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# *Small-x Physics* and the **detection** of *Ultra-High Energy Neutrinos*

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QCD and High Energy Interactions  
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# *Outline*

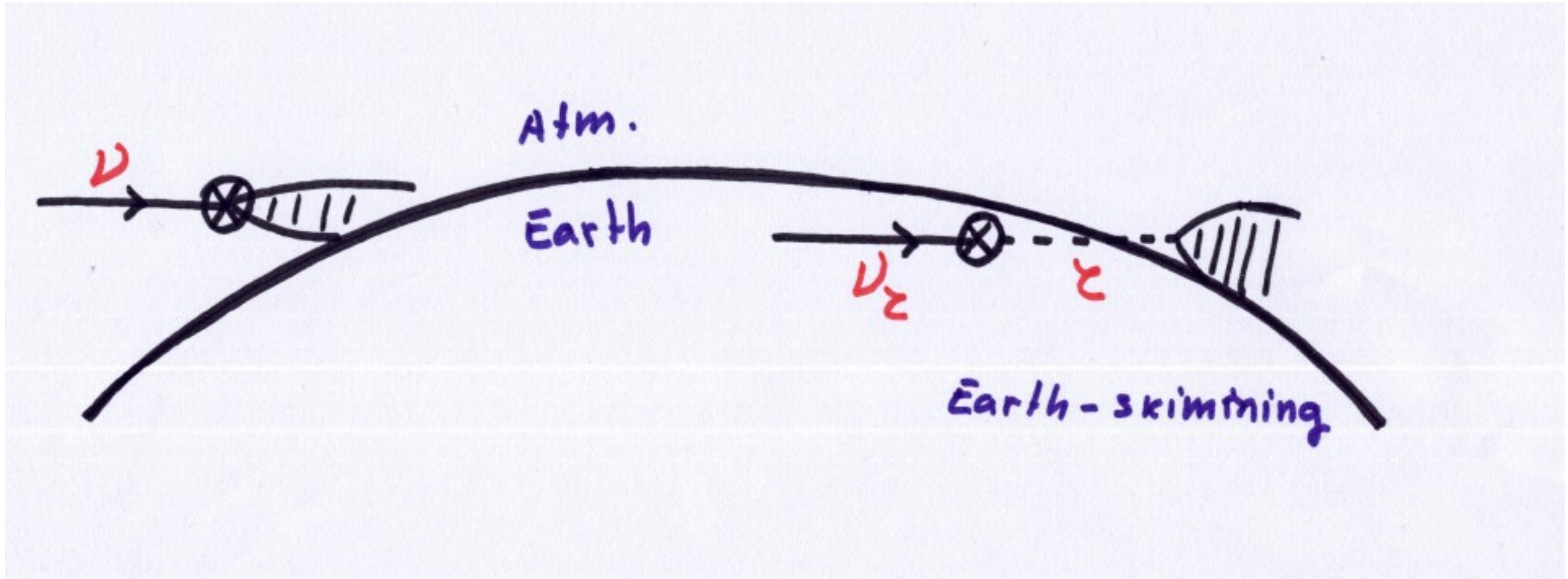
- The search for UHE cosmic neutrinos:  $\nu_\tau$
- $\tau$  energy loss at UHE
- $\nu N$  DIS CC x-section

# Detection of UHE neutrinos

- Neutrino-induced muons or showers in ice/water  
*IceCube, Antares, Nestor, Baikal*
- Neutrino-induced showers in air shower arrays  
*Pierre Auger*
  - Down-going neutrinos (inclined showers)
  - **Up-going tau-neutrinos (Earth-skimming)**

# Earth-skimming neutrinos


- $\nu_\tau$  conversion into  $\tau$  through **D.I.S.**
- $\tau$  energy loss (mostly **photonuclear**)
- $\tau$  emerges from the ground and decays in fly producing an extensive air shower



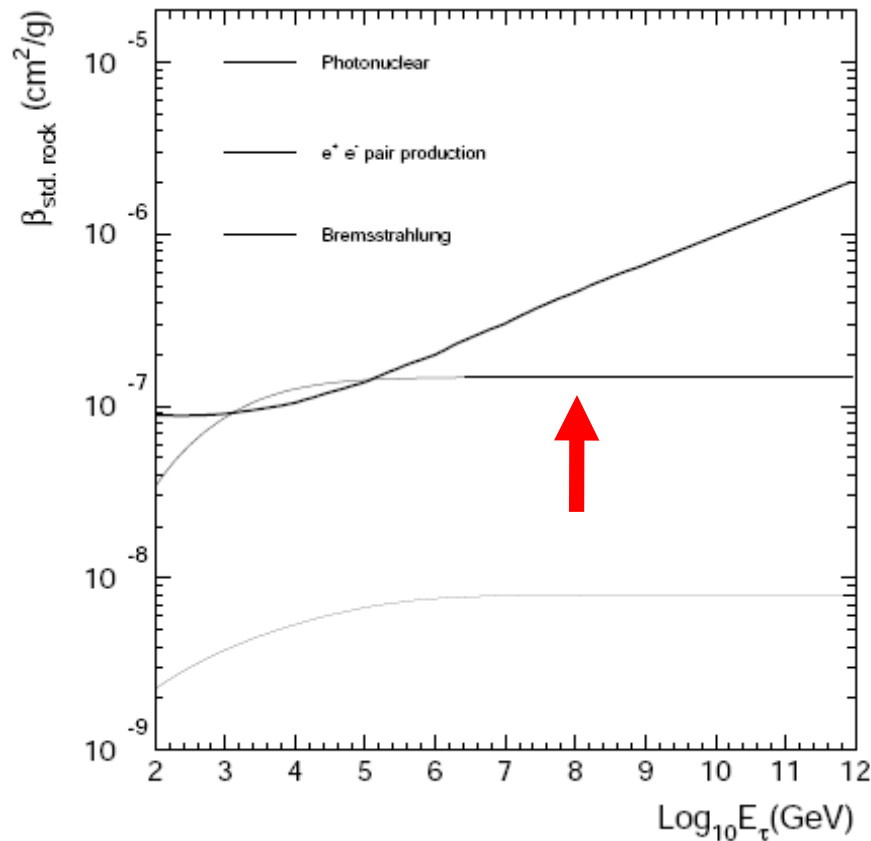
# Event rate

- tau neutrino flux
- $\nu N$  x-section (CC DIS)
- $\tau N$  x-section (photonuclear)
- $\tau$  decay probability
- Detector acceptance

# $\tau$ energy loss

$$-\left\langle \frac{dE}{dX} \right\rangle = a(E) + b(E)E$$


$b(E)$   
[g/cm<sup>2</sup>]<sup>-1</sup>



photonuclear

pair production

bremsstrahlung


From Aramo et al Astrop.Phys. 23 (2005) 65

(Standard Rock: A=22, Z=11, density=2.65 g/cm<sup>3</sup>)

# $\tau$ energy loss: photonuclear

$$b(E) = \frac{N_A}{A} \int dy y \int dQ^2 \frac{d\sigma^{lA}}{dQ^2 dy}$$

$$Q_{min}^2 = \frac{y^2 m_l^2}{1-y}, \quad Q_{max}^2 = 2m_p E y - 2m_\pi m_p - m_\pi^2,$$
$$y_{min} = \frac{2m_\pi m_p + m_\pi^2}{2m_p E}, \quad y_{max} = 1 - \frac{m_l}{E},$$


$$\frac{d\sigma^{lA}}{dQ^2 dy} = \frac{4\pi\alpha^2 F_2^A}{Q^4 y} \left[ 1 - y - \frac{Q^2}{4E^2} + \left( 1 - 2\frac{m_l^2}{Q^2} \right) \frac{y^2 + Q^2/E^2}{2(1 + R^A)} \right]$$

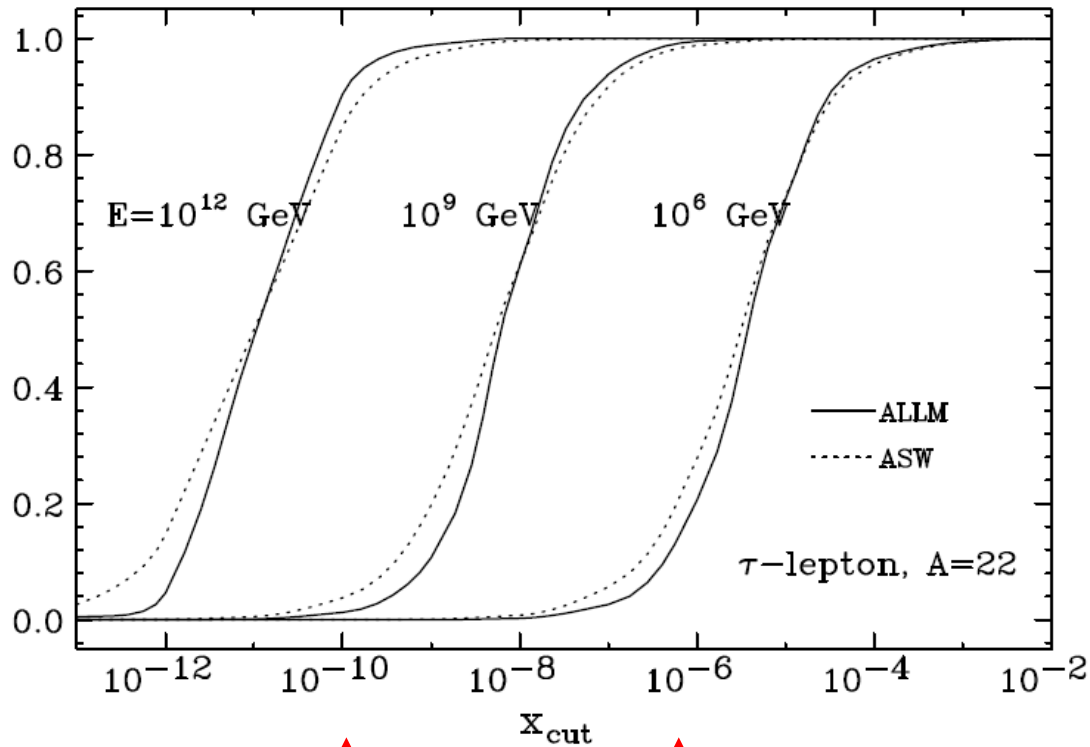
# Calculations of $b(E)$

- Generalized Vector Dominance (soft ) + Color Dipole (hard) Bezrukov, Bugaev 1981  
Bugaev, Shlepin 2003
- Regge theory (pomeron + reggeon) + pQCD (hard) Dutta et al. 2001 (ALLM)  
Butkevich, Mikheyev 2002 (CKMT)  
Kuzmin et al 2004 (CKMT)  
Petrukhin, Timashkov 2005
- Geometric scaling property found in  $\gamma^*p$  and  $\gamma^*N$  x-section data (saturation) Armesto, Merino, Parente, Zas 2008  
(ASW: Armesto, Salgado, Wiedemann 2005)



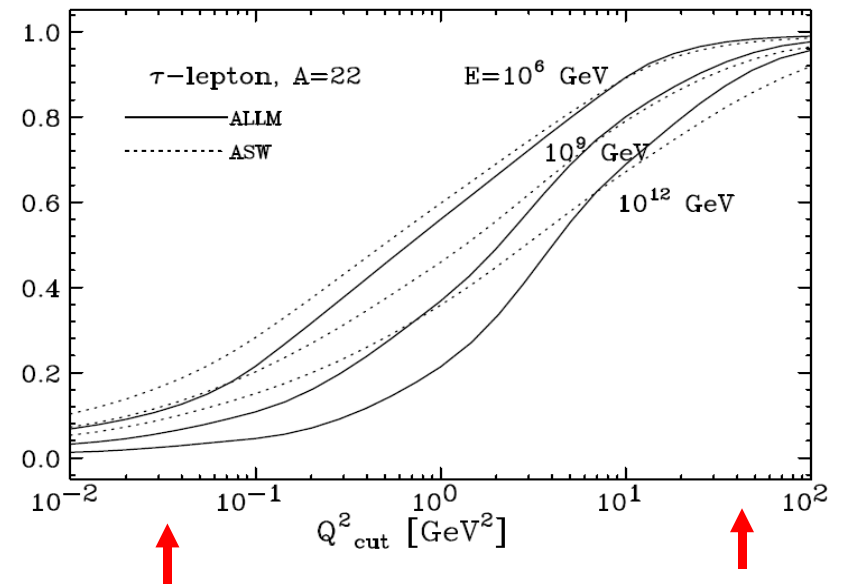
# The $\tau$ energy loss from different $X$ and $Q^2$ regions

$b(E) [x < x_{\text{cut}}]$

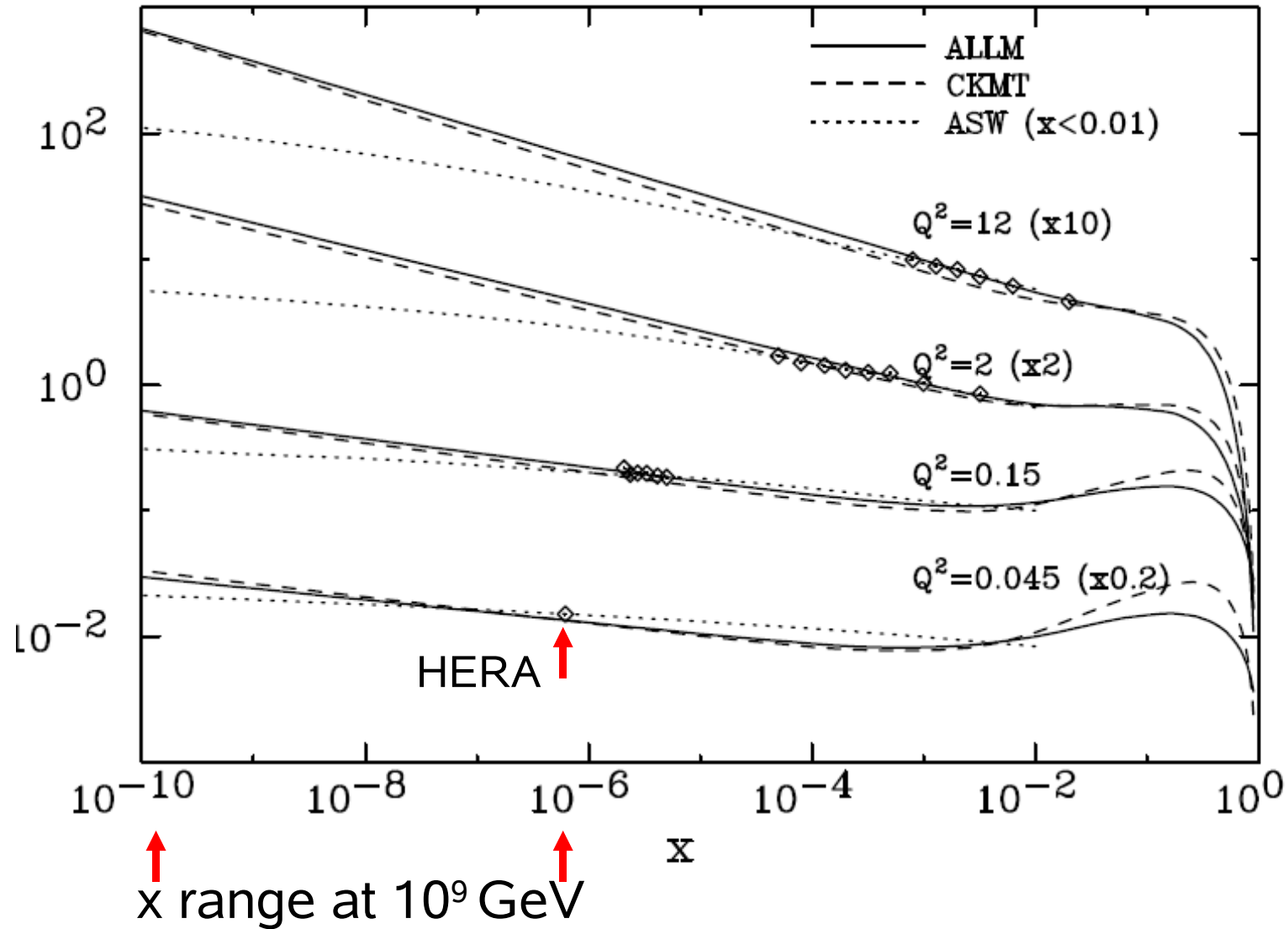


x range at  $10^9$  GeV

$b(E) [Q^2 < Q^2_{\text{cut}}]$

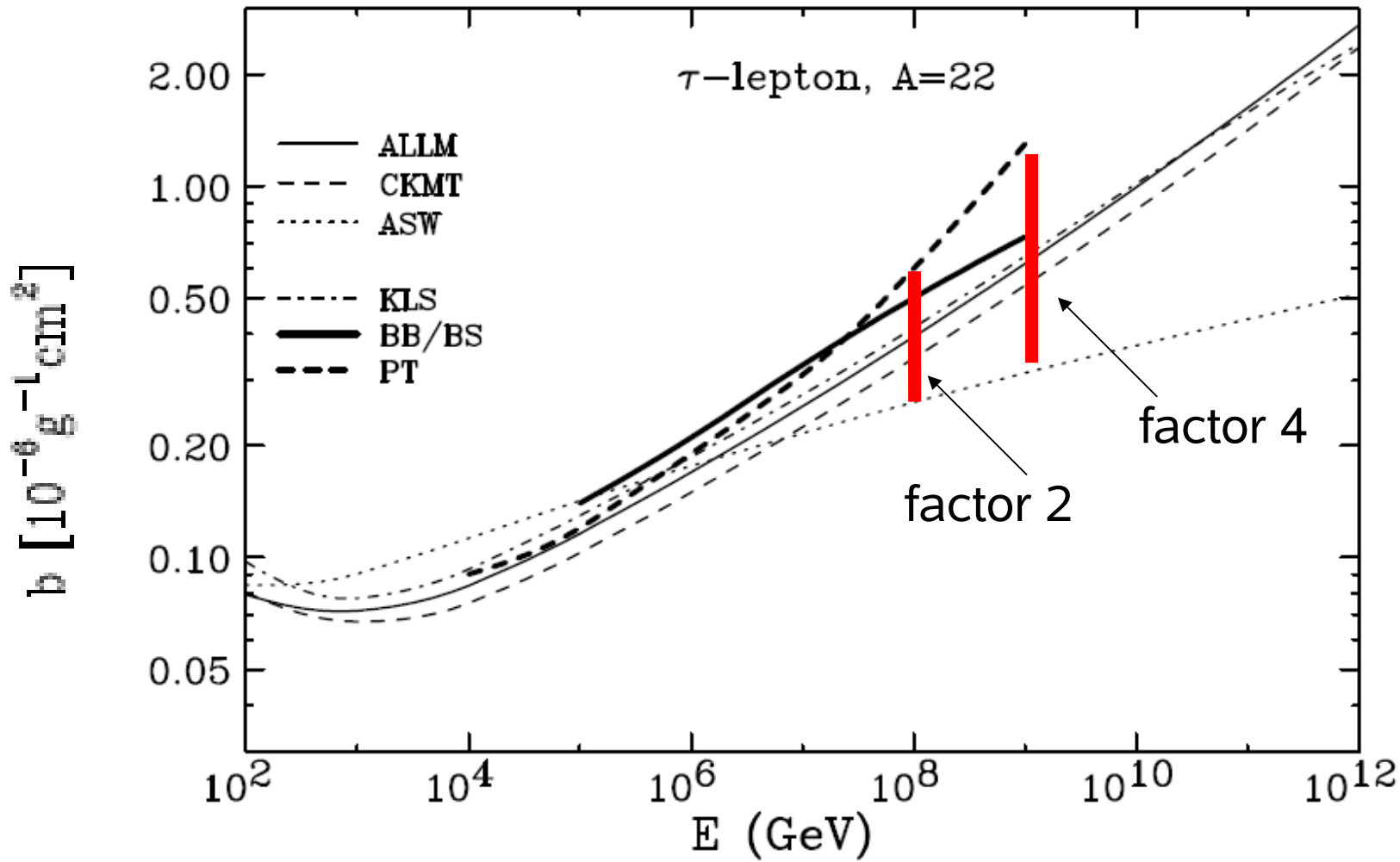


# $F_2$ proton at low and moderate $Q^2$



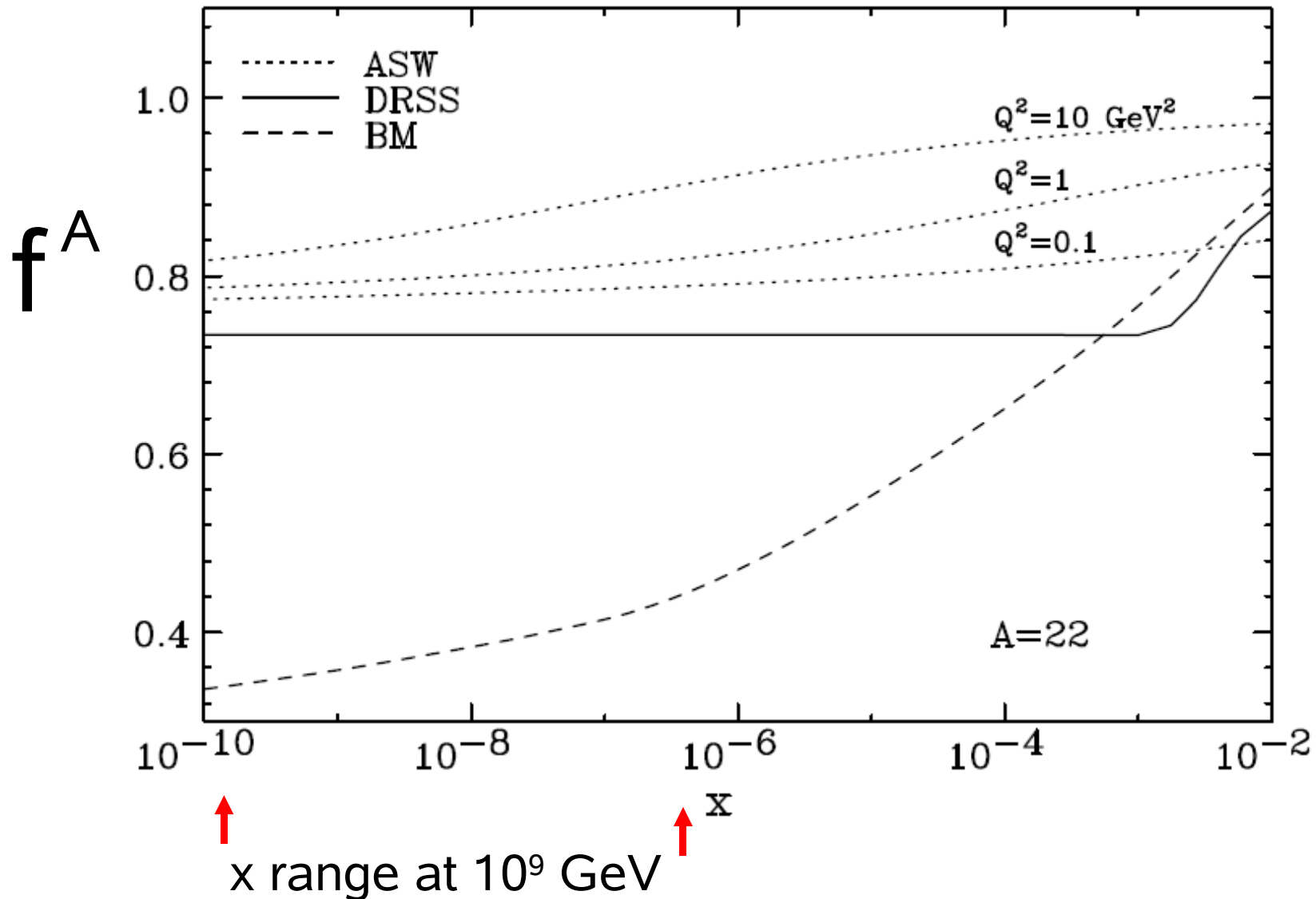
# $\tau$ energy loss

$$b(E) = -\frac{1}{E} \left\langle \frac{dE}{dX} \right\rangle = \frac{N_A}{A} \int dyy \int dQ^2 \frac{d\sigma^{lA}}{dQ^2 dy}, \quad (1)$$

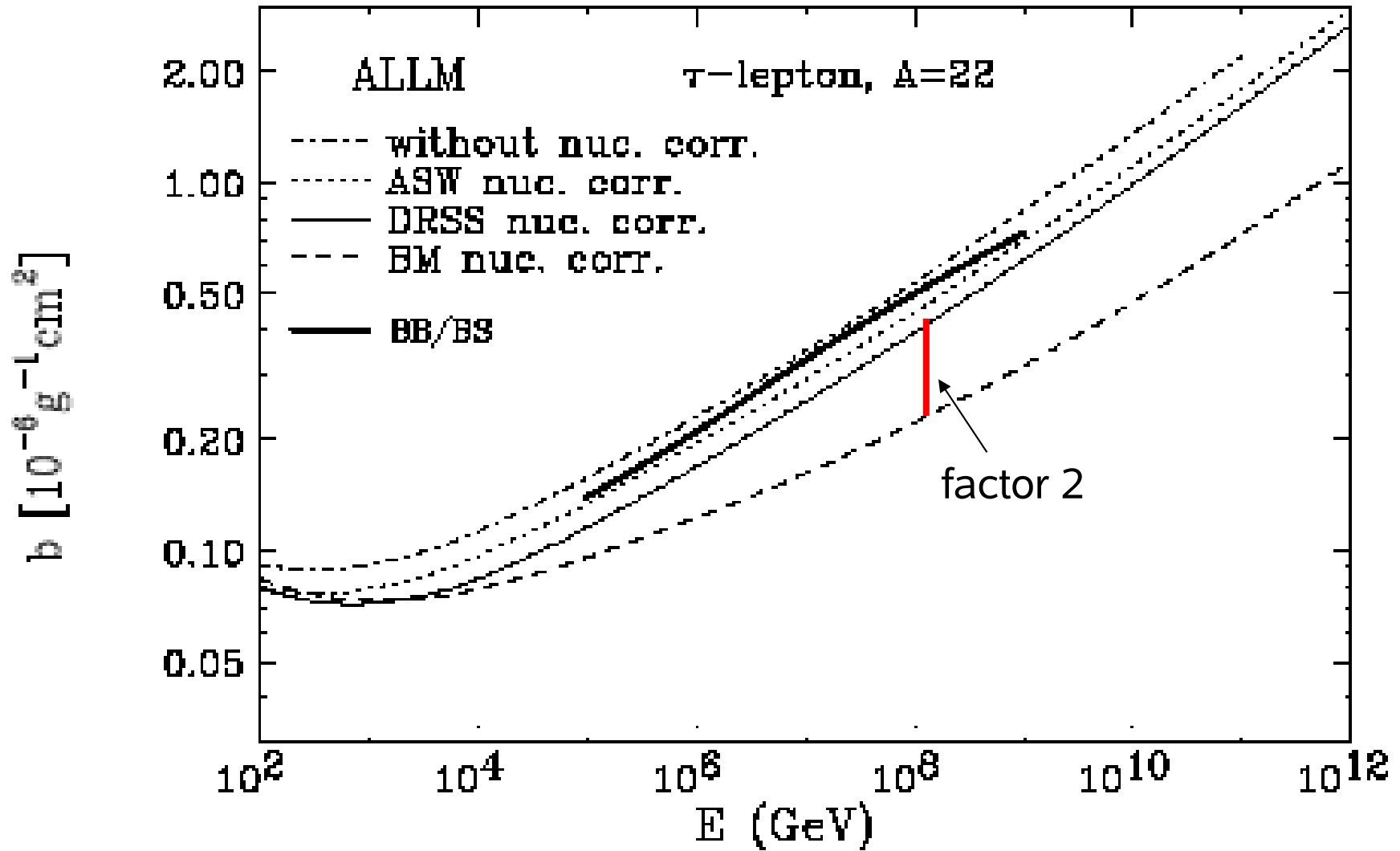


# Nuclear correction factor

$$F_2^A = f^A A F_2^P$$

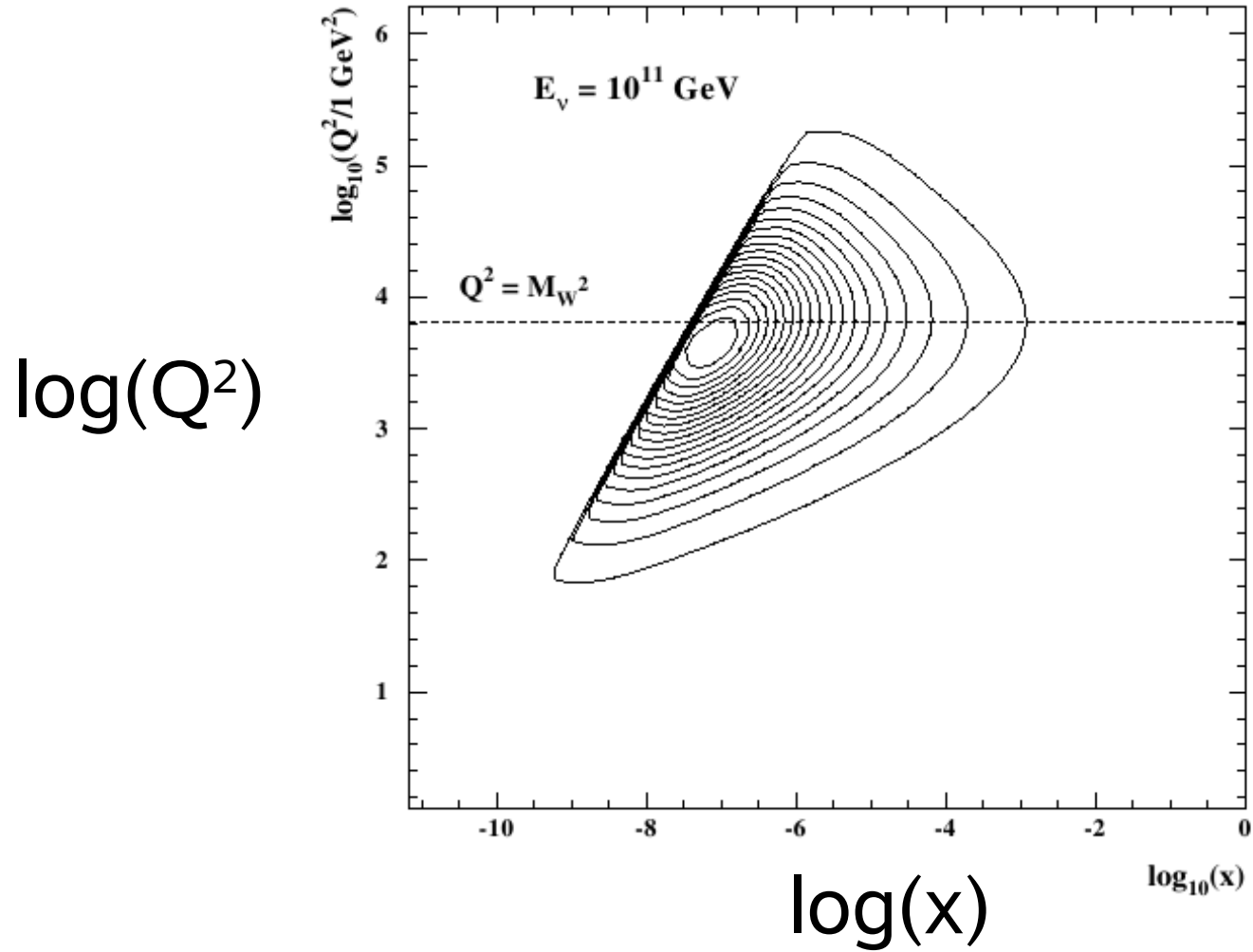


# The effect of nuclear corrections

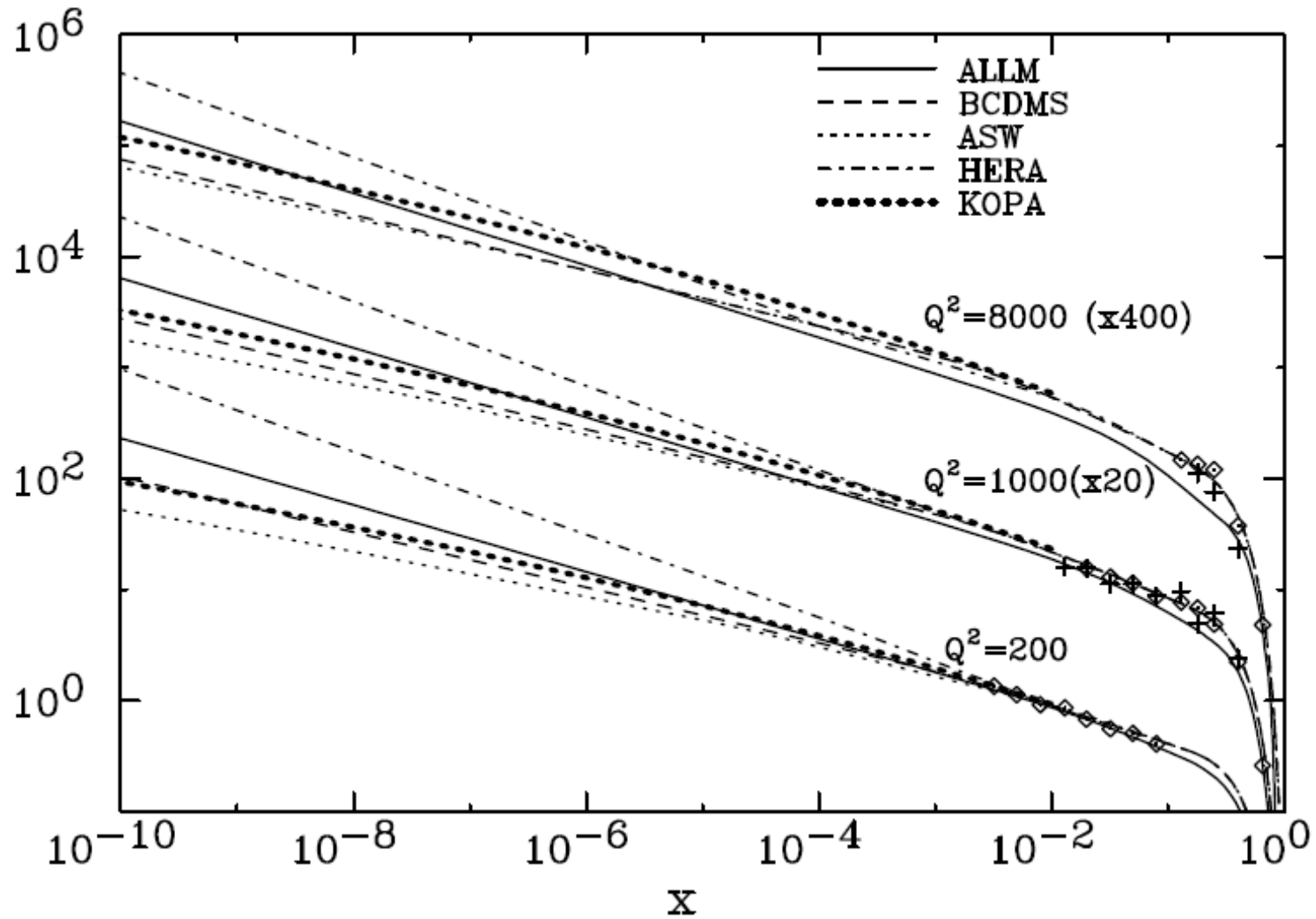


# Neutrino-nucleon CC DIS x-section

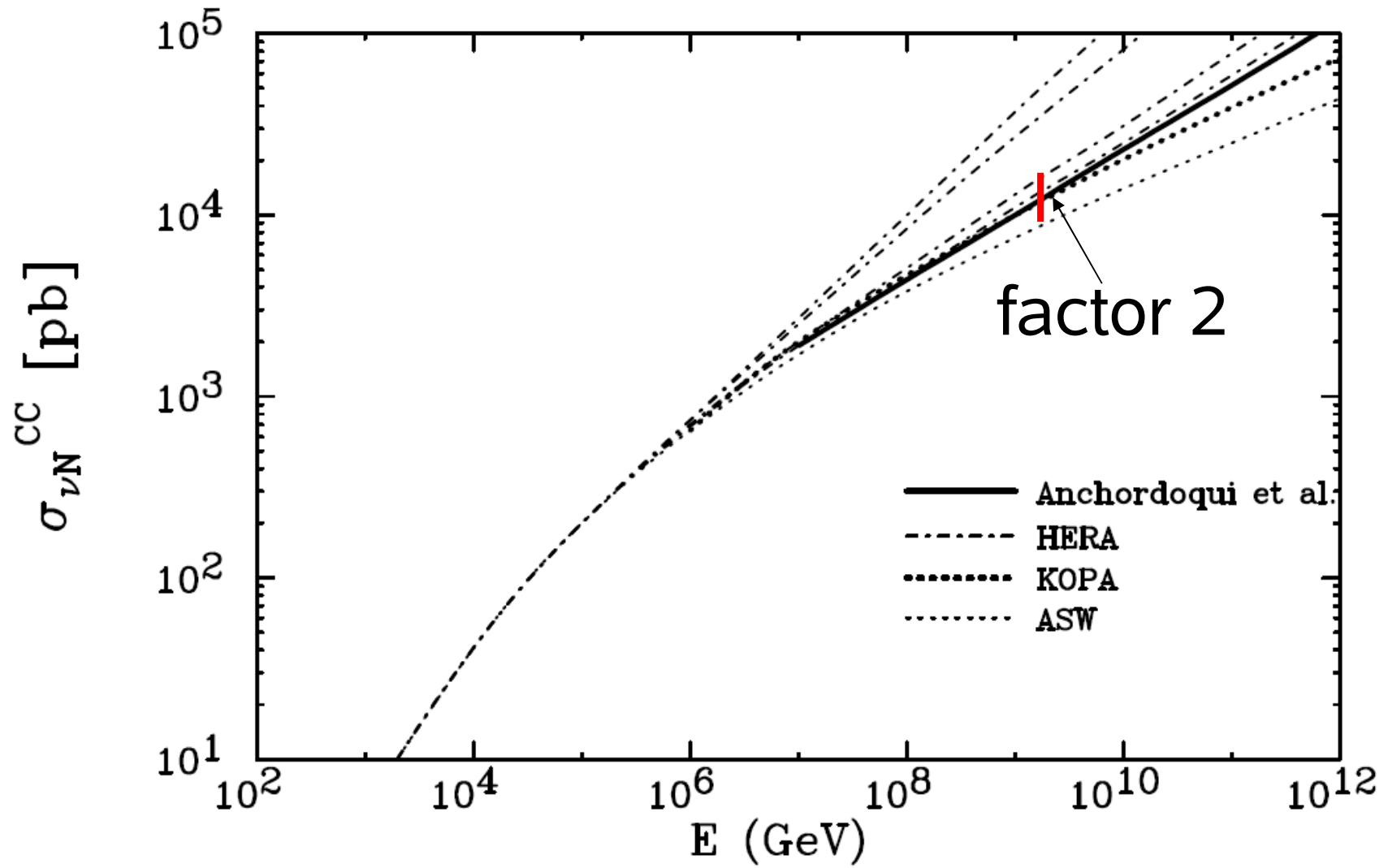
$$\frac{d\sigma_{CC}^{\nu N}}{dQ^2 dy} = \frac{G_F^2}{4\pi} \left( \frac{M_W^2}{M_W^2 + Q^2} \right)^2 \frac{F_2^{\nu N}}{y} [1 + (1 - y)^2]$$



# $F_2$ at high $Q^2$







# Summary

- The establishment of a tau neutrino bound from air shower data is affected by important systematic effects due to the uncertainty in  $F_2$  at low  $x$ .
- Several  $F_2$  models have been explored in the low  $x$  range. For the neutrino-nucleon interaction at high  $Q^2$  and for the tau interaction at low and moderate  $Q^2$ .
- For the tau energy loss at  $10^9$  GeV the difference between the two extreme predictions reaches a factor 4. For the neutrino cross section it is a factor 2.
- GVD and Regge based (+ hard pQCD component) calculations of the tau energy loss agree within a 30% and go parallel to all energies.
- A much stronger nuclear shadowing at small  $x$  lower  $b(E)$  by a factor up to 2 but is not expected to affect the neutrino cross section.