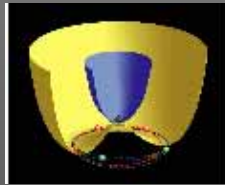
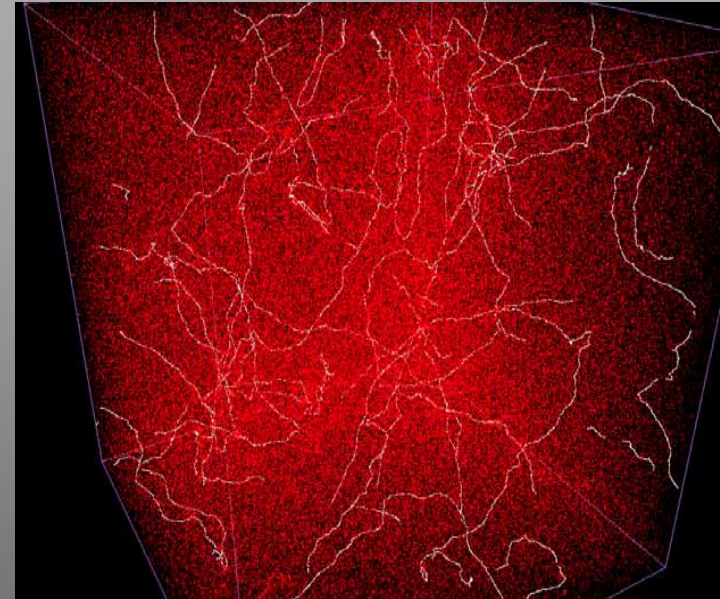
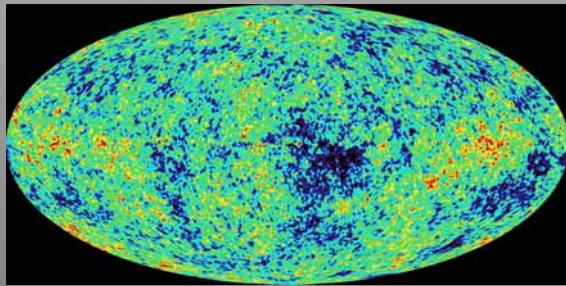


Cosmology can constrain SUSY models: The truth about the role of cosmic strings



In collaboration with Jonathan Rocher

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Within SUSY GUTs, cosmic string formation is unavoidable in all acceptable SSB schemes from the GUT down to the standard model gauge group.

Jeannerot, Rocher, Sakellariadou PRD **68**, 103514 (2003)

See talk by Jonathan Rocher

Mixed models: hybrid inflation + strings

Cosmic strings and/or embedded strings

But ...

Embedded defects:

Consider a pattern of symmetry breaking

$$G \supset H$$

To construct embedded defects, choose a sub-group

$$G_{\text{emb}} \subset G$$

such that $\dim(G_{\text{emb}}) = \dim(H_{\text{emb}})$ is non-trivial,

where
$$H_{\text{emb}} = H \setminus G_{\text{emb}}$$

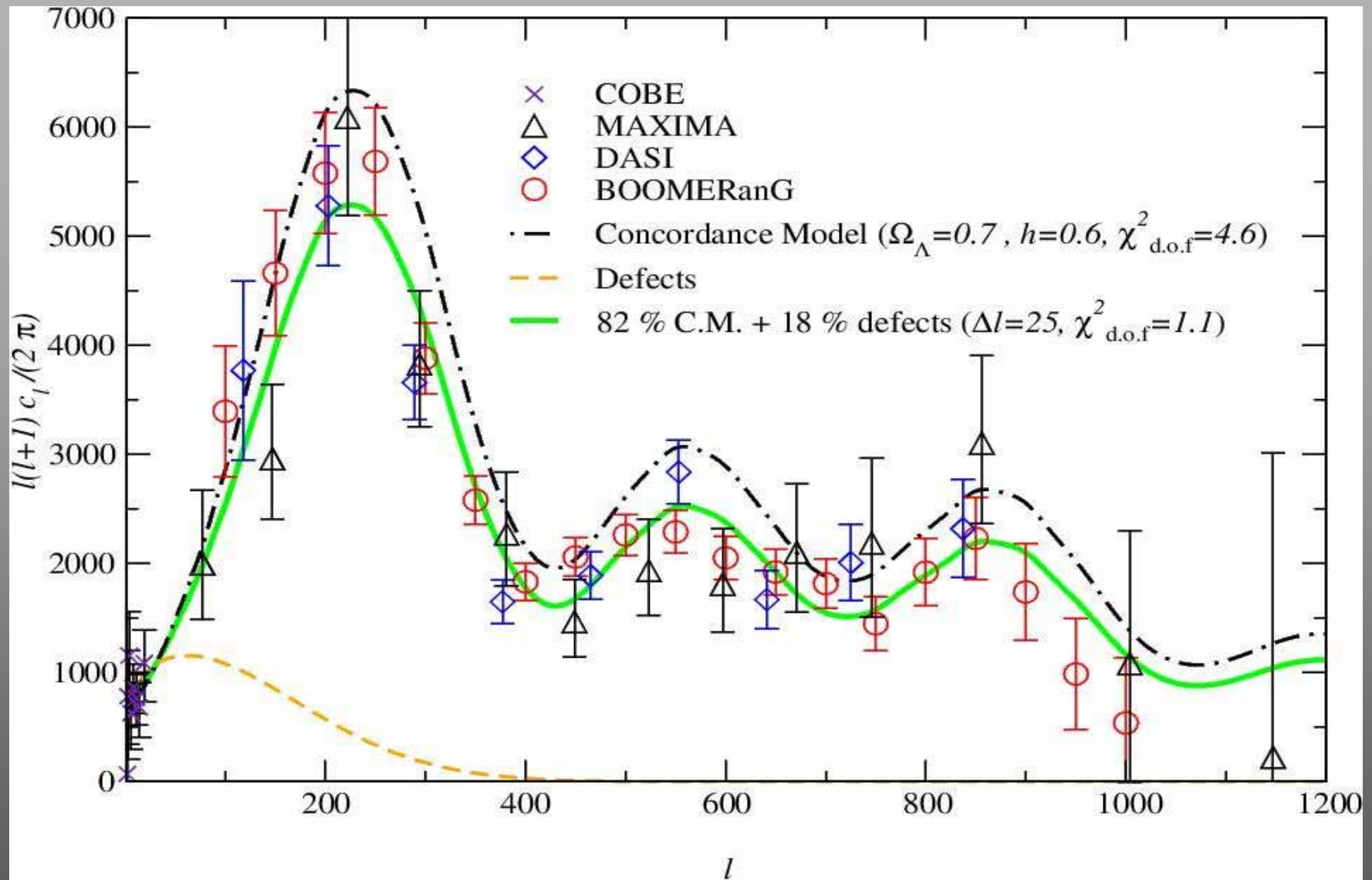
The cosmic string contribution to the CMB temperature anisotropies is highly constrained from the measurements !

$$C_\ell = \epsilon C_\ell^{\text{infl}} + (1 - \epsilon) C_\ell^{\text{CS}}$$

Already from the measurements before WMAP, the cosmic string contribution was at the most 18%.

Bouchet, Riazuelo, Peter & Sakellariadou PRD 65, 021301 (2001)

WMAP should impose even strongest constraints.



Bouchet, Riazuelo, Peter & Sakellariadou PRD 65, 021301 (2001)

Bayesian analysis in a 3dim space:

$$C_{\ell} = \alpha C_{\ell}^{\text{infl}} + \beta C_{\ell}^{\text{cs}}$$

parameters α, β, n_s with ranges:

$$\alpha \in [0.95; 1.05]; \beta \in [0.0; 0.105]; n_s \in [0.955; 1.02]$$

WMAP and SDSS data set an upper bound on the fraction of primordial fluctuations sourced by local cosmic strings, which at 99% confidence is:

$$\beta \hat{=} 0.09$$

Pogosian, Wyman, Wasserman, astro-ph/0403268

2 basic questions :

Which is the energy scale of the formed CS ?

Which is the theoretical predicted CS contribution to the $\delta T/T$?

How does this depend on the parameters of the model...

F-term inflation

Superpotential:

$$W = \hat{S}(\mathcal{D}_+ \mathcal{D}_a \hat{a} M^2)$$

Scalar potential with radiative corrections:

$$V = \hat{M}^4 \left(1 + \frac{\hat{N}^2}{32\hat{u}^2} 2 \ln \frac{jSj^2 \hat{a}^2}{\hat{E}^2} + g(z) \right)$$

$$g(z) = (z + 1)^2 \ln(1 + z^{\hat{a}^{-1}}) + (z - 1)^2 \ln(1 - z^{\hat{a}^{-1}})$$

$$z = x^2 = jSj^2 = M^2$$

Number of e-folds:

$$N_Q \approx N(S_Q) = \frac{32\dot{\alpha}^3 M^2}{\hat{\sigma}^2 N_{PI} M_{PI}^2} y_Q^2$$

with

$$y_Q^2 = \frac{R_{x_Q}^2}{1} \frac{dz}{zf(z)}$$

M_{PI} : Planck mass

$$x_Q = jS_Q j^2 = M^2$$

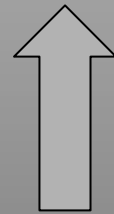
$$f(z) = (z+1) \ln(1+z^{\hat{\alpha}1}) + (z-1) \ln(1-z^{\hat{\alpha}1})$$

The relation between M and κ :

$$\frac{M}{M_{PI}} = \sqrt[3]{\frac{N_Q N_{PI} \hat{\sigma}}{32\dot{\alpha}^3 y_Q^2}}$$

Quadrupole temperature anisotropy:

$$\bar{a} \left(\frac{\hat{I}}{T} \right) Q_{\text{à tot}}^{a_2} = \bar{a} \left(\frac{\hat{I}}{T} \right) Q_{\text{à inf}}^{a_2} + \bar{a} \left(\frac{\hat{I}}{T} \right) Q_{\text{à cs}}^{a_2}$$



$$\bar{a} \left(\frac{\hat{I}}{T} \right) Q_{\text{à S}}^{a_2} + \bar{a} \left(\frac{\hat{I}}{T} \right) Q_{\text{à T}}^{a_2}$$

scalar

tensor

Using:

$$a \left(\frac{\hat{I}_T}{T}\right)_{Q\hat{a}} s^2 = \frac{32}{45} \frac{V^{3=2}}{M_{Pl}^3 V^0} \propto \frac{8}{45} \frac{N_Q M^2}{N M_{Pl}^2} x_Q^1 y_Q^1 f^1(x_Q^2)$$

inflation

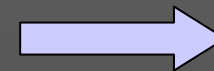
$$a \left(\frac{\hat{I}_T}{T}\right)_{Q\hat{a}} T^2 \propto 0.77 \frac{V^{1=2}}{M_{Pl}^2} \propto 0.77 \hat{M}_{Pl}^2$$

$$a \left(\frac{\hat{I}_T}{T}\right)_{Q\hat{a}} c s^2 \propto (9 \hat{a} 10) G\ddot{\omega} \quad \text{with} \quad \ddot{\omega} = 2\hat{r}' i^2$$

cosmic strings

$$\propto 18 \frac{M^2}{M_{Pl}^2}$$

and normalizing the total quadrupole temperature anisotropy to the COBE measurements, we obtain:



$$(6.3 \hat{a} \cdot 10^6)^2 = \frac{\hat{a}^2 N N_Q}{32 \hat{u}^2} y_Q^4$$

$$\hat{a} = \frac{64 N_Q}{45 N} x_Q^2 y_Q^2 f^2 (x_Q^2) + \frac{0.77 \hat{a}^2}{\hat{u}} + 324 \hat{a}$$

For given values of \hat{a} ; N ; N_Q we solve numerically the above equation and we get x_Q and thus y_Q ; M

$$N_Q = 60$$

$$N = 1; 3$$

Energy scale of inflation M as a function of the parameter κ .

The strings formed at the end of inflation have a mass proportional to the inflationary scale.

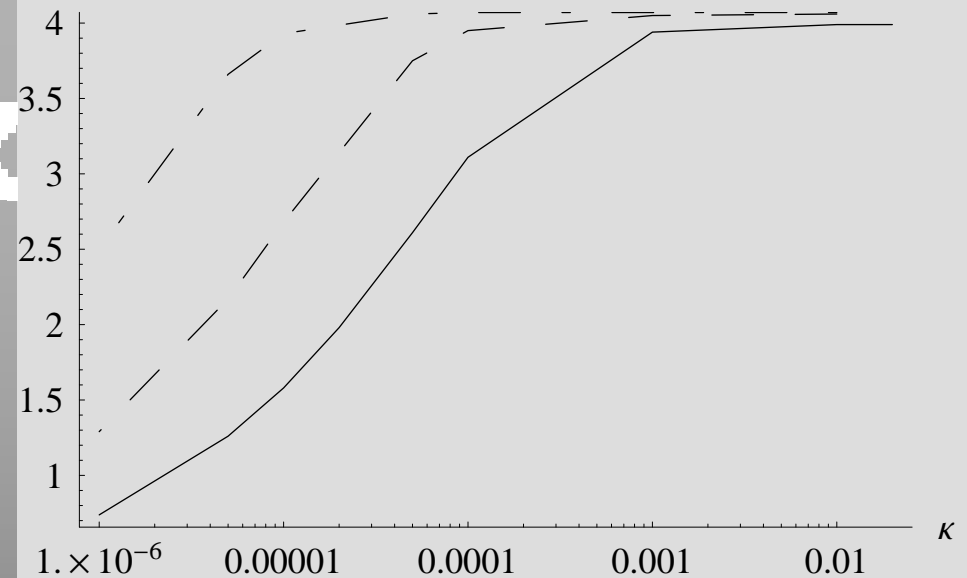
$$M \approx O(10^{15}) \text{ GeV} \ll M_{\text{Pl}}$$

→ ~~SUGRA~~

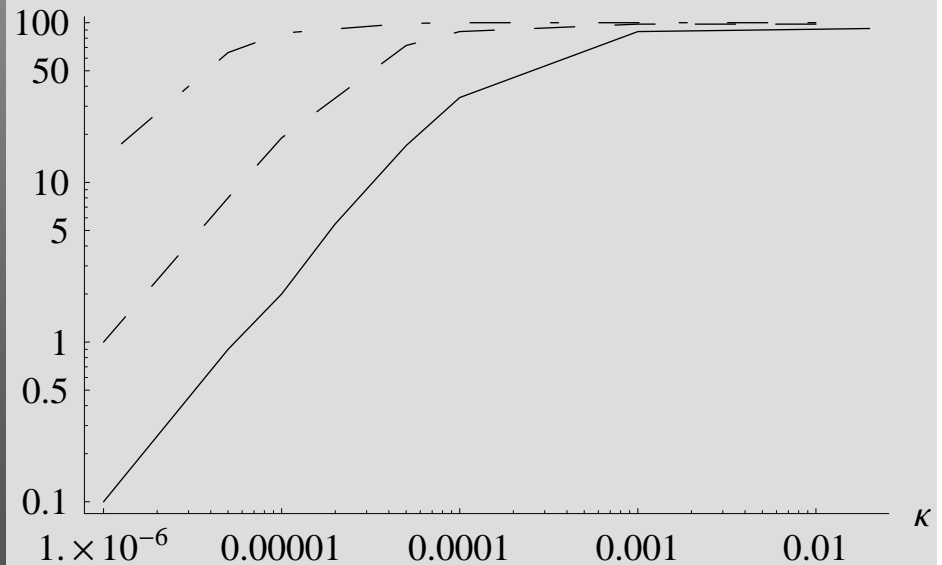
$$N = 3; 16; 126$$

Cosmic string contribution to the CMB angular power spectrum as a function of κ .

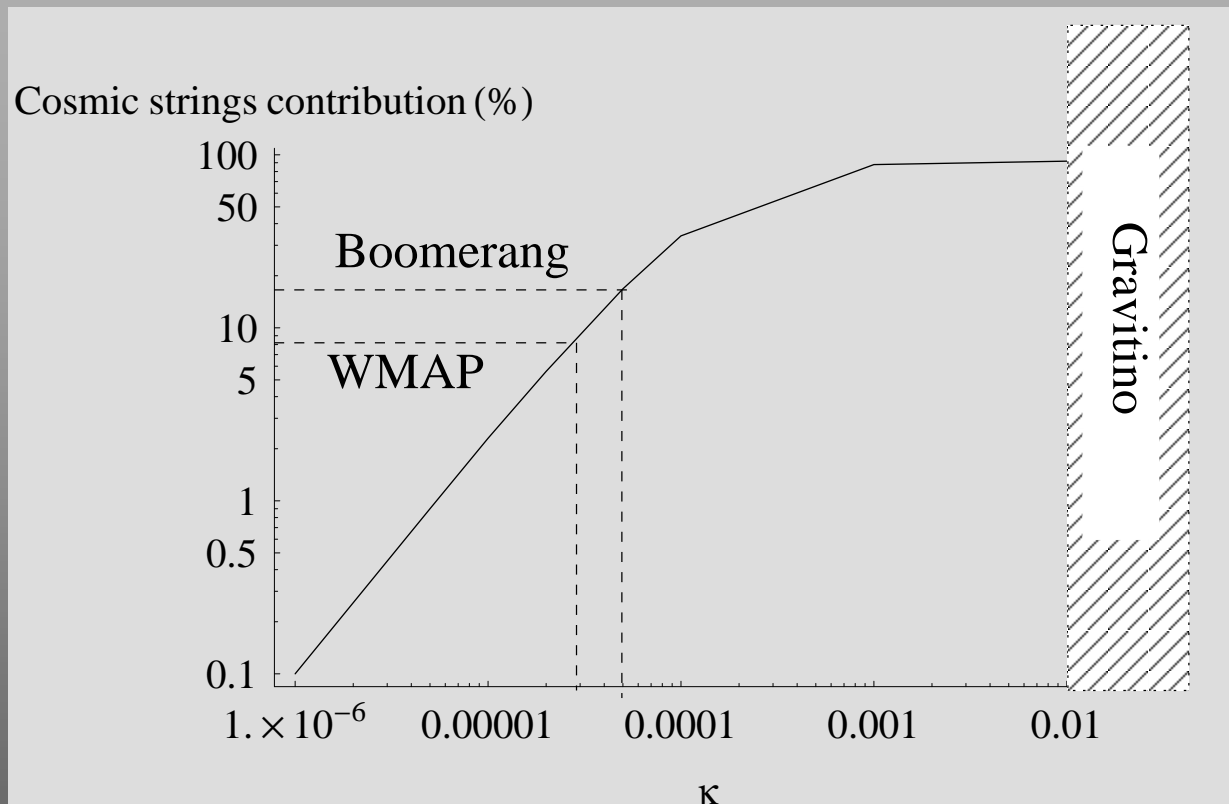
$M \times 10^{15}$ (GeV)



Contribution des cordes (%)



Constraints on the unique parameter of the model, κ , from cosmological data (i.e., CMB & gravitino)



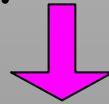
$$\hat{\kappa} \hat{\kappa} (0.8 \text{ à } 1) \hat{\kappa} 10^{\text{à } 2}$$

$$\hat{\kappa} \hat{\kappa} 3 \hat{\kappa} 10^{\text{à } 5}$$

M. Sakellariadou

The curvaton model

Embedded strings are topologically unstable and in general not dynamically stable either.



Embedded strings can play the role of the curvaton!

We can increase the allowed value of the parameter κ and still be consistent with the constraints imposed by the CMB measurements.

$$\begin{aligned}
 \left(\frac{H}{T}\right)^2 a_{\text{curv}}^2 &= a^2 \left(\frac{4 \hat{\nu}_{\text{init}}}{27 \nu_{\text{init}}}\right)^2 a_2^2 \\
 &= \frac{a_0^2 N_{\text{Q}}^2}{32 \hat{\nu}^2} y_{\text{Q}}^4 \frac{32^2}{54 \hat{\nu}^3} \hat{\nu}^2 \frac{M_{\text{Pl}}^2}{\nu_{\text{init}}} a_2^2
 \end{aligned}$$

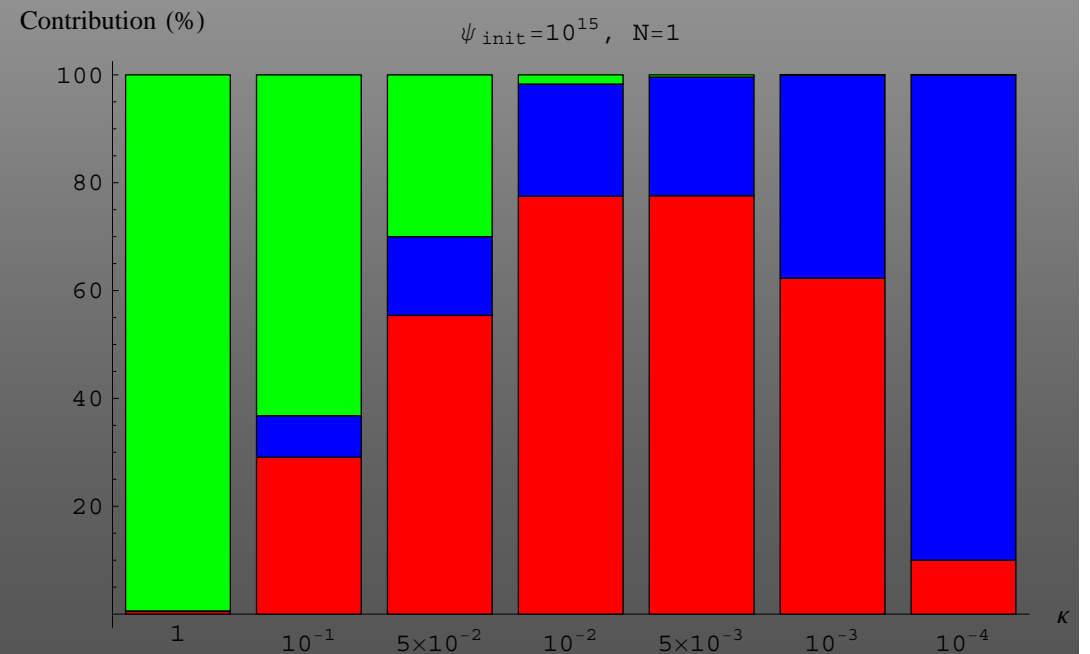
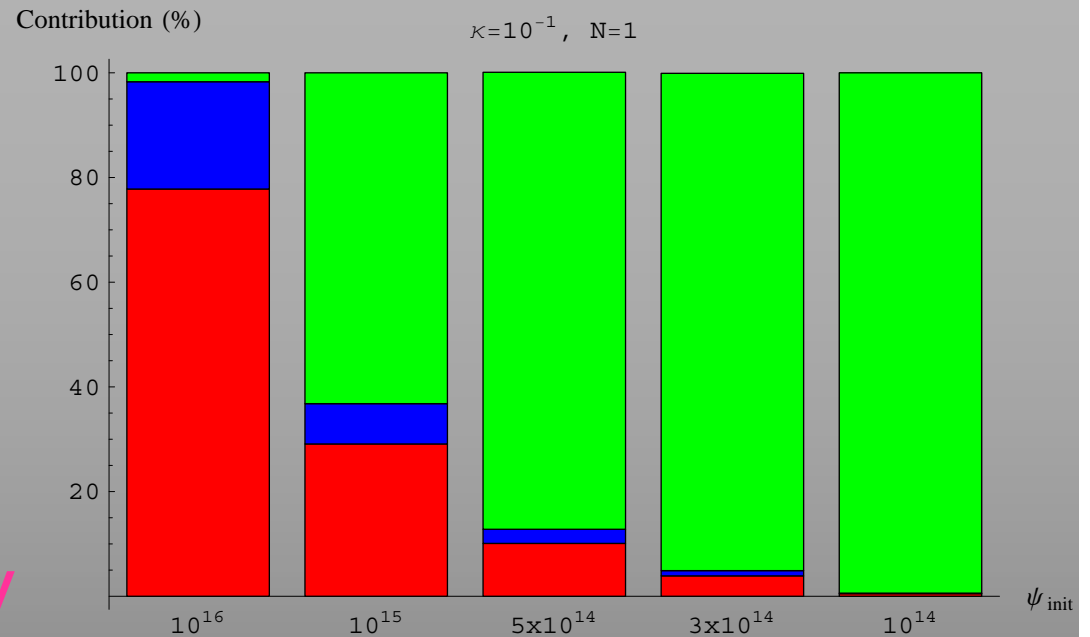
Dependence of the various CMB contributions upon the initial value of the curvaton field and the parameter κ .

$$\psi_{\text{init}} \hat{=} 3 \hat{\times} 10^{14} \frac{\hat{m}}{10^{\hat{a}+1}} \text{ GeV}$$

Topological cosmic strings

Inflaton field

Embedded strings (curvaton)



D-term inflation

Requirement: the presence of an additional U(1) factor with a nonvanishing Fayet-Iliopoulos term

$$W = \tilde{\alpha} \mathcal{D}_+ \mathcal{D}_a$$

The VEV of $\mathcal{D}_+; \mathcal{D}_a$ equals the Fayet-Iliopoulos term which also sets the scale of inflation and the energy scale of strings.

This case is much more complicated ...

The results in the literature ... are WRONG ! ...

$$V = \frac{g^2 \tilde{\phi}^2}{2} \left(1 + \frac{g^2}{16\tilde{u}^2} 2 \ln \frac{j S j^2 \tilde{\phi}^2}{\tilde{E}^2} + g(z) \right)^{\tilde{d}}$$

$$z = x^2 = \frac{\tilde{\phi}^2 j S j^2}{g^2 \tilde{\phi}}$$

$$\underbrace{\frac{\tilde{d}_1 \tilde{d}_Q \tilde{d}_2}{T}}_{6:3 \hat{a} 10^{\hat{a}6}} \tilde{\phi}^{\tilde{d}} y_Q^{\tilde{d}4} \frac{\tilde{\phi}^2 N_Q}{16\tilde{u}^2} a_2 e^{\frac{16N_Q}{45} x_Q^{\tilde{d}2}} y_Q^{\tilde{d}2} f^{\tilde{d}2}(x_Q^2) + \frac{\tilde{d}_Q \cdot 77g}{p \cdot 2\tilde{u}} \tilde{d}_2 + 324 \tilde{e} \tilde{d}$$

To reach the end of inflation

$$\frac{\tilde{\phi}}{\tilde{\phi}^2} \tilde{\phi} 10^{36}$$



$$S_Q \tilde{\phi} M_{Pl} \longrightarrow$$

SUGRA corrections are needed !

Conclusions

String formation is unavoidable at the end of inflation within SUSY GUTS.

Cosmic string contribution to the CMB data is strongly constrained.

The predicted string contribution depends on the parameters of the inflationary model (F-term / D-term).



We constraint the parameter space.