

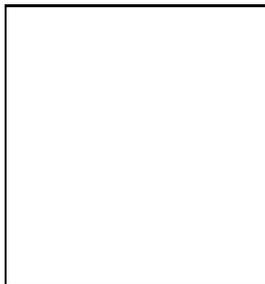
# ITERATIVE MAP-MAKING FOR SCANNING EXPERIMENTS

S. Prunet<sup>1</sup>, C. B. Netterfield<sup>2</sup>, E. Hivon<sup>3</sup>, B. P. Crill<sup>3</sup>

<sup>1</sup> *CITA, Mc Lennan Labs, University of Toronto*

<sup>2</sup> *Dept. of Astronomy, Mc Lennan Labs, University of Toronto*

<sup>3</sup> *Caltech Observational Cosmology, Pasadena*



We describe here an iterative method for jointly estimating the noise power spectrum from a scanning experiment's time-ordered data, together with the maximum-likelihood map. We test the robustness of this method on simulated datasets with colored noise, like those of bolometer receivers in CMB experiments.

## 1 Introduction

The map-making problem for CMB anisotropy measurements was first considered in the context of the COBE-DMR mission<sup>2,4</sup>. It was further extended to quick algorithms for differential measurements<sup>5</sup> taking into account  $1/f$  noise<sup>6</sup>. Given the size of the upcoming datasets ( $\geq 10^6$  pixels for *MAP*) it is essential that the map-making algorithm remains quick (typically  $O(n_d \log L_N)$  where  $n_d$  is the number of time-samples, and  $L_N$  the effective noise-filter length). Another related question, which has not been addressed so far, is how to determine the noise statistical properties from the data itself. We present a map-making method based on Wright's fast algorithm to compute iteratively a minimum variance map together with an estimate of the detector noise power spectrum which is needed to compute unbiased power spectrum estimators of the cosmological signal.

## 2 Iterative mapmaking - Application to simulations

### 2.1 Method

Following Tegmark<sup>3</sup>, we model the data stream in the following way:

$$\mathbf{d}_t = \mathbf{P}_{\text{tp}} \Delta_{\mathbf{p}} + \mathbf{n}_t \tag{1}$$

where  $\Delta_p$  is a pixelized version of the *observed* sky (i.e. convolved by the experimental beam), and  $n_t$  is the detector noise after primary deconvolution of any filter present in the instrumental chain (e.g. bolometer time constant, read-out filters). We will assume here that the experiment is a total power measurement, i.e. that the pointing matrix  $P_{tp}$  contains only one non-zero element per row. We now want to estimate the minimum variance map from this data, ie the map that minimizes  $\chi^2 = (d - P\Delta)^\dagger N^{-1} (d - P\Delta)$ . The solution is given by:

$$\tilde{\Delta} = \left( P^\dagger N^{-1} P \right)^{-1} P^\dagger N^{-1} d \quad (2)$$

A few remarks are necessary at this point. First, the matrix to be inverted is huge ( $n_{pix} \times n_{pix}$ ) so that an iterative linear solver is needed. Secondly, the noise correlation matrix  $N_{tt}^{-1}$  has to be determined from the data itself. To make this tractable we assume that, at least over subsets of the time-stream, the noise is reasonably stationary, so that the multiplication by  $N^{-1}$  becomes a convolution operator, in other words that it is diagonal in Fourier space<sup>a</sup>. We thus implemented the following algorithm:

**for each stationary noise subset**

- $n^{(j)} = d - P\tilde{\Delta}^{(j)} \Rightarrow N^{(j)-1} = \langle nn^\dagger \rangle^{-1}$
- $\tilde{\Delta}^{(j+1)} - \tilde{\Delta}^{(j)} = \left( P^\dagger W^* P \right)^{-1} P N^{(j)-1} n^{(j)}$

**endfor**

To keep the algorithm as fast as possible, we took  $W^*$  to be diagonal and constant, so that  $P^\dagger W^* P$  is diagonal, with each element being equal to the number of observations per pixel, up to a multiplicative constant. The choice of the numerical value of this constant is only important for the convergence properties of the algorithm. In this form, the algorithm is very similar to a Jacobi iterative solver (in which we would have  $\left( \text{diag}\{P^\dagger N^{-1} P\} \right)^{-1}$  instead of  $\left( P^\dagger W^* P \right)^{-1}$ ). It has the additional advantage of being very easy to compute, for very similar convergence properties. Since we made the assumption that each noise matrix is diagonal in Fourier space, all time-domain operations are done in Fourier space using FFTs, thus reducing the number of floating point operations to  $O(L_N \log L_N)$  for each subset.

The advantages of this iterative method are obvious: it is fast ( $O(n_d \log L_N)$  operations) and cheap in memory ( $O(L_N)$  storage). However, since we try to estimate both the noise power spectrum and the map in a leap-frog manner, the convergence properties must be studied by means of numerical simulations.

## 2.2 Simulations

To test the convergence properties of the algorithm, we simulated a scanning experiment on a fake  $\Lambda$ CDM map with a gaussian beam of 10.5 arcmin FWHM, and noise whose power spectrum is  $P(f) \propto (1 + f_0/f)^2$ . To simulate the contamination of low frequencies by possible scan-synchronous systematics, we high-pass filtered the generated time-stream with a step filter. This has the effect of formally replacing the pointing matrix  $P$  by  $FP$  where  $F$  is the step filter; however, looking at Eq. 2, this is completely equivalent to leaving  $P$  unchanged, and replacing  $N^{-1}$  by  $F^\dagger N^{-1} F$ , which is straightforward in Fourier space.

<sup>a</sup>This is actually only approximately true since a convolution operator is a circulant matrix, ie it assumes that the time-stream has periodic boundary conditions; this is however a rather good approximation for a time-stream much longer than the effective length of the noise-filter<sup>3</sup>

To be more specific, we used a spherical cap of  $\sim 4 \times 10^4$  HEALPix pixels of  $7'$  size<sup>1</sup> generated from an  $\Lambda$ CDM COBE-normalized power spectrum, to generate a  $\sim 6 \times 10^6$  sample time stream. The sampling rate used is 60 Hz, and the scan velocity 1 deg/s. The sky coverage of the simulation is shown in Fig. 1. The same figure shows the input map used for the simulation, as well as the error map from the naive (coadded) map-making and the iterative map-making after 10 iterations. One can see from those that the striping is much reduced in the iterative case, the only remaining feature in the error map being a large-scale mode that was lost in the high-pass filtering of the time-stream. Fig. 2 shows the (inverse) noise power spectrum estimate compared to the data power spectrum and the true power spectrum. We can see that the estimated PS is very close to the input one.

As in any iterative linear solver, the convergence of the solution, decomposed on the (noise-) matrix eigenmodes, is a function of the associated eigenvalue. Thus the noisiest pixel modes (usually at large scales since the time-stream is high-pass filtered, either by hand or as part of the algorithm if the noise is higher at low frequency) take (exponentially) more time to converge. The problem scales as the noise matrix condition number, which is a direct function of the noise power spectrum dynamical range and of the scanning strategy. We propose to accelerate the convergence of the algorithm with the implementation of a multi-grid like method where the time-stream, as well as the pixels, get rebinned. This method will be described in a future paper.

### 3 Conclusions - Perspectives

We described a fast iterative map-making method to simultaneously generate the maximum-likelihood map and the noise power spectrum from a scanning experiment time stream. We tested its convergence properties on simulations, and concluded that, except for the spatially largest (and ill conditioned) modes, the map and noise power spectra converge very quickly. We propose to cure the convergence of this large scale modes by applying a multi-grid method to be described in a future publication.

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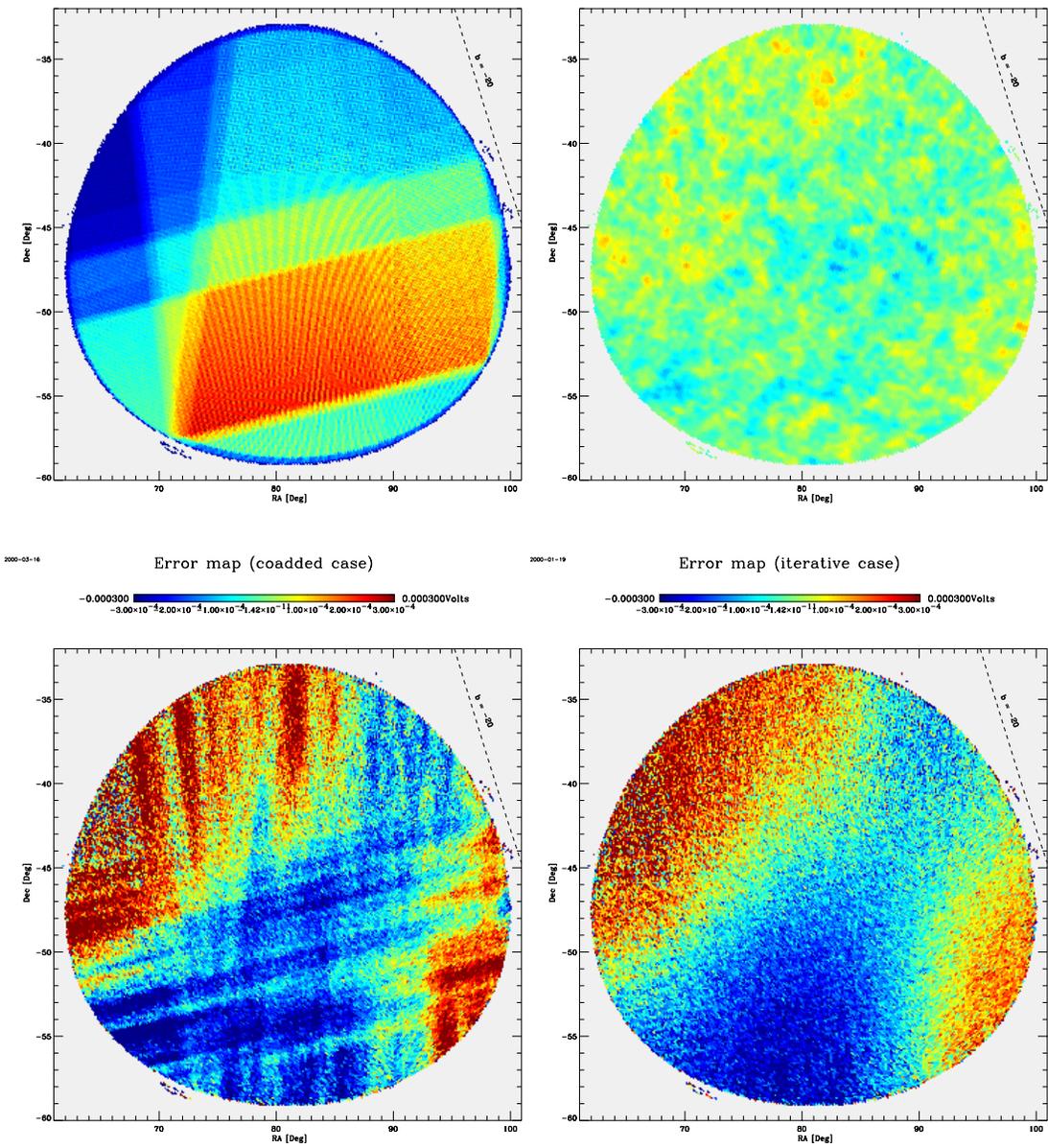


Figure 1: Simulations- Upper left: Coverage - Upper right: input map - Lower left: error map (coadded case) - Lower right: error map (iterative case).

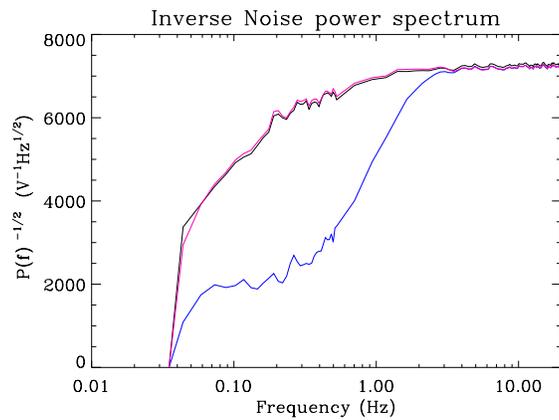


Figure 2: Inverse power spectra of noise in the time-stream. Black is the input noise, blue is the data, and red is the estimation.