \right]$$

So a_{fs} can be determined from tiny oscillatory terms in the untagged rate.

QCD corrections to a_{fs}

$$a_{fs} = \text{Im} \frac{\Gamma_{12}}{M_{12}}$$

Until recently: No QCD corrections to Γ_{12} were known.

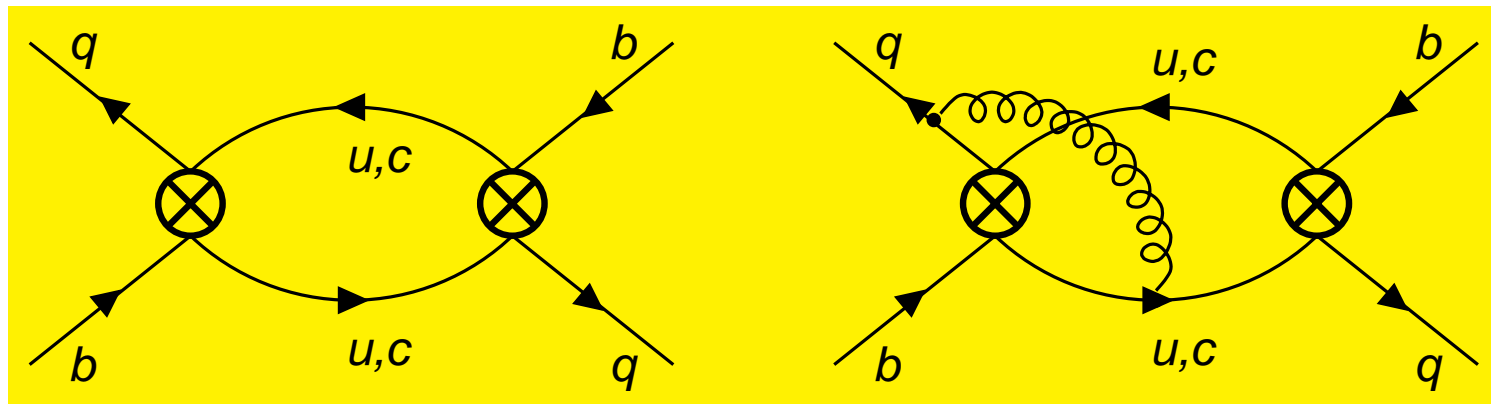
$a_{fs} \propto m_c^2 \Rightarrow$ large uncertainties from uncontrolled definition of m_c

The calculation of **next-to-leading order (NLO)** QCD corrections reduces the uncertainty from **50%** to **20%**.

Ciuchini, Franco, Lubicz, Mescia, Tarantino, hep-ph/0307344

Beneke, Buchalla, Lenz, UN, hep-ph/0308029

LO and NLO sample diagram:



Results including both **NLO QCD** and $1/m_b$ corrections:

$$a_{\text{fs}}(B_d) = 10^{-4} \left[-\frac{\sin \beta}{R_t} (12.0 \pm 2.4) + \left(\frac{2 \sin \beta}{R_t} - \frac{\sin 2\beta}{R_t^2} \right) (0.2 \pm 0.1) \right]$$

where β and R_t are one angle and one side of the **unitarity triangle**.

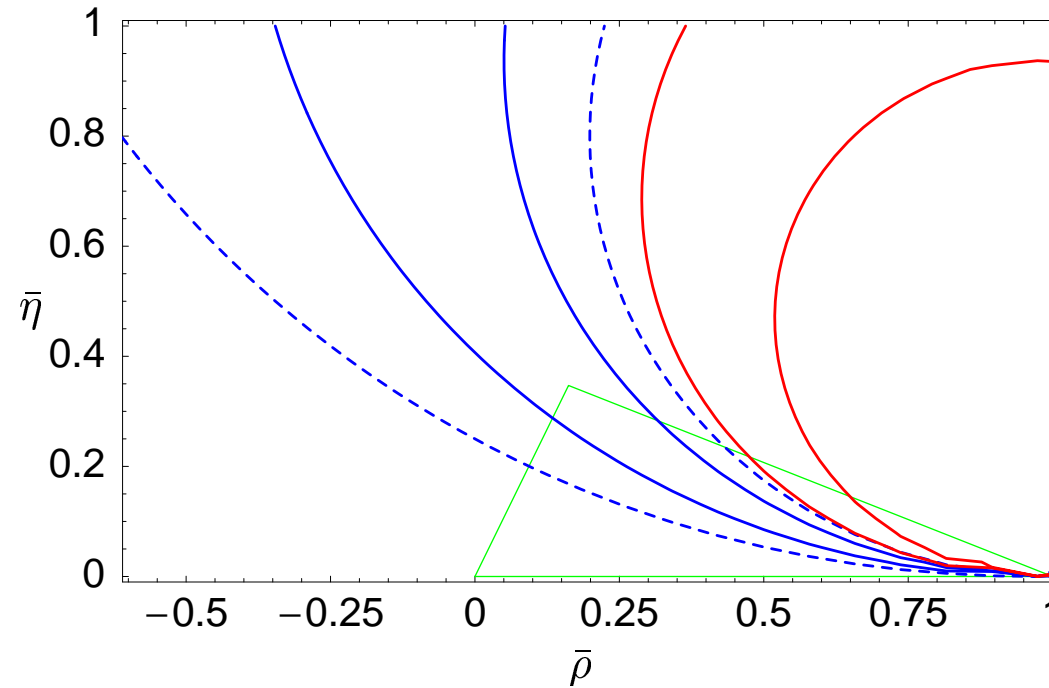
For the Standard Model fit to the unitarity triangle with $\beta = 22.4^\circ \pm 1.4^\circ$ and $R_t = 0.91 \pm 0.05$:

$$a_{\text{fs}}(B_d) = -(5.0 \pm 1.1) \cdot 10^{-4}$$

For the B_s system:

$$\begin{aligned} a_{\text{fs}}(B_s) &= (12.0 \pm 2.4) \cdot 10^{-4} |V_{us}|^2 R_t \sin \beta \\ &= (2.1 \pm 0.4) \cdot 10^{-5} \end{aligned}$$

Constraint on the $(\bar{\rho}, \bar{\eta})$ plane:



Solid curves: **NLO**

Area between red curves allowed for $a_{\text{fs}} = -10^{-3}$

Area between blue curves allowed for $a_{\text{fs}} = -5 \cdot 10^{-4}$

Area between dashed blue curves: **LO** range for $a_{\text{fs}} = -5 \cdot 10^{-4}$

As a by-product we obtain the **NLO QCD** (and $1/m_b$ corrections) to $\Delta\Gamma$:

$$\frac{\Delta\Gamma}{\Delta m} = (4.0 \pm 1.6) \cdot 10^{-3} \quad \text{for } B_d \text{ and } B_s$$

$$\frac{\Delta\Gamma}{\Gamma}(B_d) = (3.0 \pm 1.2) \cdot 10^{-3}$$

$\Delta\Gamma/\Delta m$ is proportional to the ratio $B_S/B = 1.4 \pm 0.2$ of two hadronic parameters which must be determined from lattice perturbation theory.

For $\Delta\Gamma/\Gamma$ in the B_s system **NLO QCD** corrections are known since 1998.

Beneke, Buchalla, Greub, Lenz, UN, hep-ph/9808385

Summary

Complete **NLO QCD** corrections are now known for the CP asymmetry in flavour-specific B decays and the width difference in $B-\bar{B}$ mixing.

In the Standard Model:

$$a_{\text{fs}}(B_d) = -(5.0 \pm 1.1) \cdot 10^{-4}$$

$$a_{\text{fs}}(B_s) = (2.1 \pm 0.4) \cdot 10^{-5}$$

$$\frac{\Delta\Gamma}{\Delta m} = (4.0 \pm 1.6) \cdot 10^{-3}$$

for B_d and B_s

$$\frac{\Delta\Gamma}{\Gamma}(B_d) = (3.0 \pm 1.2) \cdot 10^{-3}$$

a_{fs} can be obtained from the decay distribution of any untagged decay $B \rightarrow f$ for which λ_f is known.