

Vector Meson Production from Gauge/Gravity Duality

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work with Miguel S. Costa and Nick Evans

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Outline

Introduction

Pomeron in AdS

Vector Meson Production

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Data Analysis

Conclusions

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- ▶ However, at lower energies, once it is of order Λ_{QCD} the coupling is very strong and we cannot use pQCD.
- ▶ Our goal is to study the strong interaction at strong coupling.
- ▶ More specifically, a recent conjecture by Maldacena relating string theory on $AdS_5 \times S_5$ to $\mathcal{N} = 4SYM$ allows us to study QCD at strong coupling.

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$$A(s, t) \sim s^{\alpha(t)}, \quad \alpha(t) = \alpha(0) + \frac{\alpha' t}{2},$$

- ▶ In perturbative QCD, the propagation of the Pomeron is given by the BFKL equation.

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- ▶ Correspondence works in the limit

$$N_C \rightarrow \infty, \quad \lambda = g^2 N_C = R^4/\alpha'^2 \gg 1, \text{ fixed}$$

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$$\mathcal{K} = \frac{2(zz')^2 s}{g_0^2 R^4} \chi(s, b, z, z')$$

where

$$\chi(\tau, L) = \left(\cot\left(\frac{\pi\rho}{2}\right) + i \right) g_0^2 e^{(1-\rho)\tau} \frac{L}{\sinh L} \frac{\exp\left(\frac{-L^2}{\rho\tau}\right)}{(\rho\tau)^{3/2}}$$

- ▶ The weak and strong coupling Pomeron exchange kernels have a remarkably similar form.

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- ▶ At $t = 0$

Weak coupling:

$$\mathcal{K}(k_{\perp}, k'_{\perp}, s) = \frac{s^{j_0}}{\sqrt{4\pi\mathcal{D}\log s}} e^{-(\log k_{\perp} - \log k'_{\perp})^2 / 4\mathcal{D}\log s}$$

$$j_0 = 1 + \frac{\log 2}{\pi^2} \lambda, \quad \mathcal{D} = \frac{14\zeta(3)}{\pi} \lambda / 4\pi^2$$

Strong coupling:

$$\mathcal{K}(z, z', s) = \frac{s^{j_0}}{\sqrt{4\pi\mathcal{D}\log s}} e^{-(\log z - \log z')^2 / 4\mathcal{D}\log s}$$

$$j_0 = 2 - \frac{2}{\sqrt{\lambda}}, \quad \mathcal{D} = \frac{1}{2\sqrt{\lambda}}$$

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$$A(s, t) = 2is \int d^2l e^{-i\mathbf{l}_\perp \cdot \mathbf{q}_\perp} \int dz d\bar{z} P_{13}(z) P_{24}(\bar{z}) (1 - e^{i\chi(s, b, z, \bar{z})})$$

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- ▶ We can study different scattering processes by supplying P_{13} and P_{24} .
- ▶ For example, already applied to DIS [Brower, MD, Sarčević, Tan; Cornalba, Costa, Penedones], and DVCS [Costa, MD].

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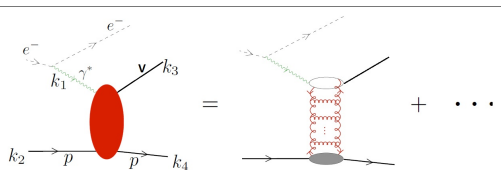
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What is Vector Meson Production?

Vector meson production occurs in the scattering between an offshell photon and a proton.

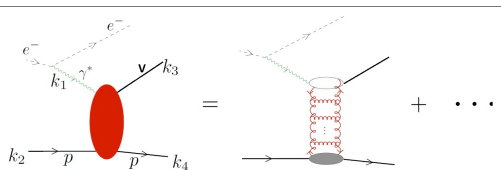
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The vector mesons consist of a quark-antiquark pair, and have the same quantum numbers as the photon, $J^{PC} = 1^{--}$. The production of the ρ^0 , ω , ϕ and J/Ψ was measured at HERA.

- ▶ We are interested in calculating the differential and exclusive cross sections

$$\frac{d\sigma}{dt}(x, Q^2, t) = \frac{|W|^2}{16\pi s^2},$$

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$$W = 2isQQ' \int dl_{\perp} e^{iq_{\perp} \cdot l_{\perp}} \int \frac{dz}{z^3} \frac{d\bar{z}}{\bar{z}^3} \Psi(z) \Phi(\bar{z}) \left[1 - e^{i\chi(S,L)} \right].$$

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- ▶ In this analysis we use

$$\Psi_n(z) = -\left(\sqrt{C \frac{\pi^2}{6}} z^2 K_n(Qz)\right) \left(\frac{\sqrt{2}}{\xi J_1(\xi)} z^2 J_n(mz)\right), \quad \Phi(\bar{z}) = \bar{z}^3 \delta(\bar{z} - z_{\star})$$

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- ▶ C is the aforementioned normalization, and g_0^2 is related to the coupling of the external states to the pomeron.

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- ▶ Similarly, the $t = 0$ result for the hard-wall model can also be written explicitly

$$\chi_{hw}(\tau, t = 0, z, \bar{z}) = \chi(\tau, 0, z, \bar{z}) + \mathcal{F}(\tau, z, \bar{z}) \chi(\tau, 0, z, z_0^2/\bar{z}).$$

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- ▶ When $t \neq 0$, we will use an approximation

$$\chi_{hw}(\tau, l, z, \bar{z}) = C(\tau, z, \bar{z}) D(\tau, l) \chi_{hw}^{(0)}(\tau, l, z, \bar{z})$$

► The function

$$\mathcal{F}(\tau, z, \bar{z}) = 1 - 4\sqrt{\pi\tau} e^{\eta^2} \operatorname{erfc}(\eta), \quad \eta = \frac{-\log(z\bar{z}/z_0^2) + 4\tau}{\sqrt{4\tau}}$$

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- ▶ It is therefore in these regions that confinement is important!

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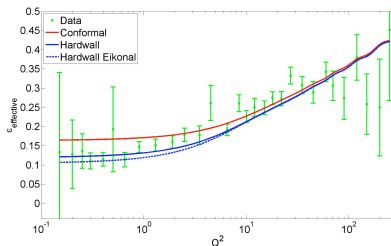
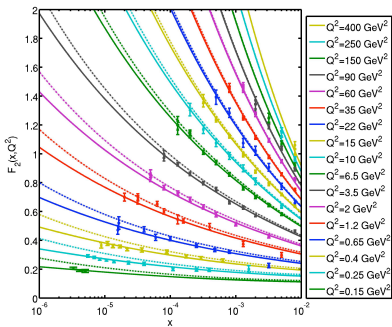
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- ▶ We will look at both the differential and total exclusive cross sections.

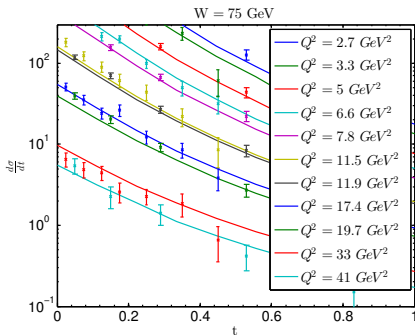
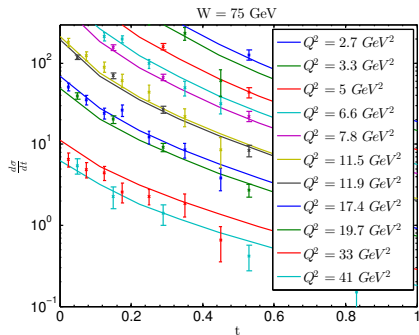
- ▶ Note that the same formalism has been applied before to DIS with good results ($\chi^2 = 1.04$ for the best model) [Brower, MD, Sarčević, Tan, 2010, Cornalba, Costa, Penedones, 2010, Brower Moriond 2011], and DVCS ($\chi^2 = 1.00$ and $\chi^2 = 0.51$ for the best models of the cross section and differential cross section respectively) [Costa, MD, 2012, MD Moriond 2012].

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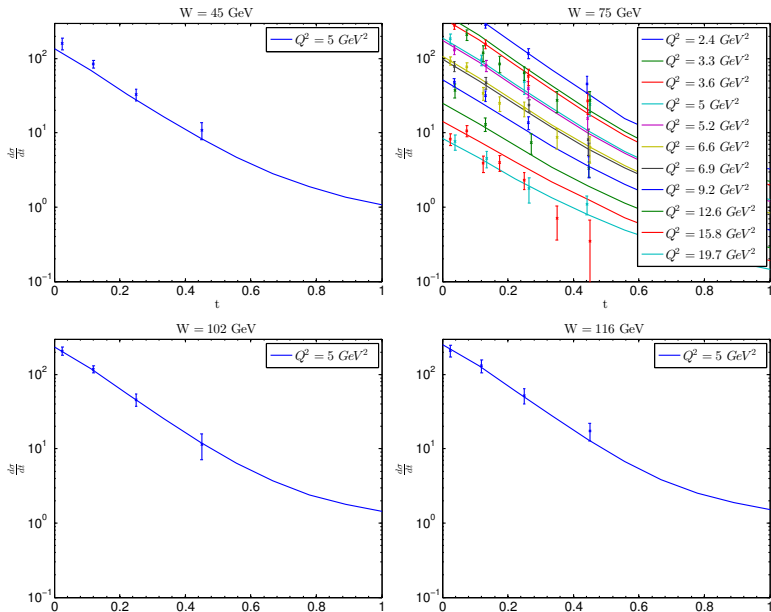


		σ				$d\sigma/dt$		
		ρ	ϕ	Ω	J/ψ	ρ	ϕ	J/ψ
C o n f o r m a l	m	0.77549	1.019445	0.78265	3.096916	0.77549	1.019445	3.096916
	N	48	27	6	38	35	21	84
	χ^2	0.9234	0.6002	0.0099	0.2844	1.7387	1.2732	2.8818
	$g0^2$	2.29	0.5742	0.2673	0.3946	0.7814	0.0805	0.3565
	$2/\sqrt{\lambda}$	0.76	0.7339	0.6416	0.697	0.6473	0.5443	0.7165
	z^*	3.4074	3.0012	1.8355	0.9823	2.1453	2.5445	2.1536
H a r d w a i l	χ^2	0.8819	0.6131	0.015	0.6285*	1.6574	1.3595	1.8442
	$g0^2$	2.0438	0.5559	0.3335	3.1893	24.4179	2.6638	9.6671
	$2/\sqrt{\lambda}$	0.758	0.7321	0.6589	0.7396	0.6946	0.5905	0.7539
	z^*	3.5947	3.6341	1.4668	2.1019	2.1847	2.5064	2.4172
	$z0$	4.8164	4.3625	7.2955	4.2519	7.6918	8.5684	4.6465

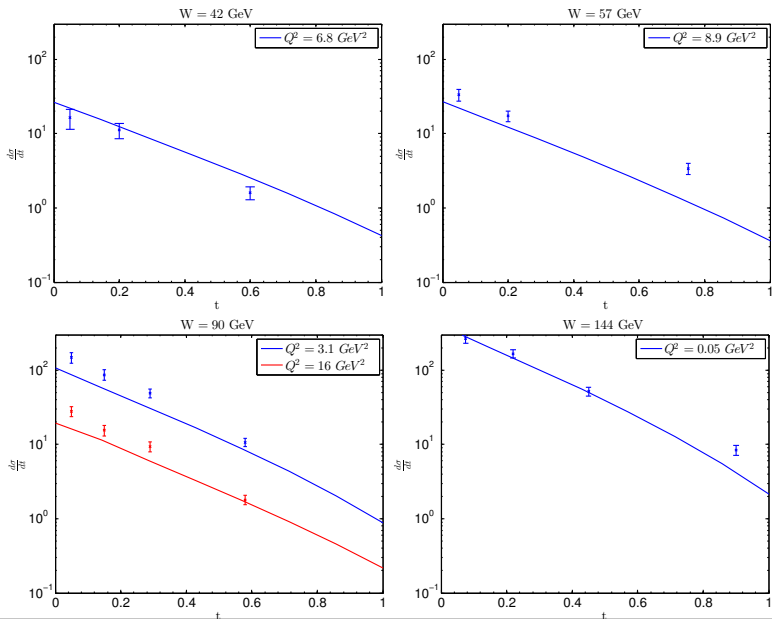
Differential cross section for the ρ meson:



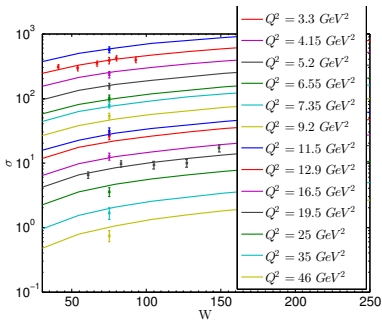
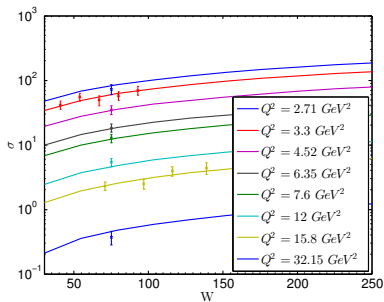
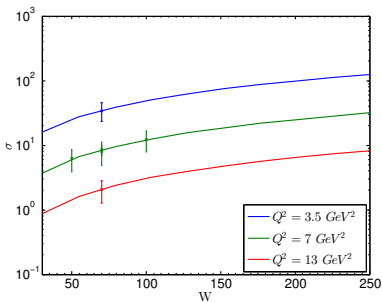
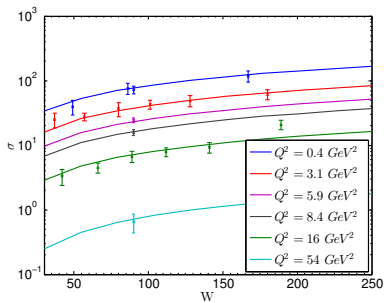
Differential cross section for the ϕ meson (hardwall model):



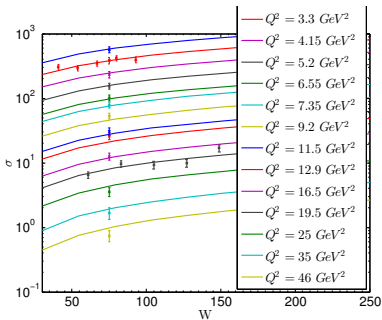
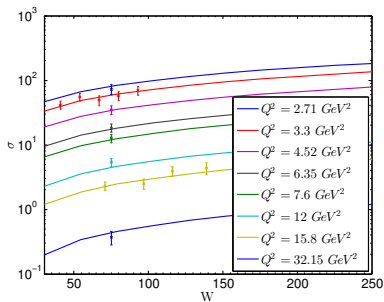
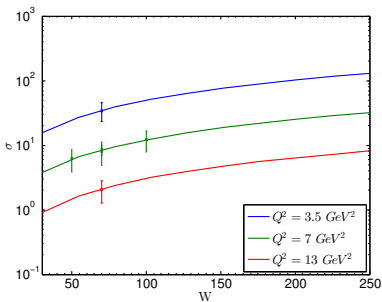
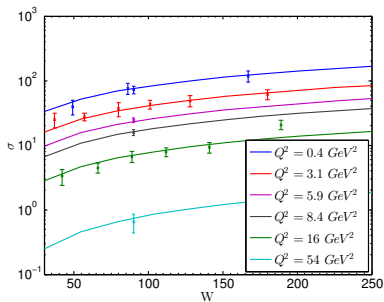
Differential cross section for the J/Ψ meson (hardwall model):



Cross sections for the conformal model:



Cross sections for the hardwall model:



Outline

Introduction

Pomeron in AdS

Vector Meson Production

Models

Data Analysis

Conclusions

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- ▶ It might therefore be possible to extend some of the insights we gain even into the weak coupling regime.
- ▶ The hard wall model, although a simple modification of AdS, seems to capture effects of confinement well. Interesting to repeat some of the calculations using a different confinement model to identify precisely what features are model independent.

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- ▶ We can also try to use a different AdS model of confinement (for example the soft wall model) and combine our methods with work by others (for example on the vector meson wavefunctions).

Thank you!