

HADRON PHYSICS AT KLOE AND KLOE-2

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KLOE data allow to study many interesting processes related to light mesons. Several items have been investigated, among them we present the recently published search for the U boson produced with η meson and the last results on the dynamics of the $\eta \rightarrow \pi^+\pi^-\gamma$ decay. The KLOE-2 project aims to extend the KLOE program with detector upgrades and increased statistics: we describe the status of the art.

1 KLOE and DAΦNE

The e^+e^- collider DAΦNE, designed to operate at the center of mass energy $\sqrt{s} \simeq 1.02$ GeV, the ϕ meson mass, has delivered to the KLOE experiment an integrated luminosity of about $2.5fb^{-1}$ on peak of the ϕ meson and also about $240pb^{-1}$ at $\sqrt{s} \simeq 1$ GeV. The KLOE detector consists of a large volume cylindrical drift chamber¹ (DC), 3.3 m length and 2 m radius, surrounded by a calorimeter² (EMC) made of lead and scintillating fibers. A superconducting coil produces an axial field $B = 0.52$ T. In the DC charged particle momenta are reconstructed with resolution $\sigma_p/p \simeq 0.4\%$, while in the EMC energy clusters are reconstructed grouping calorimeter cells close in space and in time with energy and time resolution of $\sigma_E/E = 5.7\%/\sqrt{E(GeV)}$ and $\sigma_t = 57ps/\sqrt{E(GeV)} \oplus 100$ ps.

1.1 U Boson Search: $\phi \rightarrow \eta e^+e^-/\eta U$

In recent years several unexpected astrophysical observations have failed to find a common interpretation in terms of standard astrophysical or particle sources. All these unexpected observations^a can be interpreted assuming the existence of a light hidden sector interacting with Standard Model particles through the mixing (ϵ) between a new gauge vector boson U , with mass lighter than $O(\text{GeV})$, and the photon. The U boson can be produced at e^+e^- colliders via different processes³, we present the analysis of the process $\phi \rightarrow \eta U$, where the η meson is tagged by the $\eta \rightarrow \pi^+\pi^-\pi^0$ channel. The Dalitz decay $\phi \rightarrow \eta l^+l^-$, having the same signature, is an irreducible background for the U boson search. The SND⁴ and CMD-2⁵ collaborations measured the branching fraction of $BR(\phi \rightarrow \eta e^+e^-) = O(10^{-4})$, which corresponds to a cross section of $\sigma(\phi \rightarrow \eta l^+l^-) \sim 0.7$ nb. For the signal the expected cross section is $\sigma(\phi \rightarrow \eta U) \sim 40$ fb, in the hypothesis of a mixing $\epsilon = 10^{-3}$ and a $\phi\eta\gamma^*$ transition form factor $|F_{\phi\eta}(m_U^2)|^2 = 1$. Even though the ratio between the overall cross section of the $\phi \rightarrow \eta U$ and $\phi \rightarrow \eta l^+l^-$ is not favorable to the signal, the di-lepton invariant mass should be different allowing to test the kinetic mixing parameter ϵ . We searched for the U boson in the e^+e^- final state, because the

^aFor an exhaustive list of references see the KLOE-2 paper³

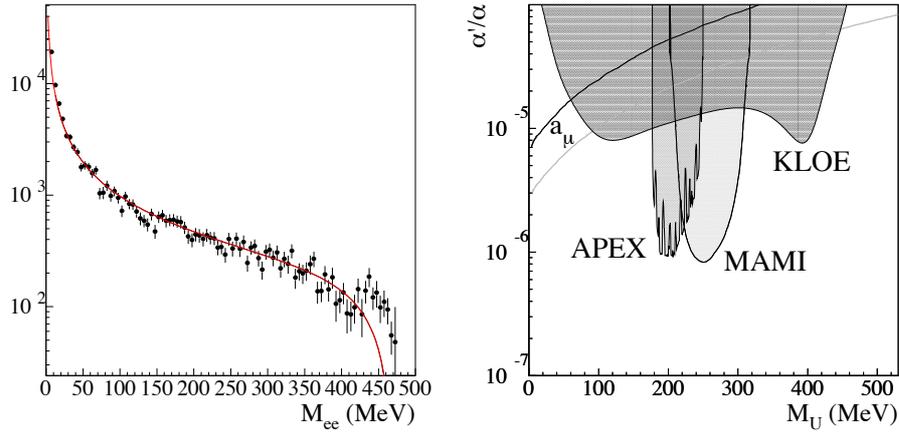


Figure 1: Left: Fit to the corrected M_{ee} spectrum for the Dalitz decays $\phi \rightarrow \eta e^+ e^-$; Right: Exclusion plot at 90% C.L. for the parameter $\alpha'/\alpha = \epsilon^2$, compared with existing limits in our region of interest.

channel $U \rightarrow e^+ e^-$ allow to search the U boson in a wider mass range and the e^\pm are easily identified using TOF technique.

The M_{ee} spectrum has been studied using an integrated luminosity of 1.5 fb^{-1} of ϕ decays: about 14,000 $\phi \rightarrow \eta e^+ e^-$, $\eta \rightarrow \pi^+ \pi^- \pi^0$ candidates are present in the analyzed data set, with a negligible background contamination.

The background shape is extracted directly from our data. A fit is performed to the M_{ee} distribution, after a bin-by-bin subtraction of $\phi \rightarrow \eta \gamma$ background and efficiency correction, using the following parametrization as from Vector Meson Dominance (VMD) model⁶

$$\frac{d\Gamma(\phi \rightarrow \eta e^+ e^-)}{dq^2} = \frac{\alpha}{3\pi} \frac{|F_{\phi\eta}(q^2)|^2}{q^2} \sqrt{1 - \frac{4m^2}{q^2}} \left(1 + \frac{2m^2}{q^2}\right) \lambda^{3/2}(m_\phi^2, m_\eta^2, m_U^2) \quad (1)$$

with $q = M_{ee}$ and the transition form factor described by:

$$F_{\phi\eta}(q^2) = \frac{1}{1 - q^2/\Lambda^2} \quad (2)$$

Free parameters of the fit are Λ and an overall normalization factor. A good description of the M_{ee} shape is obtained except at the high end of the spectrum see fig.(1.left), because of the contamination of a residual background from multi-pion events. The $\phi \rightarrow \eta U$ Monte Carlo signal has been produced according to Reece-Wang model⁷, with a flat distribution of the U mass, M_U . The sample has been used to evaluate the resolution on the $e^+ e^-$ invariant mass as a function of M_U : resolution is $\sim 2 \text{ MeV}$ for $M_U < 350 \text{ MeV}$ and then improves to 1 MeV for higher values. The upper limit on $\phi \rightarrow \eta U$ as a function of M_U has been reported in terms of the kinetic mixing parameter $\epsilon^2 = \alpha'/\alpha$, where α' is the coupling of U boson to electrons and α is the fine structure constant. We include the opening $U \rightarrow \mu^+ \mu^-$ threshold, in the hypothesis that the U boson decay only to lepton pairs and assuming equal coupling to $e^+ e^-$ and $\mu^+ \mu^-$. The smoothed exclusion plot at 90% C.L. on α'/α , see fig.(1.right), is compared with existing limits from the muon anomalous magnetic moment a_μ and from recent measurement of MAMI⁸ and APEX⁹. The gray line is where the U boson parameters should lay to account for the observed discrepancy between measured and calculated a_μ values. Our result improves existing limits in a wide mass range, resulting in an U.L. on $\alpha'/\alpha \leq 2 \times 10^{-5}$ @ 90% C.L. for $50 < M_U < 420 \text{ MeV}$. Our result excludes that the existing a_μ discrepancy is due to U boson with mass ranging between 90 and 450 MeV. Preliminary study for U -boson search looking at

$\eta \rightarrow \gamma\gamma$ and $\eta \rightarrow 3\pi^0$ looks promising and they should in principle allow to improve U.L. by a factor 2.

1.2 Light Mesons: $\eta \rightarrow \pi^+\pi^-\gamma$

The decays $\eta \rightarrow \pi^+\pi^-\gamma$ and $\eta' \rightarrow \pi^+\pi^-\gamma$ are expected to get contribution from the anomaly accounted for by the Wess Zumino Witten (WZW) term into the ChPT Lagrangian¹⁰. Those anomalous processes are referred to as box anomalies which proceed through a vector meson resonant contribution, described by VMD. According to effective theory¹⁰ the contribution of the direct term should be present together with VMD. In case of $\eta \rightarrow \pi^+\pi^-\gamma$ the ρ contribution is not dominant, this makes the partial width sensitive to the presence of the direct term. Recently CLEO¹² has measured the ratio $R_\eta = \Gamma(\eta \rightarrow \pi^+\pi^-\gamma)/\Gamma(\eta \rightarrow \pi^+\pi^-\pi^0) = 0.175 \pm 0.007_{stat} \pm 0.006_{syst}$, which differs by more than 3σ from the average result¹⁵ of previous measurements^{13,14}, $R_\eta = 0.207 \pm 0.004$. We present a preliminary measurement with the highest statistics and the smallest systematic error ever achieved.

The final state under study is $\pi^+\pi^-\gamma\gamma$, since at KLOE, the η mesons are produced together with a monochromatic recoil photon ($E_\gamma = 363$ MeV) through the radiative decay $\phi \rightarrow \eta\gamma$. In the considered data sample there are about $\simeq 25 \times 10^6$ η 's. The main background comes from $\phi \rightarrow \pi^+\pi^-\pi^0, \pi^0 \rightarrow \gamma\gamma$ decaying to the same final state. Other backgrounds are $\phi \rightarrow \eta\gamma, \eta \rightarrow \pi^+\pi^-\pi^0 \rightarrow \pi^+\pi^-3\gamma$ with one photon lost, and $\phi \rightarrow \eta\gamma, \eta \rightarrow e^+e^-\gamma$ when both electrons are mis-identified as pions. The process $\phi \rightarrow \eta\gamma$ with $\eta \rightarrow \pi^+\pi^-\pi^0$ represents a good control sample, due to the similar topology. Moreover the ratio $\Gamma(\eta \rightarrow \pi^+\pi^-\gamma)/\Gamma(\eta \rightarrow \pi^+\pi^-\pi^0)$ is not affected by the uncertainties on the luminosity, the $\phi \rightarrow \eta\gamma$ partial width and the ϕ production cross section cancel in the ratio. We use the same preselection as for the $\eta \rightarrow \pi^+\pi^-\gamma$ signal. Concerning the control sample we select $N(\eta \rightarrow \pi^+\pi^-\pi^0) = 1190 \cdot 10^3$, with a selection efficiency of $\varepsilon = 0.2277 \pm 0.0002$ and a background contamination of 0.65%; concerning the signal we select $N(\eta \rightarrow \pi^+\pi^-\gamma) = 204950 \pm 450$ with $\varepsilon = 0.2131 \pm 0.0004$ and a background contamination of 10%. Combining our results we obtain the ratio:

$$R_\eta = \frac{\Gamma(\eta \rightarrow \pi^+\pi^-\gamma)}{\Gamma(\eta \rightarrow \pi^+\pi^-\pi^0)} = 0.1856 \pm 0.0005_{stat} \pm 0.0028_{syst} \quad (3)$$

Our measurement is in agreement with the most recent result from CLEO¹², $R_\eta = 0.175 \pm 0.007_{stat} \pm 0.006_{syst}$. Combining our measurement with the world average value¹¹ $\Gamma(\eta \rightarrow \pi^+\pi^-\pi^0) = (295 \pm 16)$ eV, we get $\Gamma(\eta \rightarrow \pi^+\pi^-\gamma) = (55 \pm 3)$ eV, which is in agreement with the value expected taking into account the direct term¹⁰, providing a strong evidence in favor of the box anomaly.

The $M_{\pi^+\pi^-}$ dependence of decay width has been parametrized in different approaches, in which VMD has been implemented in effective Lagrangians^{10,16}. Recently a model independent method, based on ChPT and dispersive analysis, has been developed¹⁷. In this approach, the relative strength between tree level and resonance contribution are not fixed. The function proposed to describe the partial width as function of $s_{\pi\pi} = m_{\pi\pi}^2$ is the following:

$$\frac{d\Gamma(\eta \rightarrow \pi^+\pi^-\gamma)}{ds_{\pi\pi}} = |AP(s_{\pi\pi})F(s_{\pi\pi})|^2 \Gamma_0(s_{\pi\pi}) \quad (4)$$

where A is a normalization factor; $\Gamma_0(s_{\pi\pi}) = \frac{1}{3 \cdot 2^{11} \cdot \pi^3 m_\eta^3} (m_\eta^2 - s_{\pi\pi})^3 s_{\pi\pi} \sigma(s_{\pi\pi})^3$ with $\sigma(s_{\pi\pi}) = \sqrt{1 - 4m_\pi^2/s_{\pi\pi}}$; $F_V(s_{\pi\pi})$ is the pion vector form factor, the function $P(s_{\pi\pi}) = 1 + \alpha s_{\pi\pi}$ is reaction specific. For more details see Stoll's paper¹⁷.

The α parameter was measured also by the WASA@COSY collaboration¹⁸: $\alpha = (1.89 \pm 0.25_{stat} \pm 0.59_{syst} \pm 0.02_{th})$ GeV⁻². Fig.2 shows the observed $M_{\pi^+\pi^-}$ spectrum, background

subtracted, compared with the theoretical prediction of eq.(4) with the value $\alpha = (1.31 \pm 0.08_{stat} \pm 0.40_{syst} \pm 0.02_{th})$ coming as output of the fit to the $M_{\pi\pi}$ shape, corrected for acceptance and smearing.

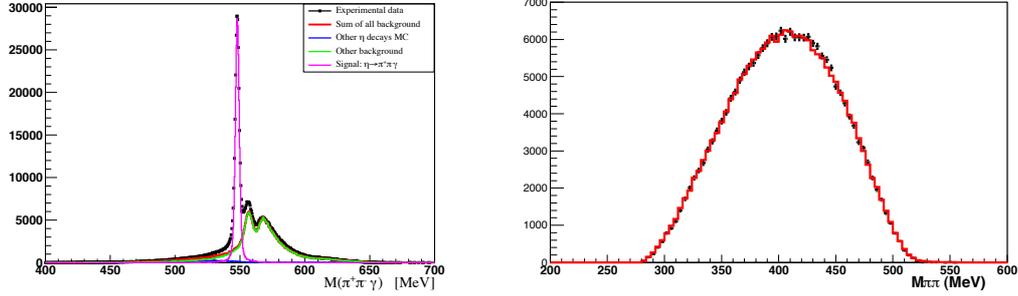


Figure 2: Left: The $\pi^+\pi^-\gamma\eta$ invariant mass distribution: Data-MC comparison. Dots are data, Magenta is MC signal $\eta \rightarrow \pi^+\pi^-\gamma$, Red is all MC background contribution; Right: Measured spectrum $m_{\pi\pi}$ (dots); histogram is the prediction from eq.(4) with α as from the output of the fit, corrected for acceptance and experimental resolution

1.3 KLOE-2

High statistic samples of light mesons produced at KLOE allowed to perform precision measurement and to look for very rare decays. A new DAΦNE interaction region, with large beam crossing angle and sextupoles for crab waist, improved the performance of the collider: a factor 3 in the luminosity has been gained. Minimal detector upgrade for first KLOE-2 run are already available as taggers to detect momentum of leptons in $e^+e^- \rightarrow e^+e^-\gamma^*\gamma^* \rightarrow e^+e^-X$. Work is in progress to insert an Inner Tracker, a 4 layers of cylindrical triple GEM, and new calorimeters around the beam pipe to increase acceptance for γ 's from interaction point. Nowadays a new data taking at KLOE-2 is waiting for stable run conditions.

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