

$e^+e^- \rightarrow J/\psi + \eta_c$ at B factory energy in Bethe-Salpeter framework

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BS framework as one way to investigate the **virtuality and **relative move.** of the compositio particles in bound for produc. process**

- Experimental results
- Theoretical explanations: consistent & in.
- Introduction on the BS framework
- The calculation procedure
- Disucssion on the investigation

Data and LO Calculations

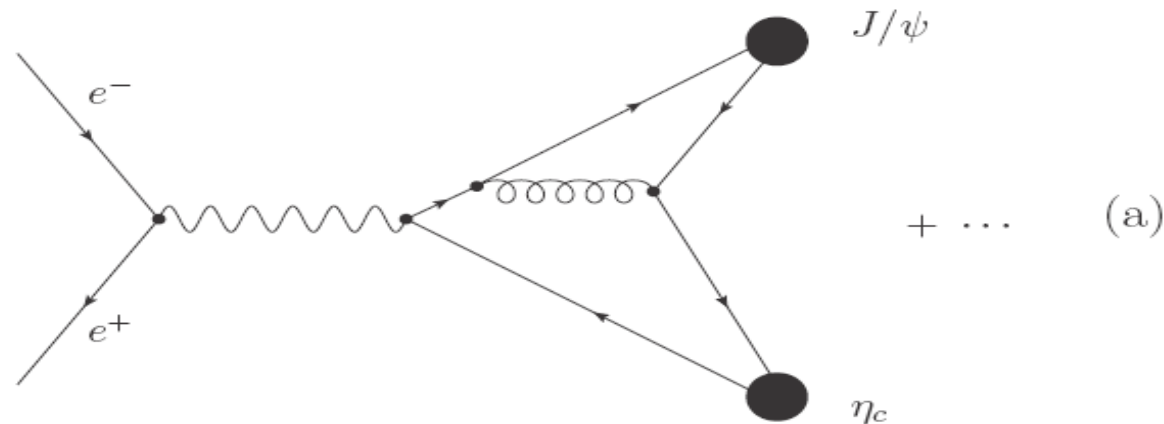
LO=leading order; leading opinion

$$\sigma(e^+ + e^- \rightarrow J/\psi + \eta_c) \times \mathcal{B}_{>2} = 25.6 \pm 2.8 \pm 3.4 \text{ fb (Belle)}$$

$$2.3 \sim 5.5 \text{ fb,}$$

Braaten & Lee,
Liu, He & Chao
Hagiwara, Kou & Qiao

$$\sigma(e^+ + e^- \rightarrow J/\psi + \eta_c) \times \mathcal{B}_{>2} = 17.6 \pm 2.8_{-2.1}^{+1.5} \text{ fb (BaBar),}$$

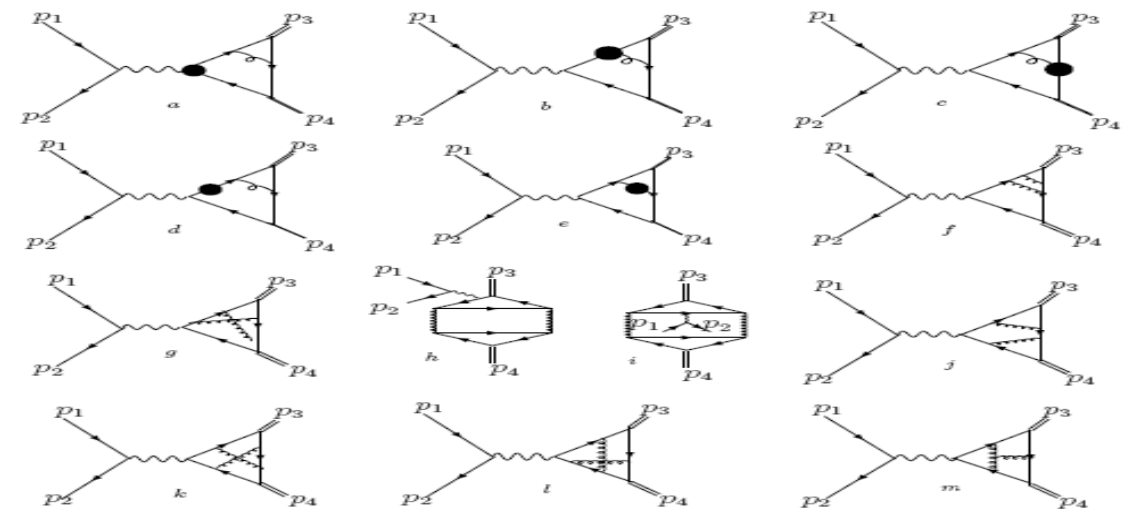


NRQCD

Higher order (NLO)

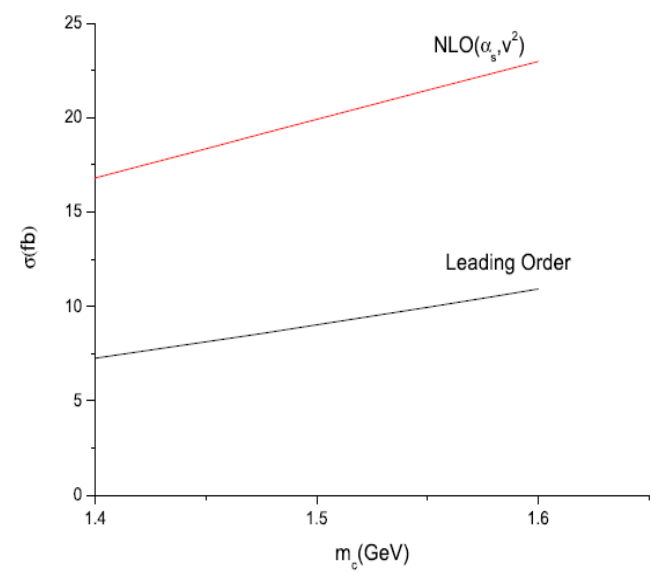
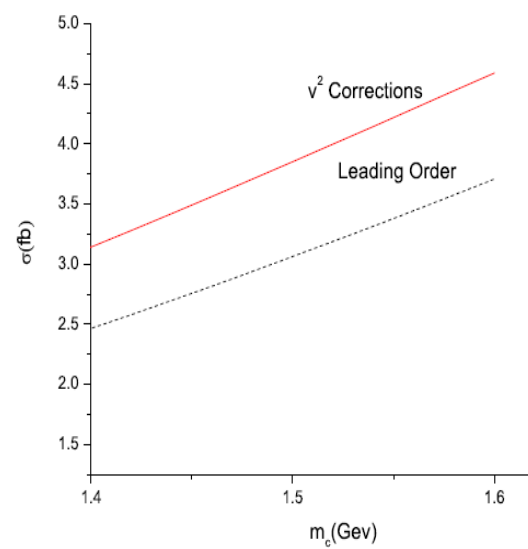
Zhang, Gao & Chao
Gong & Wang*

NRQCD



relativistic correction

He, Fan & Chao*
Bodwin, Lee & Yu

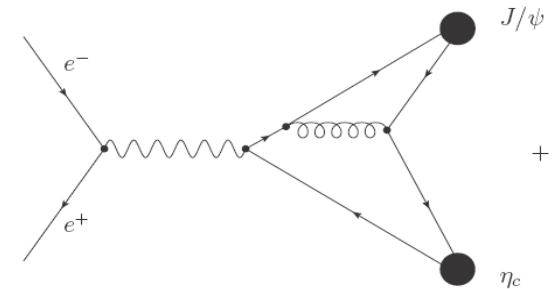


2012-3-12

Rencontres de Moriond (QCD)

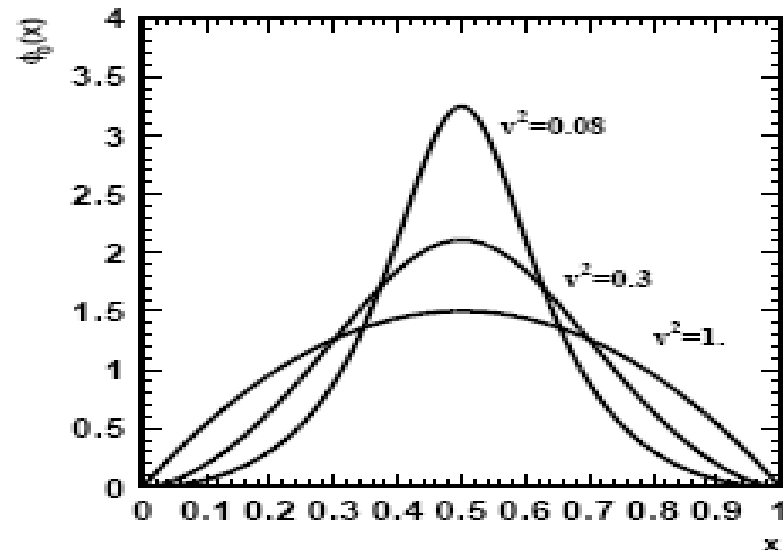
Relativistic approaches (on shell)

- Lightcone WF (Ma & Si, Bondar & Chernyak*)
- QCD sum rule (Sun, Wu, Zuo & Huang)
- BS (Guo, Ke, Li & Wu; Mengesha & Bhatnagar)

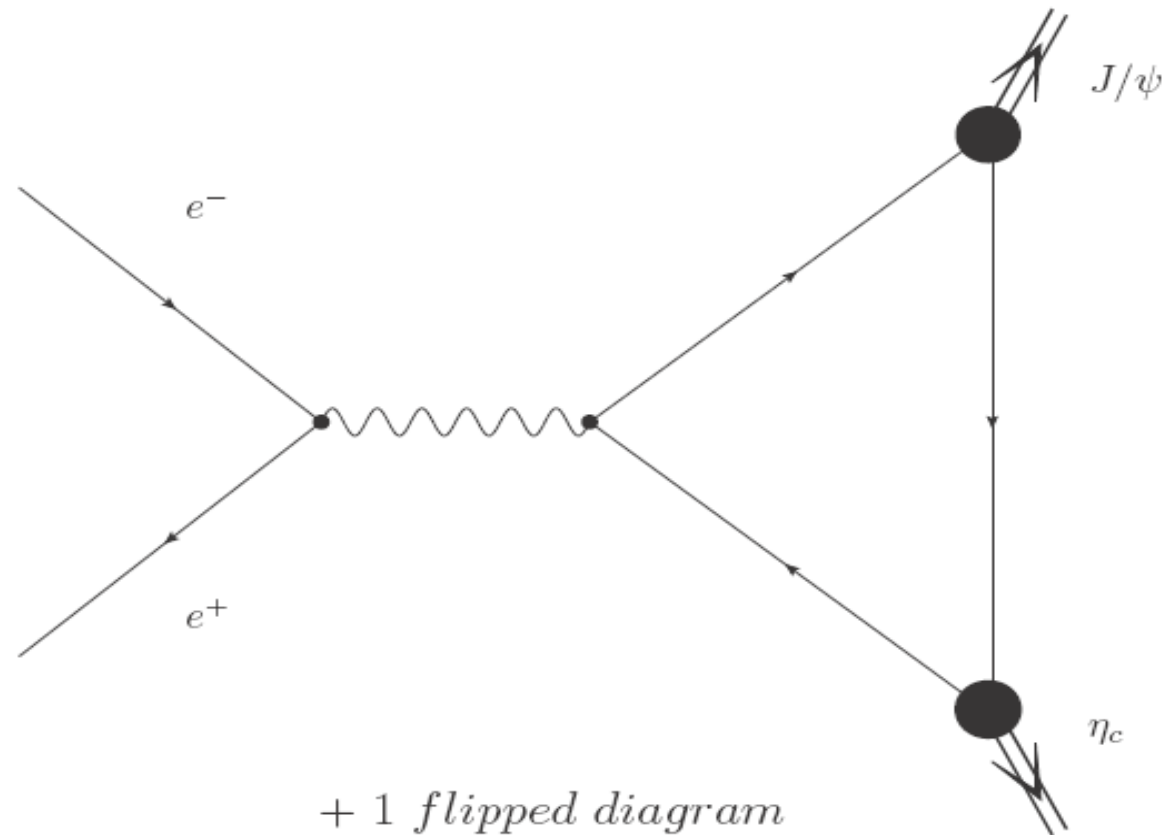


however!

They take the 'NR'
approximation
to some level
(heavy quark limit)



Both virtuality & relative move. taken into account, can start from the **lowest order**



Key points for Calculation

- **The vertices c-cbar-bound state**

the momentum distribution

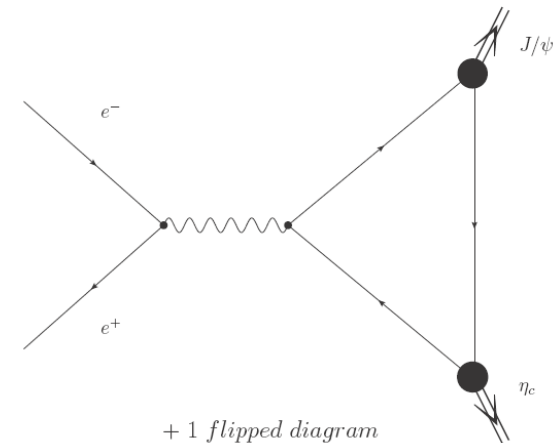
the Dirac structure

(both contain relativistic informations)

- **The loop integrals**

pole of the propagators

the other 3-dimensions



The coupling vertex and the BS WF

$$\chi(P, q) = \int \frac{d^4x}{(2\pi)^4} e^{-iqx} \frac{1}{\sqrt{3}} \delta_{ij} \langle 0 | T \psi^i\left(\frac{x}{2}\right) \bar{\psi}^j\left(-\frac{x}{2}\right) | B \rangle$$

$$= S_F(p_1) \Gamma(\hat{q}) S_F(-p_2)$$

by definition

$$\mathcal{L}_I = \bar{\psi}^i(x) \Theta_{ij} \left(i \frac{\partial}{\partial x} \right) \delta_{ij} \psi^j(x) B(x)$$

$$p_1 = \frac{P}{2} + q,$$

$$p_2 = \frac{P}{2} - q,$$

$$P = p_1 + p_2,$$

$$q = \frac{1}{2}(p_1 - p_2)$$

Reduction Formula

Some more details in CIA approx.

$$\hat{q}_\mu = q_\mu - \frac{q \cdot P}{P^2} \phi(\hat{q}) \sim e^{-\hat{q}^2/2\beta^2} \quad \begin{aligned} \Gamma(\hat{q}) &= D'(\hat{q})\phi(\hat{q}) \\ &= \tilde{\Gamma}ND(\hat{q})\phi(\hat{q}) \end{aligned}$$

First analyzed by Llewellyn-Smith:

$$\begin{aligned} \tilde{\Gamma}_J &= i\not{\epsilon}A_0 + \not{\epsilon}\not{P}\frac{A_1}{M} + (q\varepsilon - \not{\epsilon}\not{q})\frac{A_2}{M} + i(\not{\epsilon}\not{P}\not{q} - \not{\epsilon}\not{q}\not{P} + 2i(q\varepsilon)\not{P})\frac{A_3}{M^2} \\ &\quad + (q\varepsilon)\frac{A_4}{M} + iq\varepsilon\not{P}\frac{A_5}{M^2} - iq\varepsilon\not{q}\frac{A_6}{M^2} + (q\varepsilon)(\not{P}\not{q} - \not{q}\not{P})\frac{A_7}{M^3} \end{aligned}$$

$$\tilde{\Gamma}_y = \gamma_5 B_0 - i\gamma_5\not{P}\frac{B_1}{M} - i\gamma_5\not{q}\frac{B_2}{M} - \gamma_5(\not{P}\not{q} - \not{q}\not{P})\frac{B_3}{M^2}$$

see, Bhatnagr & LSY

Successful in describing decays

	f_P^0	f_P^1	f_P^2	f_P^3	$ f_P^{LO} $	$ f_P^{NLO} $	$f_P^{LO}(\%)$	$f_P^{NLO}(\%)$	$f_P = f_P^{LO} + f_P^{NLO}$
π	0.110	-0.154	0.000	0.175	0.044	0.175	25%	75%	0.130
K	0.202	-0.104	0.025	0.039	0.098	0.064	60%	40%	0.164
D	0.271	-0.097	0.010	0.009	0.174	0.019	90%	10%	0.194
D_S	0.426	-0.156	0.013	0.013	0.270	0.026	91%	9%	0.296
B	0.345	-0.125	0.005	0.003	0.220	0.008	96%	4%	0.228

	f_π	f_K	f_D	f_{D_S}	f_B
BSE (3.5% average error) present paper	0.130	0.164	0.194	0.296	0.228
BSE [5]				0.248	
SDE [6]		0.164			
Lattice [15]			0.208 ± 0.004	0.241 ± 0.003	
QCD-SR [16]			0.20 ± 0.02	0.23 ± 0.02	
Exp. Results [13]	0.1300 ± 0.0001	0.159 ± 0.001	0.22 ± 0.02	0.29 ± 0.03	
Babar+Belle Collaboration [14]					0.24 ± 0.04

see Bhatnagar, LSY & Mahecha
Rencontres de Moriond (QCD)

$$\begin{aligned}\sigma &= \int \frac{|\mathbf{M}|^2}{F} dR_2 \\ &= \int \frac{|\mathbf{M}|^2}{2s} \frac{k}{32\pi^2 E} 4E^2 (2\pi)^6 \sin\theta d\theta d\varphi\end{aligned}$$

$$\begin{aligned}\mathbf{M} &= 2l_\mu^{(e)} \frac{ie^2 e_Q}{q^2} \frac{1}{-(2\pi)^2} N_J N_y \int \frac{d^4 p}{(2\pi)^4} \frac{D(\hat{q}_J) D(\hat{q}_y) \phi(\hat{q}_J) \phi(\hat{q}_y)}{((p - P_y)^2 - m^2)(p^2 - m^2)((p + P_J)^2 - m^2)} \\ &\quad \times \text{Tr}[\gamma^\mu (\not{p} - \not{P}_y + m) \Gamma_y (\not{p} + m) \Gamma_J (\not{p} + \not{P}_J + m)]\end{aligned}\quad (3.20)$$

$$N_J^{-2}(\beta) = \int_0^\infty d\hat{q} 16\pi^3 e^{-\frac{\hat{q}^2}{\beta^2}} \hat{q}^2 (m^2 + \hat{q}^2) (-M^2 + 4(m^2 + \hat{q}^2))^2 A$$

$$N_y^{-2}(\beta) = \int_0^\infty d\hat{q} 16\pi^3 e^{-\frac{\hat{q}^2}{\beta^2}} \hat{q}^2 (m^2 + \hat{q}^2) (-M^2 + 4(m^2 + \hat{q}^2))^2 B$$

$$\begin{aligned}
A = & \frac{2\pi}{3(m^2 + \hat{q}^2)^{1.8}} + \frac{4m^2\pi}{3M^2(m^2 + \hat{q}^2)^{1.8}} - \frac{2\pi}{3M^2(m^2 + \hat{q}^2)^{0.8}} \\
& + \frac{4m^2\pi(M^2 - 12(m^2 + \hat{q}^2))}{3(m^2 + \hat{q}^2)^{1.8}(M^2 - 4(m^2 + \hat{q}^2))^2} + \frac{M^2\pi(M^2 - 12(m^2 + \hat{q}^2))}{3(m^2 + \hat{q}^2)^{1.8}(M^2 - 4(m^2 + \hat{q}^2))^2} \\
& + \frac{4\pi\left(-\frac{M^2}{2} + (m^2 + \hat{q}^2)\right)}{3M^2(m^2 + \hat{q}^2)^{1.8}} + \frac{4\pi}{(m^2 + \hat{q}^2)^{0.8}(-M^2 + 4(m^2 + \hat{q}^2))} \\
& - \frac{16m^2\pi}{3M^2(m^2 + \hat{q}^2)^{0.8}(-M^2 + 4(m^2 + \hat{q}^2))} \tag{3.24}
\end{aligned}$$

$$\begin{aligned}
B = & -\frac{1}{M^2}\left(\frac{\pi(-2m^2M^2 - \frac{M^4}{2} - 6M^2\hat{q}^2)(-M^2(m^2 + \hat{q}^2)^{0.8} + 12(m^2 + \hat{q}^2)^{1.8})}{2(m^2 + \hat{q}^2)^2(M^2 - 4(m^2 + \hat{q}^2))^2}\right. \\
& + 2\pi(4m^2M^2 - M^4 - 4M^2\hat{q}^2)\left(-\frac{M^2}{8(m^2 + \hat{q}^2)^{1.8}(M^2 - 4(m^2 + \hat{q}^2))^2}\right. \\
& \left. + \frac{5}{2(m^2 + \hat{q}^2)^{0.8}(M^2 - 4(m^2 + \hat{q}^2))^2}\right) + 8M^4\pi\left(-\frac{0.25}{M^6(M^2 - 2M(m^2 + \hat{q}^2)^{0.8})^2}\right. \\
& \left. + \frac{M^2 + 2M(m^2 + \hat{q}^2)^{0.8}}{8M^6(m^2 + \hat{q}^2)^{0.8}} + \frac{M^4 - 2M^2(m^2 + \hat{q}^2)}{16M^6(m^2 + \hat{q}^2)^{1.8}}\right) \\
& + \frac{1}{4}\pi\left(\frac{2(M^2 + 2M(m^2 + \hat{q}^2)^{0.8})^2}{M^2(m^2 + \hat{q}^2)^{0.8}(M + 2(m^2 + \hat{q}^2)^{0.8})^2}\right. \\
& + \frac{2M(-m + M)(M^2 - 2M(m^2 + \hat{q}^2)^{0.8})^2}{(m^2 + \hat{q}^2)(-(m^2 + \hat{q}^2) + (M - (m^2 + \hat{q}^2)^{0.8})^2)^2} \\
& - \frac{4M^2(M^2 - 2M(m^2 + \hat{q}^2)^{0.8})}{(m^2 + \hat{q}^2)(-(m^2 + \hat{q}^2) + (M - (m^2 + \hat{q}^2)^{0.8})^2)} \\
& \left. - \frac{M(M^2 - 2M(m^2 + \hat{q}^2)^{0.8})^2}{(m^2 + \hat{q}^2)^{1.8}(-(m^2 + \hat{q}^2) + (M - (m^2 + \hat{q}^2)^{0.8})^2)}\right) \tag{3.25}
\end{aligned}$$

$$| \mathbf{M} |^2 = 4L_{\mu\nu}^{(e)} L^{\mu\nu}_{(h)} \frac{e^4 e_Q^2}{(2\pi)^4 q^4} N_J^2 N_y^2 | I(\beta) |^2$$

$$I(\beta) = \int \frac{d^4 p}{(2\pi)^4} \frac{D(\hat{q}_J) D(\hat{q}_y) \phi(\hat{q}_J, \beta) \phi(\hat{q}_y, \beta)}{((p - P_y)^2 - m^2)(p^2 - m^2)((p + P_J)^2 - m^2)}$$

$$I(\beta) = \frac{1}{2(2\pi)^3} \int_0^{+\infty} d\vec{p}_\perp^2 \int_{-\infty}^{+\infty} dp^- \int_{-\infty}^{+\infty} dp^+ F(p^+, p^-, \vec{p}_\perp^2, \beta) \frac{1}{2p^+ p^- - \vec{p}_\perp^2 - m^2}$$

$$\times \frac{1}{2(p^+ + P_J^+)(p^- + P_J^-) - \vec{p}_\perp^2 - m^2} \frac{1}{2(p^+ - P_y^+)(p^- - P_y^-) - \vec{p}_\perp^2 - m^2}$$

$$I(\beta) = I_1(\beta) + I_2(\beta)$$

$$I_1(\beta) = \frac{1}{2(2\pi)^3} \int_0^{+\infty} d\vec{p}_\perp^2 \int_0^{P_y^-} dp^- \int_{-\infty}^{+\infty} dp^+$$

$$I_2(\beta) = \frac{1}{2(2\pi)^3} \int_0^{+\infty} d\vec{p}_\perp^2 \int_{-P_J^-}^0 dp^- \int_{-\infty}^{+\infty} dp^+$$

β	I^2	N^2	I^2/N^2	$\sigma(fb)$
1/2	3.39×10^{-14}	1.29×10^6	2.63×10^{-20}	1.4×10^{-6}
$1/\sqrt{2}$	7.90×10^{-9}	1.47×10^7	5.37×10^{-16}	2.9×10^{-2}
1	1.36×10^{-6}	1.52×10^8	8.95×10^{-15}	4.9×10^{-1}
1.15	8.00×10^{-6}	2.83×10^8	2.83×10^{-14}	1.5
1.3	9.80×10^{-6}	8.06×10^7	1.22×10^{-13}	6.6
$\sqrt{2}$	9.80×10^{-5}	1.08×10^9	9.09×10^{-14}	5.0
1.5	1.00×10^{-4}	3.36×10^9	2.97×10^{-14}	1.6
1.6	3.59×10^{-4}	8.85×10^9	4.05×10^{-14}	2.2
1.7	1.08×10^{-3}	1.89×10^{10}	5.73×10^{-14}	3.1
1.8	2.95×10^{-3}	3.40×10^{10}	8.69×10^{-14}	4.7
1.9	7.32×10^{-3}	4.96×10^{10}	1.48×10^{-13}	8.1
2	1.05×10^{-2}	4.72×10^{10}	2.23×10^{-13}	12.2
2.1	3.41×10^{-2}	2.50×10^{10}	1.37×10^{-12}	74.6
2.25	8.59×10^{-2}	5.73×10^{11}	1.50×10^{-13}	8.2
2.5	1.81×10^{-1}	6.59×10^{12}	2.75×10^{-14}	1.5
3	1.41	2.08×10^{14}	6.81×10^{-15}	0.4
4	26.35	3.07×10^{16}	8.57×10^{-16}	0.05

Conclusion

- Taking into account the virtual quark state in the bound state in a full relativistic framework
- The lowest order contributing sounding cross section
- A consistent framework to higher precision (higher order) and for global analysis