

# Pati–Salam and lepton universality in B decays

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based on **JHEP 1812 (2018) 103** with **Daniele Teresi**

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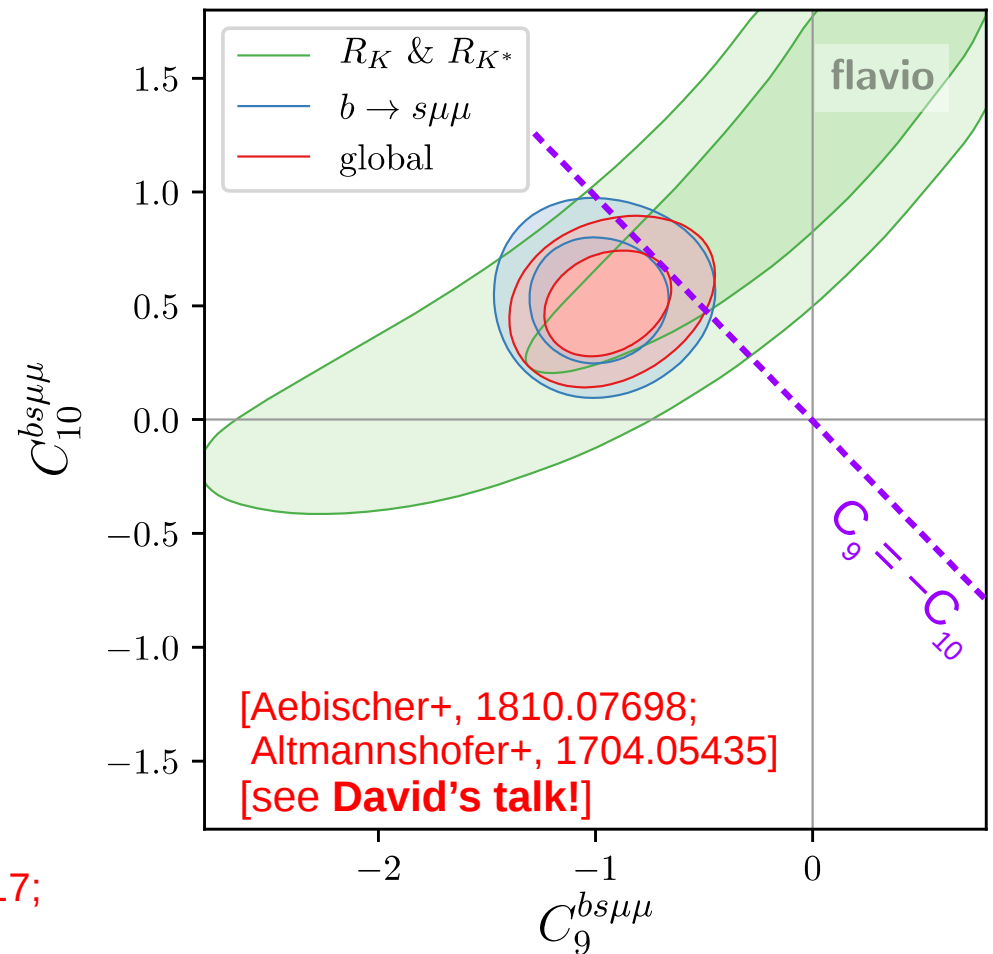
$$R_{K^{(*)}} \equiv \frac{\Gamma(B \rightarrow K^{(*)} \mu\mu)}{\Gamma(B \rightarrow K^{(*)} ee)}$$

- Neutral current, *loop-level* SM.
- **LHCb** anomaly.
- Good operator ( $C_9 = -C_{10}$ )

$$\frac{1}{(31 \text{ TeV})^2} \bar{s} \gamma_\alpha P_L b \bar{\mu} \gamma^\alpha P_L \mu.$$

[Capdevila+, '17; Altmannshofer+, '17; Geng+, '17; Ciuchini+, '17; D'Amico+, '17;...]

- Tree-level UV: **Z'** or **leptoquark**. [[arxiv.org/list/hep-ph/](https://arxiv.org/list/hep-ph/)]
- High NP scale = heavy new boson = easy to have.
- Also look for  $B_s - \bar{B}_s$  and di-muons at LHC.



$$R_{D^{(*)}} \equiv \frac{\Gamma(\bar{B} \rightarrow D^{(*)} \tau \nu)}{\Gamma(\bar{B} \rightarrow D^{(*)} \ell \nu)}$$

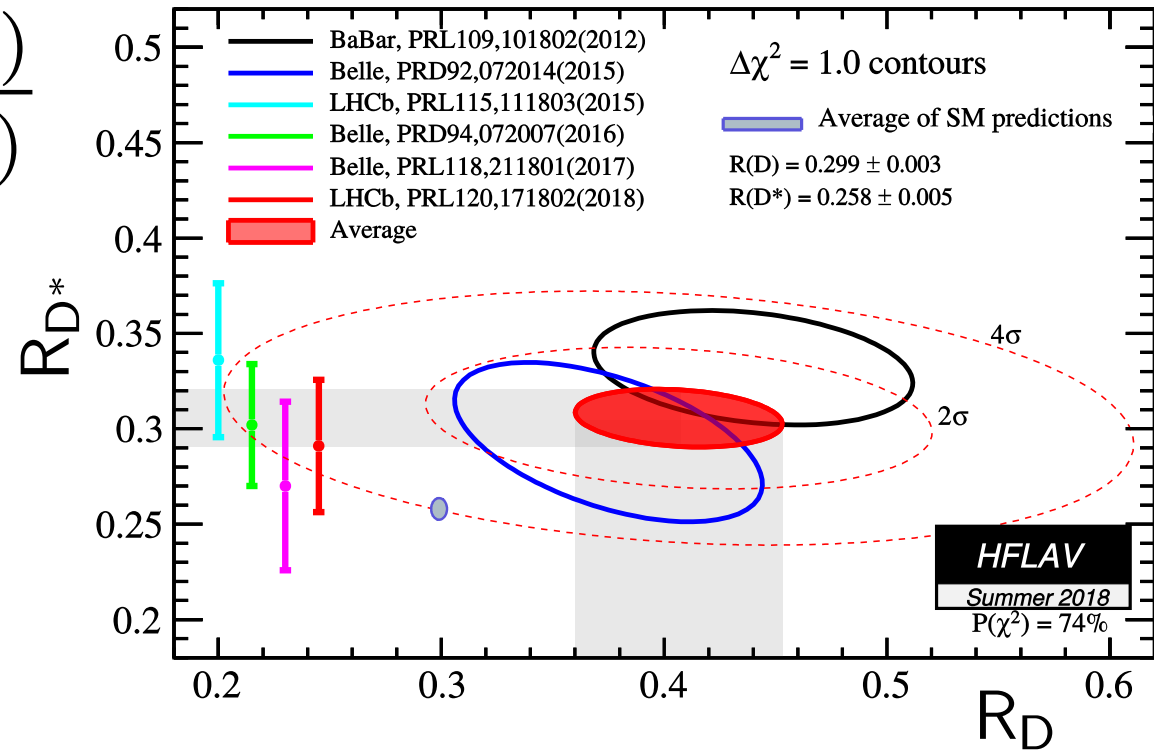
- Charged current, *tree-level* SM.
- LHCb, BaBar, Belle.
- Good operator

$$\frac{1}{(2.4 \text{ TeV})^2} \bar{c} \gamma_\alpha P_L b \bar{\tau} \gamma^\alpha P_L \nu.$$

[HFLAV; Bernlochner+, '17; Di Luzio+, '17;...]

- ~~$H^+$~~ ,  $W'$  or leptoquark.
- Low scale, many constraints:  $B \rightarrow K \nu \nu$ ,  $B_c \rightarrow \tau \nu$ .

[Li+, 1605.09308; Alonso+, 1611.06676]



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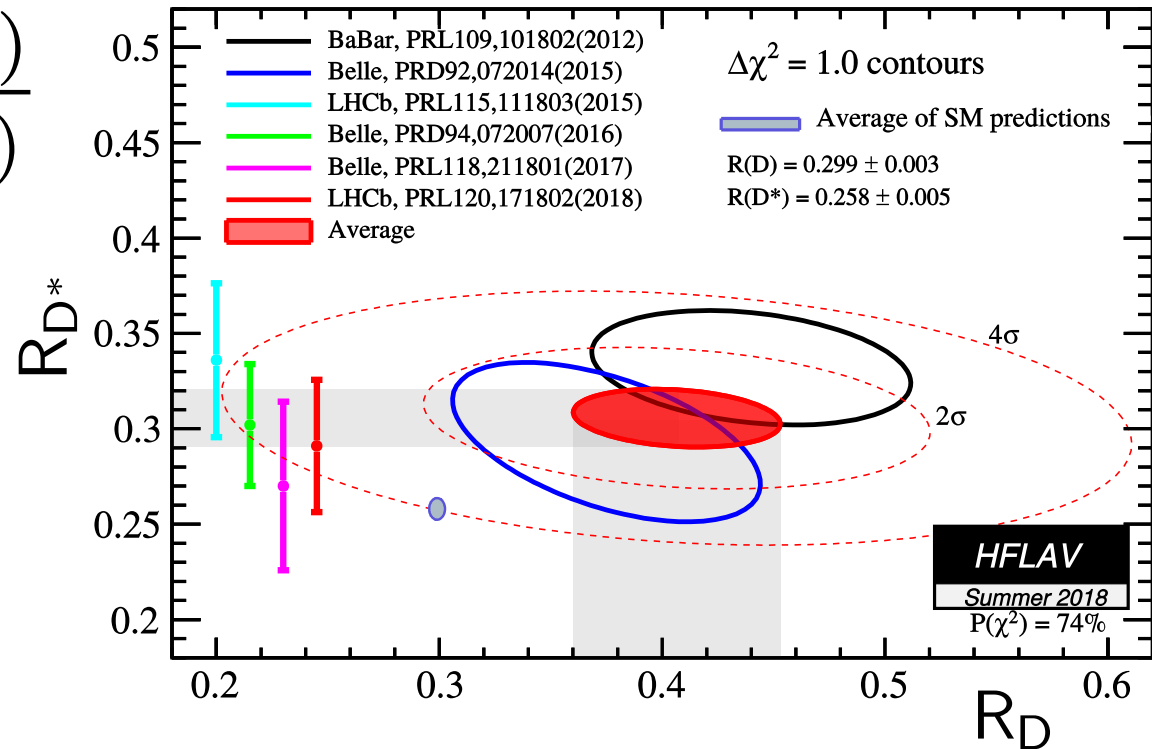
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- ~~H<sup>+</sup>~~, W' or leptoquark.

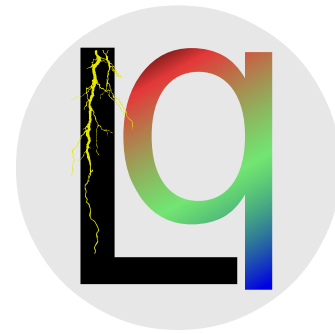
- Low scale, many constraints:  $B \rightarrow K \nu \nu$ ,  $B_c \rightarrow \tau \nu$ .
- Light (< 100 MeV) sterile neutrino? Evades  $B \rightarrow K \nu \nu$ .



or right-handed neutrino!

[He, Valencia, '12/'17; Greljo+, '18; Asadi+, '18; Robinson+, '18; Bečirević+, '16; Azatov+, '18]

# Leptoquarks?



- Bosons that couple to quarks & leptons, e.g.

$$L = y_{ij} Q_i S_3 L_j + z_{ij} Q_i S_3 Q_j + \text{h.c.}$$

- Leads to **proton decay!**  
Impose  $U(1)_B$  or  $U(1)_{e,\mu,\tau}$ ?

[Barr, Freire '90; JH, Hambye, PRL 2018]

- For **B anomalies**:

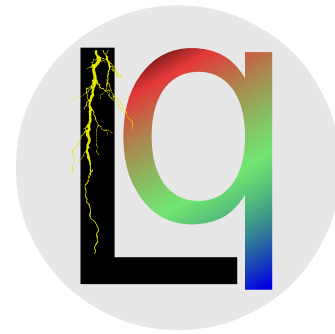
$S_1, S_3, R_2, U_1, U_3$  . [see Andrei's talk!]

- Too *ad hoc*?  
Leptoquarks part of GUTs  
as multiplet partners!

LQ	SM rep	spin	p decay
$S_1$	$(\bar{\mathbf{3}}, \mathbf{1}, +1/3)$	0	yes
$\bar{S}_1$	$(\bar{\mathbf{3}}, \mathbf{1}, -2/3)$	0	yes
$\tilde{S}_1$	$(\bar{\mathbf{3}}, \mathbf{1}, +4/3)$	0	yes
$S_3$	$(\bar{\mathbf{3}}, \mathbf{3}, +1/3)$	0	yes
$R_2$	$(\mathbf{3}, \mathbf{2}, +7/6)$	0	no
$\tilde{R}_2$	$(\mathbf{3}, \mathbf{2}, +1/6)$	0	yes
$U_1$	$(\mathbf{3}, \mathbf{1}, +2/3)$	1	no
$\bar{U}_1$	$(\mathbf{3}, \mathbf{1}, -1/3)$	1	no
$\tilde{U}_1$	$(\mathbf{3}, \mathbf{1}, +5/3)$	1	no
$U_3$	$(\mathbf{3}, \mathbf{3}, +2/3)$	1	no
$V_2$	$(\bar{\mathbf{3}}, \mathbf{2}, +5/6)$	1	yes
$\tilde{V}_2$	$(\bar{\mathbf{3}}, \mathbf{2}, -1/6)$	1	yes

[Review: Doršner+, '16]

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- For **B anomalies**:

$S_1, S_3, R_2, U_1, U_3$  . [see **Andrei's talk!**]

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Leptoquarks part of GUTs  
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LQ	SM rep	$R_K^{(*)}$	$R_D^{(*)}$
$S_1$	$(\bar{\mathbf{3}}, \mathbf{1}, +1/3)$	(✓)	✓
$\bar{S}_1$	$(\bar{\mathbf{3}}, \mathbf{1}, -2/3)$		
$\tilde{S}_1$	$(\bar{\mathbf{3}}, \mathbf{1}, +4/3)$		
$S_3$	$(\bar{\mathbf{3}}, \mathbf{3}, +1/3)$	✓	
$R_2$	$(\mathbf{3}, \mathbf{2}, +7/6)$	✓	or ✓
$\tilde{R}_2$	$(\mathbf{3}, \mathbf{2}, +1/6)$		
$U_1$	$(\mathbf{3}, \mathbf{1}, +2/3)$	✓	✓
$\bar{U}_1$	$(\mathbf{3}, \mathbf{1}, -1/3)$		
$\tilde{U}_1$	$(\mathbf{3}, \mathbf{1}, +5/3)$		
$U_3$	$(\mathbf{3}, \mathbf{3}, +2/3)$	✓	
$V_2$	$(\bar{\mathbf{3}}, \mathbf{2}, +5/6)$		
$\tilde{V}_2$	$(\bar{\mathbf{3}}, \mathbf{2}, -1/6)$		

[Angelescu+, 1808.08179]

# Pati-Salam

$$SU(4)_C \times SU(2)_L \times SU(2)_R$$

[Pati, Salam, '74]

- Fermions:  $\Psi_L \sim (4, 2, 1) \rightarrow Q_L \oplus L_L$ ,  
 $\Psi_R \sim (4, 1, 2) \rightarrow u_R \oplus d_R \oplus \ell_R \oplus N_R$ .

- $SU(4)_C \rightarrow SU(3)_C \times U(1)_{B-L}$   
generates massive  $U_1$ .

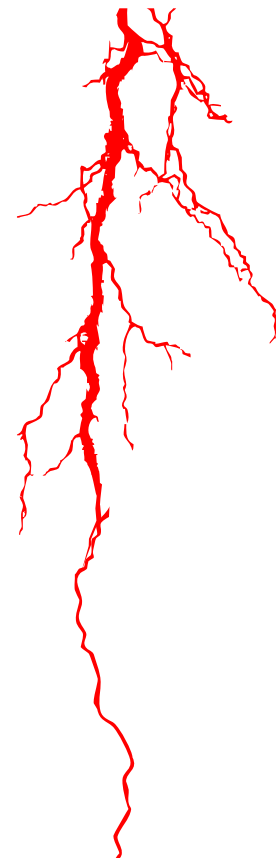
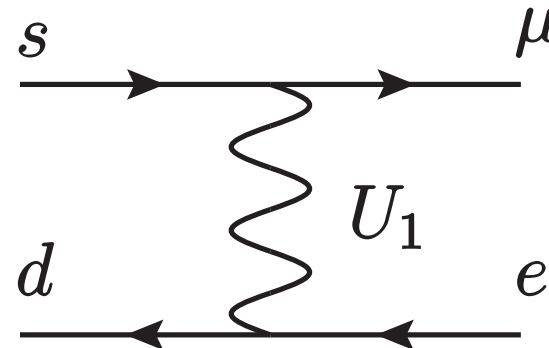
- $K_L \rightarrow \mu e$  sets bound

$$m_{U_1} \gtrsim 1000 \text{ TeV.}$$

[Valencia, Willenbrock, '94; Smirnov, '07/'18]

- Need extra work to lower  $U_1$  and get LFUV.

[Calibbi+, '17; Di Luzio+, '17; Blanke, Crivellin, '18; Bordone+, '18]



How about **scalar** leptoquarks?

# Pati-Salam $SU(4)_C \times SU(2)_L \times SU(2)_R \rightarrow SM$

- Breaking (and  $N_R$  mass!) via scalars

$$\Delta_R \sim (\overline{\mathbf{10}}, \mathbf{1}, \mathbf{3}) \supset S_1 \oplus \tilde{S}_1 \oplus \bar{S}_1 \oplus \delta_1,$$

$$\Delta_L \sim (\overline{\mathbf{10}}, \mathbf{3}, \mathbf{1}) \supset S_3 \oplus \delta_3.$$

type-I seesaw

type-II seesaw

$$\mathcal{L} = \bar{\Psi}_L^c y_L^L \Delta_L \Psi_L + \bar{\Psi}_R^c y_R^R \Delta_R \Psi_R + \text{h.c.}$$

no proton decay!

- Parity  $X_L \leftrightarrow X_R$  requires  $\Delta_L$  and sets  $y_L = y_R$ .
- Only **one (symmetric) coupling matrix!**

PS relates couplings of different LQs & to neutrinos!



# Type-II seesaw $\leftrightarrow R_{K(*)}$

$$\bullet \bar{\Psi}_L^c y^L \Delta_L \Psi_L \supset \underbrace{\left( \bar{Q}_L^c y^L L_L S_3 \right)}_{R_{K(*)}} + \frac{1}{\sqrt{2}} \underbrace{\left( \bar{L}_L^c y^L L_L \delta_3 \right)}_{\text{type-II seesaw}}$$

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$$\bullet M_\nu \simeq -\sqrt{2} \langle \delta_3 \rangle V_L^* y^L V_L^\dagger \xRightarrow{*} y^L \propto \begin{pmatrix} 0.05 & 0.06 & -0.10 \\ 0.06 & \boxed{1} & \boxed{0.74} \\ -0.10 & 0.74 & 0.97 \end{pmatrix}$$

- $m_{S_3} \sim 30 \text{ TeV}$  can give  $R(K)$ !

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- $m_{S_3} \sim 30 \text{ TeV}$  can give  $R(K)$ !
- $\mu \rightarrow e$  conversion too large, use CP phases to suppress.

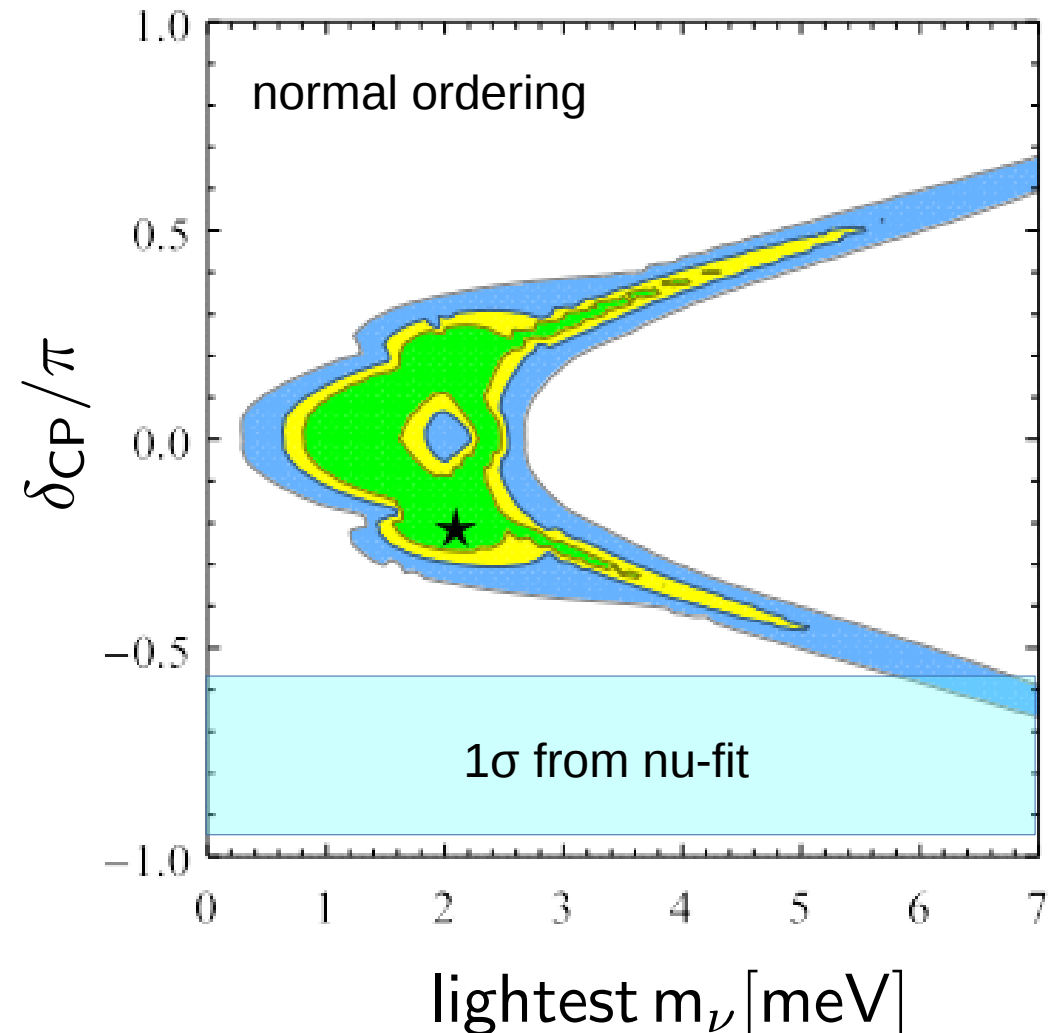


\*) normal ordering,  $m_1 = 0$ ,  $\delta_{CP} = \alpha = \beta = 0$ .

# $R_{K^{(*)}}$ & type-II seesaw

- $R(K)$  and  $\mu \rightarrow e$  fix neutrino parameters!
- Fixing  $\langle \delta_3 \rangle = 50$  meV :

parameter	best fit
$M_{S_3}$	11.4 TeV
$m_{\nu_1}$	2.1 meV
$\delta_{CP}$	$-0.2\pi$
$\alpha$	$0.55\pi$
$\beta$	$0.95\pi$
$\chi^2$	2.45

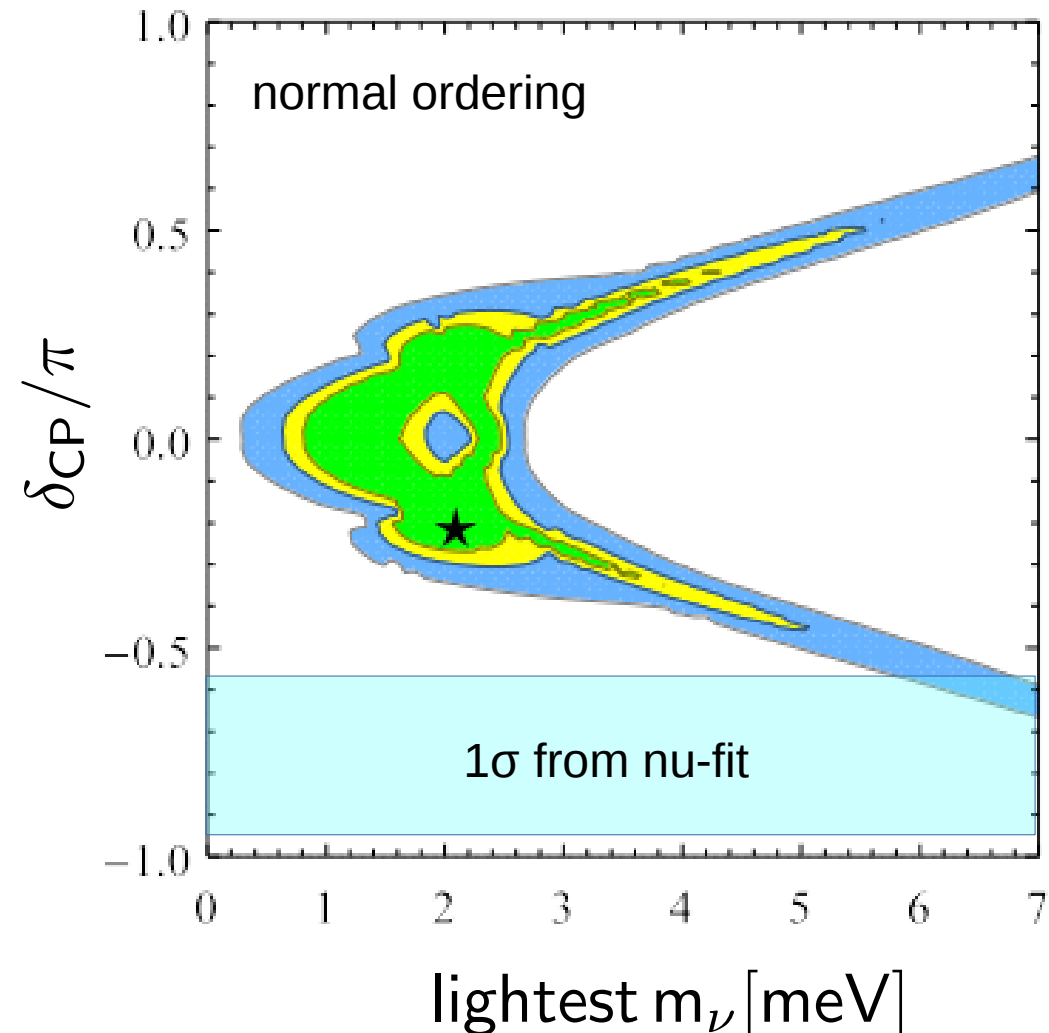


Should show up in next-gen  $\mu \rightarrow e$ !

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Non-trivial, only works for type-II seesaw with normal ordering!

# What about $R_{D^{(*)}}$ ?

- Difficult, our  $S_1$  LQs couple to  $N_R$ , not  $\nu_L$ .
- Could use the  $R_2$  LQs from the EWSB, but too flexible.
- Forget neutrino connection and assume **one light  $N_R$** .

- $S_3$  gives  $R(K)$ ,  $S_1$  gives  $R(D)$ ,  
same coupling matrix  $y^L = y^R \simeq \begin{pmatrix} 0 & 0 & 0 \\ 0 & y_{22} & y_{23} \\ 0 & y_{23} & y_{33} \end{pmatrix}$ .

- Fixing  $M_{S_1} = 1 \text{ TeV}$  :

parameter	best fit
$M_{S_3}$	6.5 TeV
$y_{22}$	0.034
$y_{23}$	1.17
$\chi^2$	2.4

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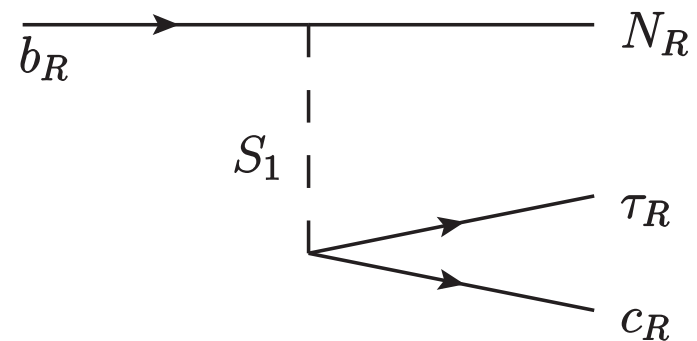
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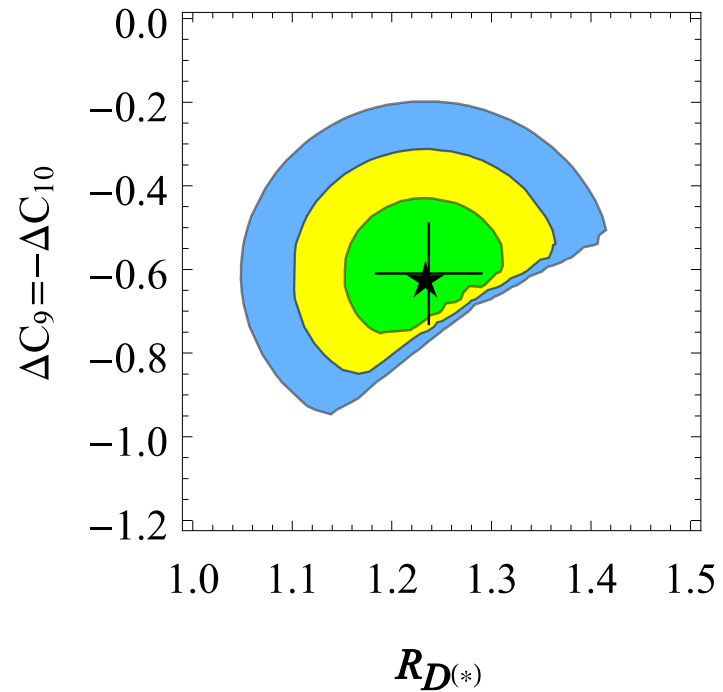
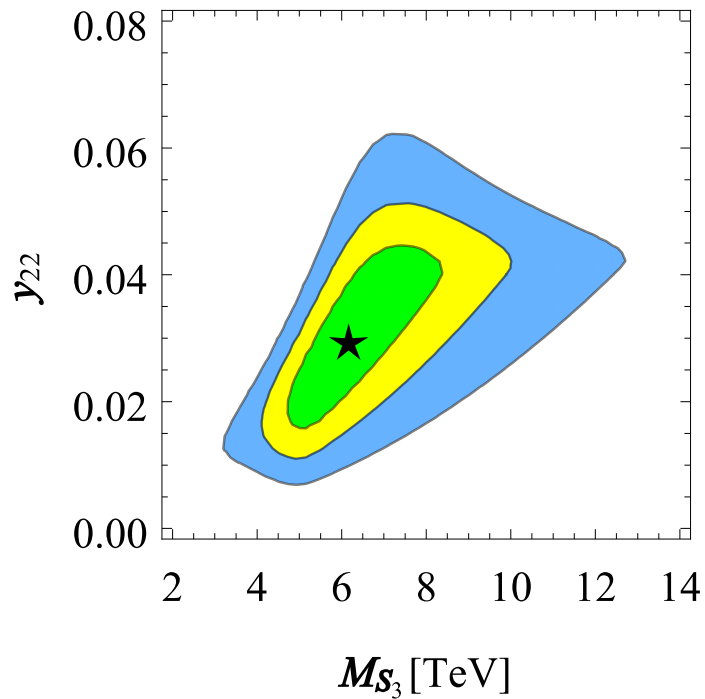
- $S_3$  gives  $R(K)$ ,  $S_1$  gives  $R(D)$ ,  
same coupling matrix

$$y^L = y^R \simeq \begin{pmatrix} 0 & 0 & 0 \\ 0 & y_{22} & y_{23} \\ 0 & y_{23} & \cancel{y_{33}} \end{pmatrix} \begin{matrix} R(K) \\ R(D) \end{matrix}$$

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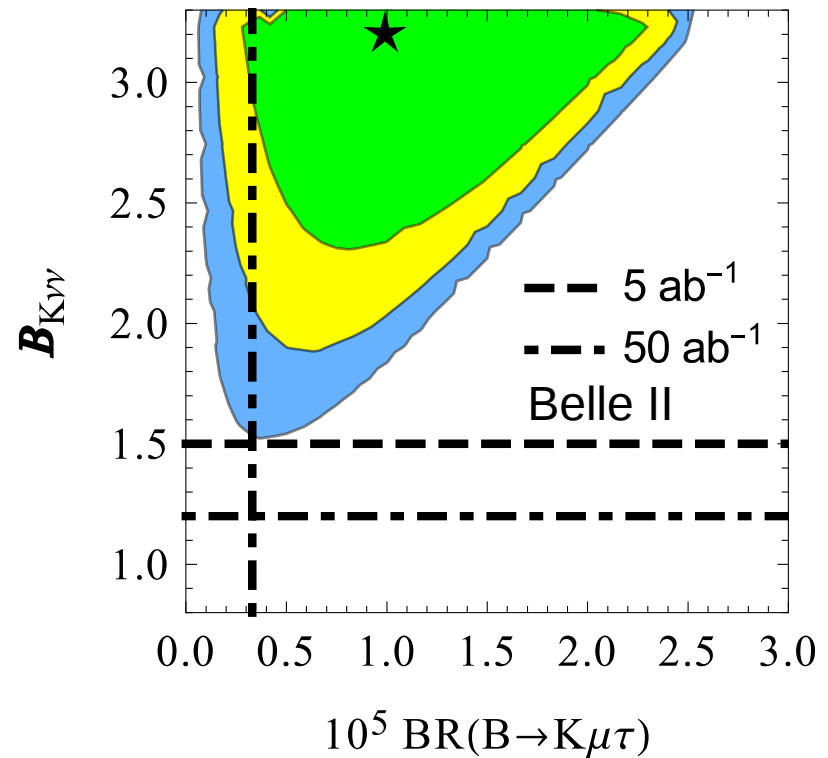
- $S_1$  testable in Belle-II:

$$B \rightarrow K\nu\nu \text{ \& \ } B \rightarrow K\mu\tau$$

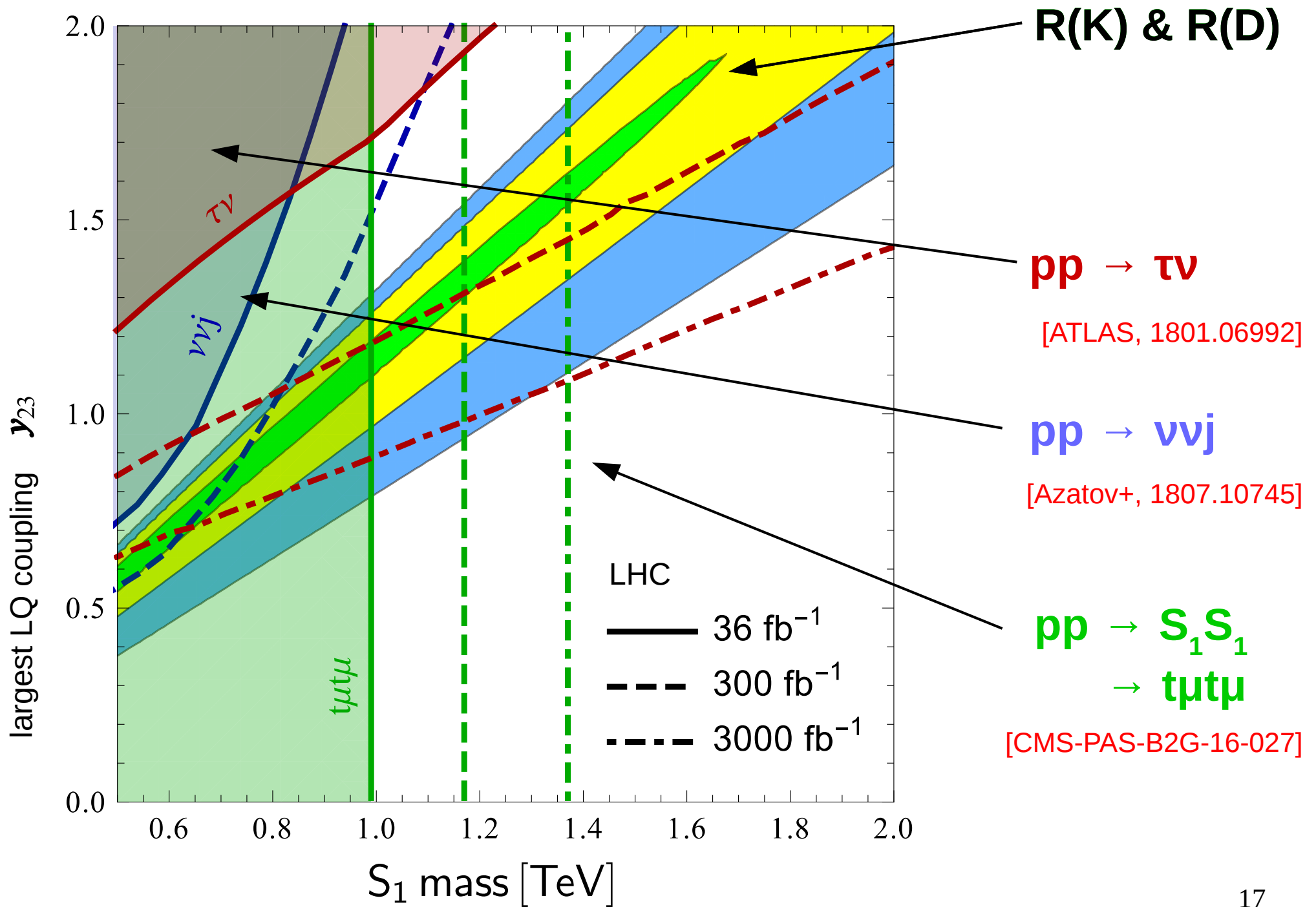
and at the LHC:

$$pp \rightarrow \tau\nu, \nu\nu j, t\mu t\mu.$$

- $\nu = N_R$  here.







# Conclusions

- Lepton non-universality in B decays very intriguing.
- New physics in the form of  $Z'$ ,  $W'$  or **leptoquarks!**
- LQ explanation fits surprisingly nicely into **Pati-Salam**:
  - Required for symmetry breaking & seesaw.
  - Automatically chiral & no proton decay.
  - Pati-Salam **relates couplings** of LQs and seesaw.
  - Parity relates  $S_3$  and  $S_1$  couplings.
- **Testable**:  $pp \rightarrow \tau\nu, \nu\nu j, t\mu t\mu$ ,  
 $BR(S_1 \rightarrow t\mu) \simeq BR(S_1 \rightarrow c\tau) \simeq BR(S_1 \rightarrow b \text{ inv})$ .



# Backup

# Fermion masses

- $(\bar{4}, 2, 1) \otimes (4, 1, 2) = (1, 2, 2) \oplus (15, 2, 2)$

- Just complex  $(1, 2, 2)$  :  $m_d = m_\ell, \quad m_u = m_N^{\text{Dirac}}$  [Volkas, '95]

- Diagonalization:

$$u_{L,R} \rightarrow V_{L,R}^\dagger u_{L,R}, \quad \nu_L \rightarrow V_L^\dagger \nu_L, \quad N_R \rightarrow V_R^\dagger N_R$$

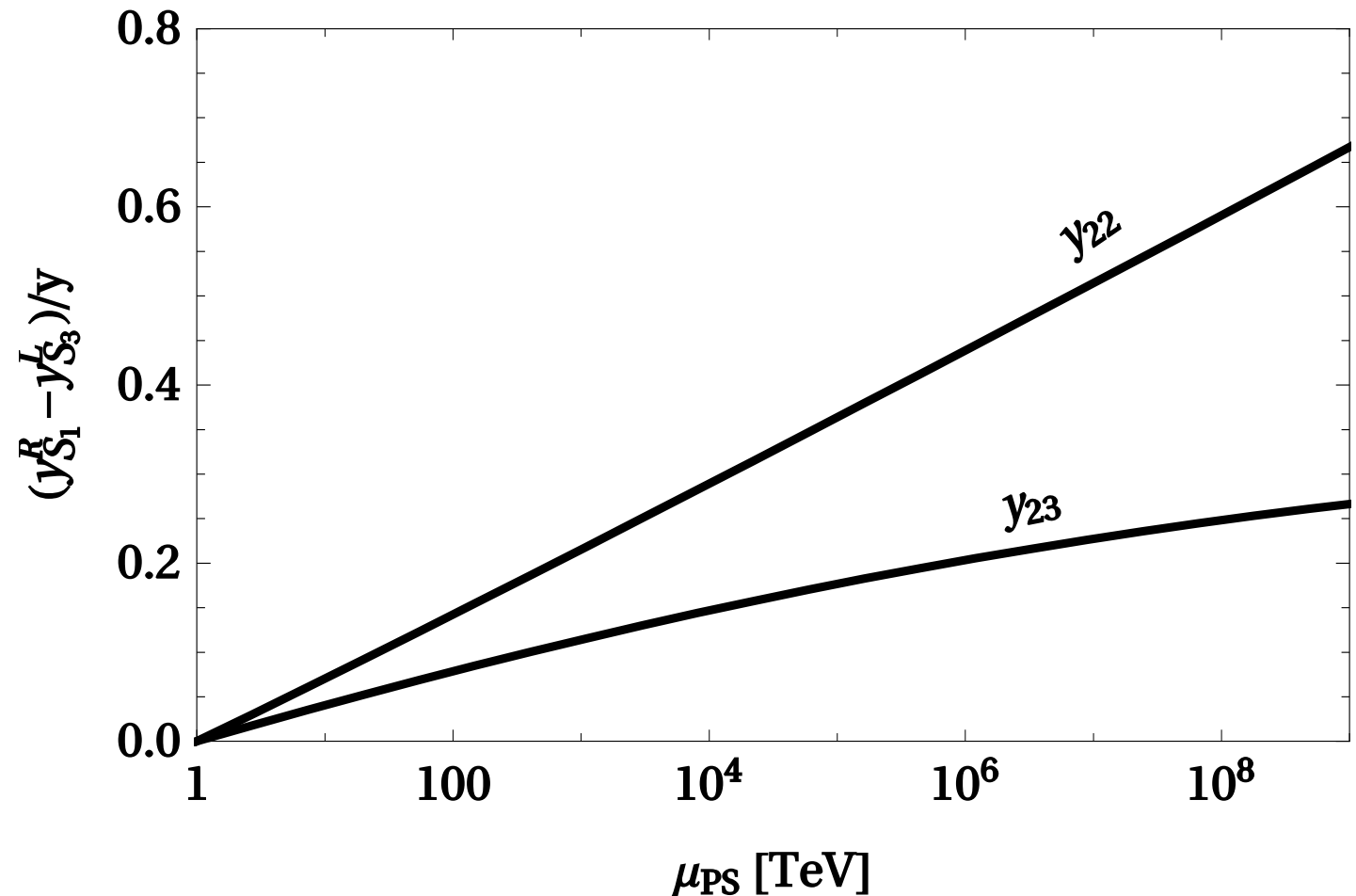
with parity relation  $V_R \simeq V_L$ . [Maiezza+, '10]

- Adding  $(15, 2, 2)$  gives freedom (4HDM) and  $R_2$  LQs.

$$(1, 2)_{\frac{1}{2}} \oplus (1, 2)_{-\frac{1}{2}} \oplus (3, 2)_{\frac{1}{6}} \oplus (\bar{3}, 2)_{-\frac{1}{6}} \oplus (3, 2)_{\frac{7}{6}} \oplus (\bar{3}, 2)_{-\frac{7}{6}} \oplus (8, 2)_{\frac{1}{2}} \oplus (8, 2)_{-\frac{1}{2}}$$

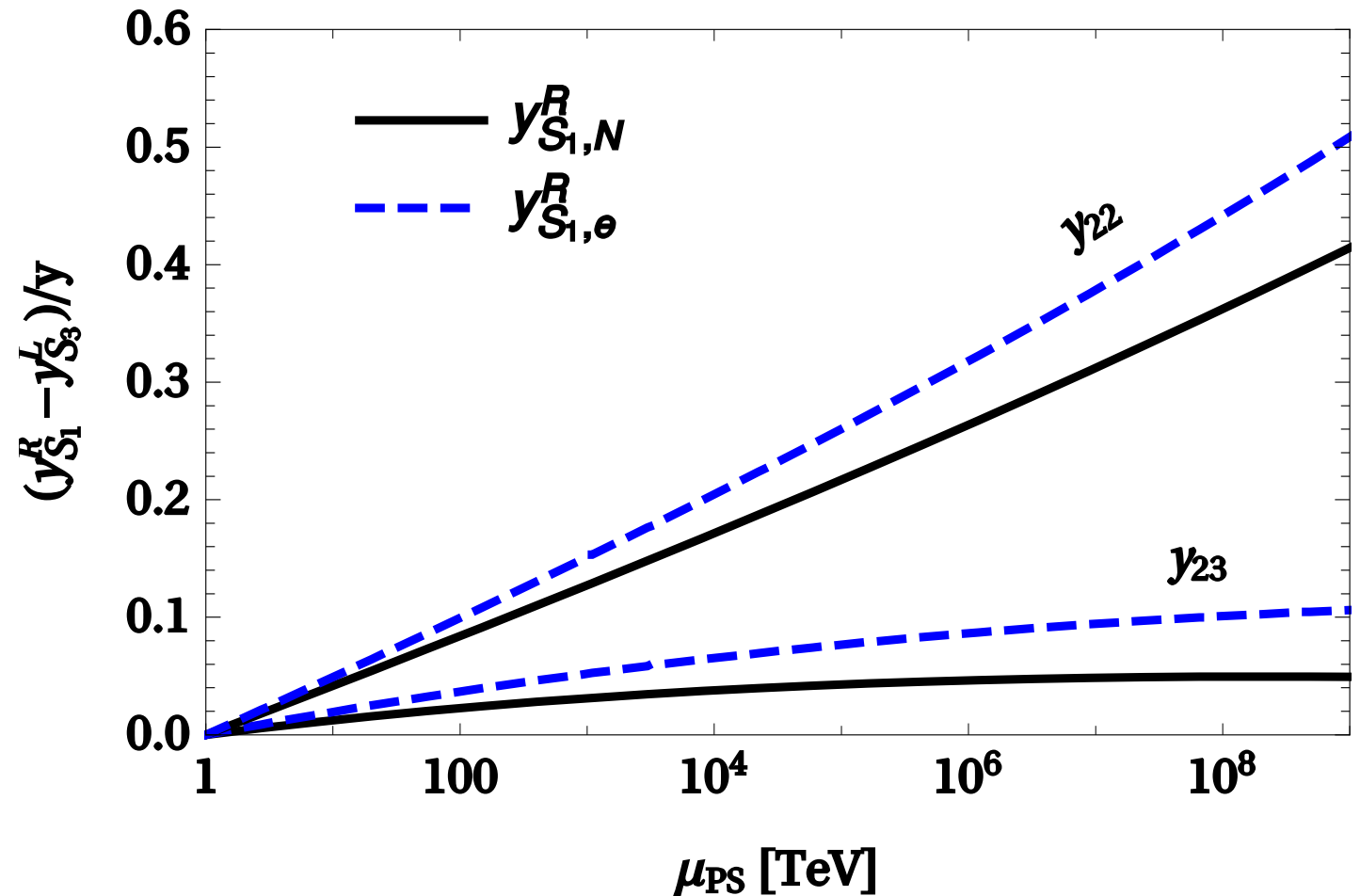
# RGEs

- The Pati-Salam relations  $y = y^T$  and  $y_L = y_R$  are broken!
- RGEs depend strongly on all other particle masses.
- Heavy  $N_R$ :



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# Unification

- Parity requires  $g_L = g_R$  at the PS scale. Possible?
- RGEs depend strongly on all other particle masses.
- Lowest order:  $\frac{3}{5} \alpha_L^{-1}(\mu_{PS}) = \alpha_1^{-1}(\mu_{PS}) - \frac{2}{5} \alpha_C^{-1}(\mu_{PS})$

$$\Rightarrow \mu_{PS} \sim 5 \times 10^{13} \text{ GeV.}$$

- More light states change this.

