Theoretical developments in jet physics at hadron colliders

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- in collaboration with A. Gehrmann-De Ridder, T. Gehrmann, N.Glover

INCLUSIVE JET AND DIJET CROSS SECTIONS

- look at the production of jets of hadrons with large transverse energy in
  - inclusive jet events $pp \to j + X$
  - exclusive dijet events $pp \to 2j$
- cross sections measured as a function of the jet p_T, rapidity y and dijet invariant mass m_{jj} in double differential form

state of the art:

- dijet production is completely known in NLO QCD [Ellis, Kunszt, Soper '92], [Giele, Glover, Kosower '94], [Nagy '02]
- NLO+Parton shower [Alioli, Hamilton, Nason, Oleari, Re '11]
- threshold corrections [Kidonakis, Owens '00]

Goal:

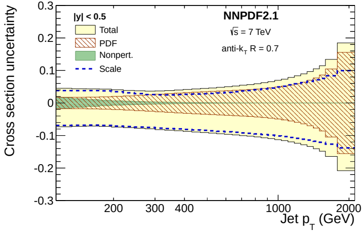
- obtain the jet cross sections at NNLO accuracy in double differential form

$$\frac{d^2\sigma}{dp_T d|y|} \qquad \frac{d^2\sigma}{dm_{jj} dy^*}$$

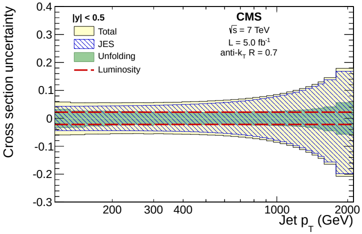
this talk:

- NNLO inclusive jet and dijet cross section (gluons only, leading colour)

Theoretical vs experimental uncertainties



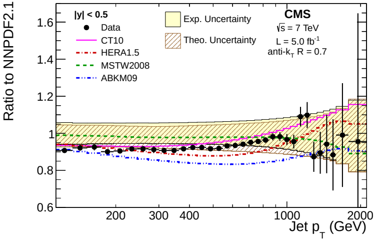
relative theoretical uncertainties for the inclusive jet production (NLO theory input) [CMS, arXiv:1212.6660]



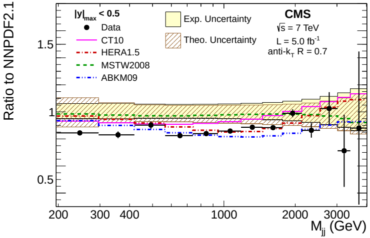
relative experimental uncertainties for the inclusive jet production [CMS, arXiv:1212.6660]

- residual uncertainty due to scale choice at NNLO expected at $\approx$ few percent level
- jet energy scale uncertainty has been determined to less than 5% for central jets $\rightarrow$ expect steady improvement with higher statistics
- theoretical prediction with the same precision as the experimental data $\rightarrow$ need pQCD predictions at NNLO accuracy

INCLUSIVE JET AND DIJET CROSS SECTIONS



(NLO theory input)
[CMS, arXiv:1212.6660]

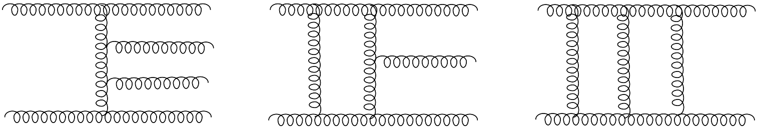


(NLO theory input)
[CMS, arXiv:1212.6660]

Phenomenological applications with NNLO:

- data can be used to constrain parton distribution functions
- size of NNLO correction important for precise determination of PDF's
- inclusion of jet data in NNLO parton distribution fits requires NNLO corrections to jet cross sections
- α_s determination from hadronic jet observables limited by the unknown higher order corrections

$pp \to 2j$ at NNLO: gluonic contributions



$A_6^{(0)}(gg \to gggg)$ $A_5^{(1)}(gg \to ggg)$ $A_4^{(2)}(gg \to gg)$

[Berends, Giele '87], [Mangano, Parke, Xu '87], [Britto, Cachazo, Feng '06]
[Bern, Dixon, Kosower '93]
[Anastasiou, Glover, Oleari, Tejeda-Yeomans '01],[Bern, De Freitas, Dixon '02]

$$\mathrm{d}\hat{\sigma}_{NNLO} = \int_{\mathrm{d}\Phi_4} \mathrm{d}\hat{\sigma}^{RR}_{NNLO} + \int_{\mathrm{d}\Phi_3} \mathrm{d}\hat{\sigma}^{RV}_{NNLO} + \int_{\mathrm{d}\Phi_2} \mathrm{d}\hat{\sigma}^{VV}_{NNLO}$$

- explicit infrared poles from loop integrations
- implicit poles in phase space regions for single and double unresolved gluon emission
- procedure to extract the infrared singularities and assemble all the parts in a parton-level generator

NNLO ANTENNA SUBTRACTION

$$
\begin{aligned}
\mathrm{d}\hat{\sigma}_{NNLO} &= \int_{\mathrm{d}\Phi_4} \left(\mathrm{d}\hat{\sigma}^{RR}_{NNLO} - \mathrm{d}\hat{\sigma}^{S}_{NNLO} \right) \\
&+ \int_{\mathrm{d}\Phi_3} \left(\mathrm{d}\hat{\sigma}^{RV}_{NNLO} - \mathrm{d}\hat{\sigma}^{T}_{NNLO} \right) \\
&+ \int_{\mathrm{d}\Phi_2} \left(\mathrm{d}\hat{\sigma}^{VV}_{NNLO} - \mathrm{d}\hat{\sigma}^{U}_{NNLO} \right)
\end{aligned}
$$

- $\mathrm{d}\hat{\sigma}^{S}_{NNLO}$: real radiation subtraction term for $\mathrm{d}\hat{\sigma}^{RR}_{NNLO}$
- $\mathrm{d}\hat{\sigma}^{T}_{NNLO}$: one-loop virtual subtraction term for $\mathrm{d}\hat{\sigma}^{RV}_{NNLO}$
- $\mathrm{d}\hat{\sigma}^{U}_{NNLO}$: two-loop virtual subtraction term for $\mathrm{d}\hat{\sigma}^{VV}_{NNLO}$
- subtraction terms constructed using the antenna subtraction method at NNLO for hadron colliders $\rightarrow$ presence of initial state partons to take into account
- contribution in each of the round brackets is finite, well behaved in the infrared singular regions and can be evaluated numerically

NNLO ANTENNA SUBTRACTION

[A. Gehrmann-De Ridder, T. Gehrmann, N. Glover, JP]

Implementation checks (gluons only channel at leading colour in $pp \to 2j$):

- subtraction terms correctly approximate the matrix elements in all unresolved configurations of partons j, k

$$\boxed{\mathrm{d}\hat{\sigma}^{RR,RV}_{NNLO} \xrightarrow{\forall \{j,k\},\{j\}\to 0} \mathrm{d}\hat{\sigma}^{S,T}_{NNLO}}$$

- local (pointwise) analytic cancellation of all infrared explicit ϵ-poles when integrated subtraction terms are combined with one, two-loop matrix elements

$$\boxed{\mathcal{P}oles\left(\mathrm{d}\hat{\sigma}^{RV}_{NNLO} - \mathrm{d}\hat{\sigma}^{T}_{NNLO}\right) = 0}$$

$$\boxed{\mathcal{P}oles\left(\mathrm{d}\hat{\sigma}^{VV}_{NNLO} - \mathrm{d}\hat{\sigma}^{U}_{NNLO}\right) = 0}$$

- process independent NNLO subtraction scheme
- singularities in intermediate steps cancel analytically
- allows the computation of multiple differential distributions in a single program run

NUMERICAL SETUP

[A. Gehrmann-De Ridder, T. Gehrmann, N. Glover, JP]

- jets identified with the anti-k_T jet algorithm
- jets accepted at rapidities $|y| < 4.4$
- leading jet with transverse momentum $p_t > 80$ GeV
- subsequent jets required to have at least $p_t > 60$ GeV
- MSTW2008nnlo PDF
- dynamical factorization and renormalization central scale equal to the leading jet p_T ($\mu_R = \mu_F = \mu = p_{T1}$)

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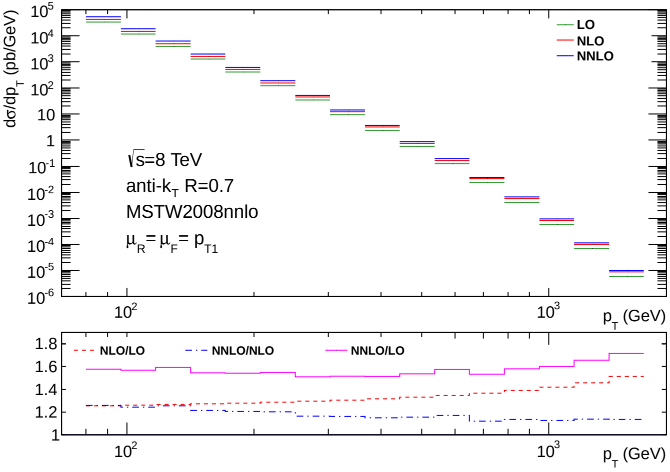
Integrated cross section results (gluons only channel) with scale variations

	$\sigma^{8TeV-LO}_{incl.jet}(pb)$	$\sigma^{8TeV-NLO}_{incl.jet}(pb)$	$\sigma^{8TeV-NNLO}_{incl.jet}(pb)$
$\mu = 0.5p_{T1}$	$(12.586 \pm 0.001) \times 10^5$	$(11.299 \pm 0.001) \times 10^5$	$(15.33 \pm 0.03) \times 10^5$
$\mu = p_{T1}$	$(9.6495 \pm 0.001) \times 10^5$	$(12.152 \pm 0.001) \times 10^5$	$(15.20 \pm 0.02) \times 10^5$
$\mu = 2.0p_{T1}$	$(7.5316 \pm 0.001) \times 10^5$	$(11.824 \pm 0.001) \times 10^5$	$(15.21 \pm 0.01) \times 10^5$

- NNLO result increased by about 25% with respect to the NLO cross section

INCLUSIVE JET p_T DISTRIBUTION

[A. Gehrmann-De Ridder, T. Gehrmann, N. Glover, JP]



- NNLO effect stabilizes the NLO k-factor growth with p_T

INCLUSIVE JET p_T DISTRIBUTION

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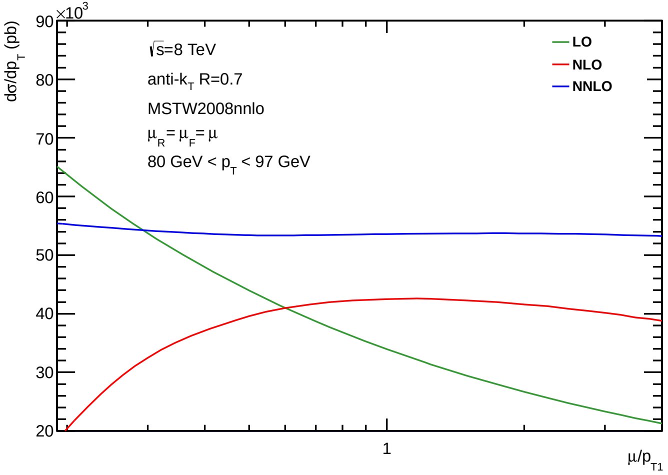
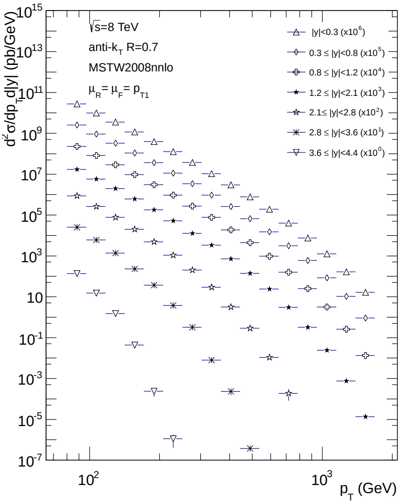
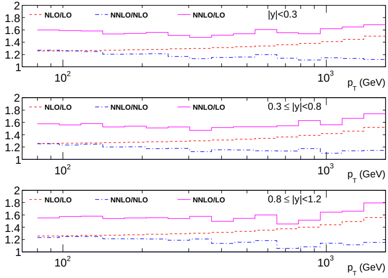


- flat scale dependence at NNLO

INCLUSIVE JET p_T DISTRIBUTION $\boxed{R = 0.7}$

- double differential inclusive jet p_T distribution at NNLO (gluons only)



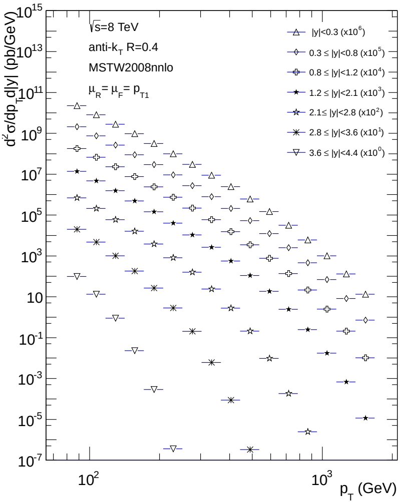


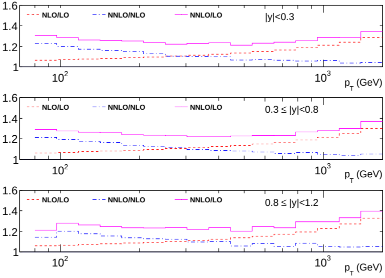
double differential k-factors

- NNLO result varies between 25% to 12% with respect to the NLO cross section
- similar behaviour between the rapidity slices

INCLUSIVE JET p_T DISTRIBUTION $R = 0.4$

- double differential inclusive jet p_T distribution at NNLO (gluons only)



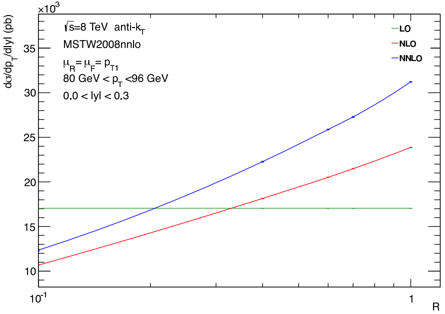


double differential k-factors

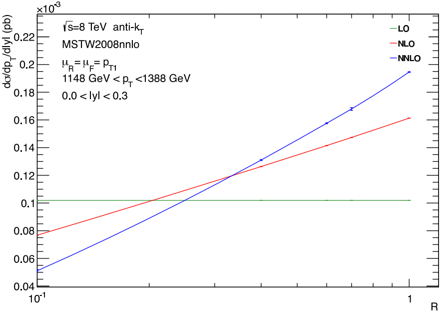
- NNLO result varies between 20% to 5% with respect to the NLO cross section
- similar behaviour between the rapidity slices

INCLUSIVE JET p_T DISTRIBUTION

$80\text{GeV} < p_T < 96\text{GeV}$



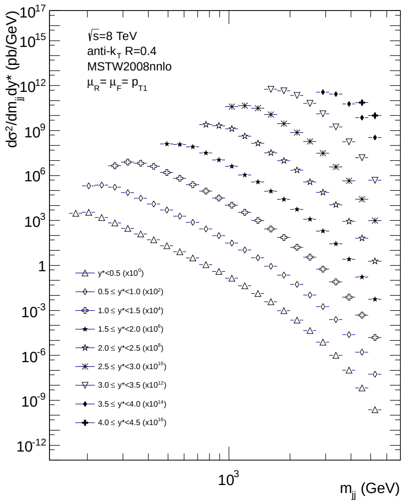
$1148\text{GeV} < p_T < 1388\text{GeV}$

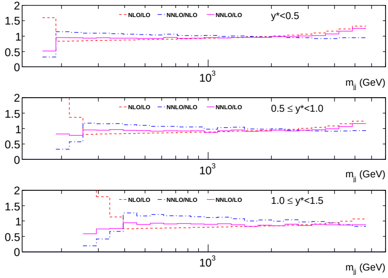


- inclusive jet cross section versus R
- can the NNLO cross section describe data for different values of R simultaneously at low and high jet p_T?

EXCLUSIVE DIJET MASS DISTRIBUTION $R = 0.4$

- double differential dijet mass distribution at NNLO (gluons only)





double differential k-factors

- NNLO corrections up to 20% with respect to the NLO cross section
- corrections increase slightly for large $y^* = 1/2|y_1 - y_2|$

CONCLUSIONS

- antenna subtraction method generalised for the calculation of NNLO QCD corrections for exclusive collider observables with partons in the initial-state
- non-trivial check of analytic cancellation of infrared singularities between double-real, real-virtual and double-virtual corrections
- proof-of principle implementation of the $gg \to gg$ contribution to $pp \to 2j$ at NNLO in the new NNLOJET parton-level generator
- computation of multiple differential distributions at NNLO in a single program run $\to$ experimentalists input welcome

Future work:

- go beyond gluons only leading colour approximation
- include remaining partonic subprocesses involving quarks