

# Scalar Dark Matter candidate in Multi-Higgs Models

Dorota Sokołowska

University of Warsaw, Faculty of Physics,  
Institute of Theoretical Physics

Moriond QCD, La Thuile, 26.03.2017

based on

**JHEP 1612 (2016) 014** and work in progress  
with A. Cordero-Cid, J. Hernandez-Sanchez, V. Keus,  
S. F. King, S. Moretti, D. Rojas

# The Standard Model

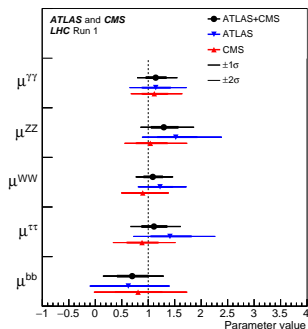
A rigorously tested Theory of Fundamental Interactions

From the LHC:

- a Higgs particle found in 2012
- no significant deviation from the SM
- no sign of New Physics

But no explanation for:

- Dark Matter
- neutrino masses
- baryon asymmetry and baryogenesis
- extra source of CP violation
- vacuum stability
- ...



JHEP 08 (2016) 045

## Dark Matter

Evidence for Dark Matter at diverse scales:

- **galaxy scales**: rotational speeds of galaxies
- **cluster scales**: gravitational lensing at galaxy clusters
- **horizon scales**: anisotropies in the CMB

⇒ **around 25 % of the Universe is:**

- cold
- non-baryonic
- neutral
- very weakly interacting

⇒ **Weakly Interacting Massive Particle**

- **stable due to the discrete symmetry**

$$\underbrace{\text{DM DM} \rightarrow \text{SM SM}}_{\text{pair annihilation}}, \quad \underbrace{\text{DM} \not\rightarrow \text{SM}, \dots}_{\text{stable}}$$

- annihilation cross-section  $\langle\sigma v\rangle \propto \text{EW interaction}$
- thermal evolution of DM density – a fixed value after freeze-out

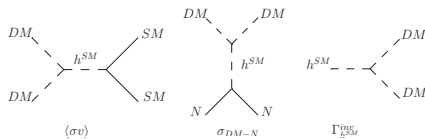
## Higgs-portal DM

Simplest realisation: the SM with  $\Phi_{SM} + Z_2$ -odd scalar  $S$ :

$S \rightarrow -S$ , SM fields  $\rightarrow$  SM fields

$$\mathcal{L} = \mathcal{L}_{SM} + \frac{1}{2}(\partial S)^2 - \frac{1}{2}m_{DM}^2 S^2 - \lambda_{DM} S^4 - \lambda_{SSh} \Phi_{SM}^2 S^2$$

SM sector  $\overset{\text{Higgs}}{\longleftrightarrow}$  DM sector



given by the same coupling

Strong constraints from relic density + direct detection + Higgs decays

$\Rightarrow$  modified Higgs-portal-type DM candidates in multi-scalar models

in this talk: focus on **DM phenomenology** in the  $I(2+1)HDM$  i.e.

*3HDM with Two Inert and One Higgs doublet*

## I(2+1)HDM

$Z_2$ -symmetry in I(2+1)HDM:

$$\phi_1 \rightarrow -\phi_1, \quad \phi_2 \rightarrow -\phi_2, \quad \phi_3 \rightarrow \phi_3, \quad \text{SM fields} \rightarrow \text{SM fields}$$

$Z_2$ -invariant potential:

$$\begin{aligned}
 V = \sum_i^3 & \left[ -|\mu_i^2|(\phi_i^\dagger \phi_i) + \lambda_{ii}(\phi_i^\dagger \phi_i)^2 \right] + \sum_{ij}^3 \left[ \lambda_{ij}(\phi_i^\dagger \phi_i)(\phi_j^\dagger \phi_j) + \lambda'_{ij}(\phi_i^\dagger \phi_j)(\phi_j^\dagger \phi_i) \right] \\
 & + \left( -\mu_{12}^2(\phi_1^\dagger \phi_2) + \lambda_1(\phi_1^\dagger \phi_2)^2 + \lambda_2(\phi_2^\dagger \phi_3)^2 + \lambda_3(\phi_3^\dagger \phi_1)^2 + h.c. \right) \\
 & + \left( \lambda_4(\phi_3^\dagger \phi_1)(\phi_2^\dagger \phi_3) + \lambda_5(\phi_1^\dagger \phi_2)(\phi_3^\dagger \phi_3) + \lambda_6(\phi_1^\dagger \phi_2)(\phi_1^\dagger \phi_1) \right. \\
 & \quad \left. + \lambda_7(\phi_1^\dagger \phi_2)(\phi_2^\dagger \phi_2) + \lambda_8(\phi_3^\dagger \phi_1)(\phi_3^\dagger \phi_2) + h.c. \right)
 \end{aligned}$$

- 21 parameters in  $V$
- $\mu_{12}^2, \lambda_1, \lambda_2, \lambda_3$  are complex
- Yukawa interaction: "Model I"-type (only  $\phi_3$  couples to fermions)
- explicit  $Z_2$ -symmetry

## Parameters of $V$

- $\mu_3^2 = v^2 \lambda_{33} = m_h^2/2$  fixed from extremum conditions
- "dark democracy":  $\mu_1^2 = \mu_2^2$ ,  $\lambda_{13} = \lambda_{23}$ ,  $\lambda'_{13} = \lambda'_{23}$ ,  $\lambda_3 = \lambda_2$ , e.g.  

$$\lambda_2(\phi_2^\dagger \phi_3)^2 + \lambda_3(\phi_3^\dagger \phi_1)^2 + h.c. \rightarrow \lambda_2 \left( (\phi_2^\dagger \phi_3)^2 + (\phi_3^\dagger \phi_1)^2 + h.c. \right)$$
- $\left( \lambda_4(\phi_3^\dagger \phi_1)(\phi_2^\dagger \phi_3) + \lambda_5(\phi_1^\dagger \phi_2)(\phi_3^\dagger \phi_3) + \dots \right)$ : no new phenomenology  
 $\Rightarrow \lambda_{4-8} = 0$
- $\lambda_1, \lambda_{11,22,12}, \lambda'_{12}$  – self-interactions of inert doublets

### 21 parameters $\rightarrow$ 7 important parameters

- $\mu_2^2$  – mass scale of inert particles
- $\mu_{12}^2 = |\mu_{12}^2| e^{i\theta_{12}}$ ,  $\lambda_2 = |\lambda_2| e^{i\theta_2}$  – mass splittings and CPv phase
- $\lambda_2, \lambda_{23}, \lambda'_{23}$  – DM-Higgs coupling

## DM in I(2+1)HDM

$Z_2$ -invariant vacuum state:

$$\phi_1 = \begin{pmatrix} H_1^+ \\ \frac{H_1^0 + iA_1^0}{\sqrt{2}} \end{pmatrix}, \quad \phi_2 = \begin{pmatrix} H_2^+ \\ \frac{H_2^0 + iA_2^0}{\sqrt{2}} \end{pmatrix}, \quad \phi_3 = \begin{pmatrix} G^+ \\ \frac{v+h+iG^0}{\sqrt{2}} \end{pmatrix}$$

- $\phi_3$  – SM-like doublet with SM-like Higgs  $h$
- $Z_2$ -odd doublets  $\phi_1$  and  $\phi_2$  mix:

$$S_1 = \frac{\alpha H_1^0 + \alpha H_2^0 - A_1^0 + A_2^0}{\sqrt{2\alpha^2 + 2}}, \quad S_2 = \frac{-H_1^0 - H_2^0 - \alpha A_1^0 + \alpha A_2^0}{\sqrt{2\alpha^2 + 2}}$$

$$S_3 = \frac{\beta H_1^0 - \beta H_2^0 + A_1^0 + A_2^0}{\sqrt{2\beta^2 + 2}}, \quad S_4 = \frac{-H_1^0 + H_2^0 + \beta A_1^0 + \beta A_2^0}{\sqrt{2\beta^2 + 2}}$$

$$H_1^\pm = \frac{e^{\pm i\theta_{12}/2}}{\sqrt{2}} (S_1^\pm - S_2^\pm), \quad H_2^\pm = \frac{e^{\mp i\theta_{12}/2}}{\sqrt{2}} (S_1^\pm + S_2^\pm)$$

- 4 neutral and 4 charged  $Z_2$ -odd particles (double the IDM)
- $S_1$  – DM candidate, other dark particles heavier

## Physical Parameters

Parameters of  $V$ :  $\mu_2^2, |\lambda_2|, |\mu_{12}^2|, \lambda_{23}, \lambda'_{23}, \theta_{12}, \theta_2$

**Physical parameters:**

**DM mass:**

$$m_{S_1}$$

**Mass splittings:**

$$\delta_{12} = m_{S_2} - m_{S_1}$$

$$\delta_{1c} = m_{S_1^\pm} - m_{S_1}$$

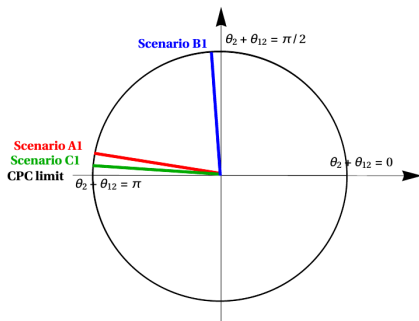
$$\delta_c = m_{S_2^\pm} - m_{S_1^\pm}$$

**Higgs-DM coupling:**

$$g_{S_1 S_1 h}$$

**CPv phases:**

$$\theta_{12}, \theta_2$$





## Benchmark scenarios

This talk:  $m_{S_1} < m_Z$ :

$$A1 : \delta_{12} = 125 \text{ GeV}, \delta_{1c} = 50 \text{ GeV}, \delta_c = 50 \text{ GeV}, \theta_2 = \theta_{12} = 1.5$$

$$m_{S_1} < m_{S_{2,3,4}}, m_{S_{1,2}^\pm} \text{ (no coannihilation)}$$

$$B1 : \delta_{12} = 125 \text{ GeV}, \delta_{1c} = 50 \text{ GeV}, \delta_c = 50 \text{ GeV}, \theta_2 = \theta_{12} = 0.82$$

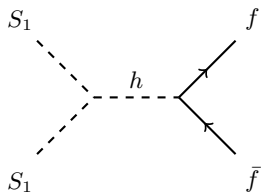
$$m_{S_1} \approx m_{S_3} < m_{S_{2,4}}, m_{S_{1,2}^\pm}$$

$$C1 : \delta_{12} = 12 \text{ GeV}, \delta_{1c} = 100 \text{ GeV}, \delta_c = 1 \text{ GeV}, \theta_2 = \theta_{12} = 1.57$$

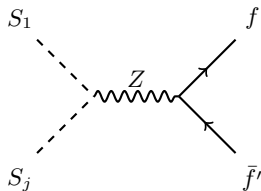
$$m_{S_1} \approx m_{S_3} \approx m_{S_4} \approx m_{S_2} < m_{S_{1,2}^\pm}$$

*Checked against experimental and theoretical constraints:  
details in JHEP 1612 (2016) 014*

## DM Annihilation for light DM

**Higgs-mediated annihilation**

depends on  $m_{S_1}$  and  $g_{S_1 S_1 h}$

**Z-mediated coannihilation**

depends on  $m_{S_j} - m_{S_1}$ :

**A1**: no coannihilation

**B1**:  $S_1 S_3$  coannihilation only

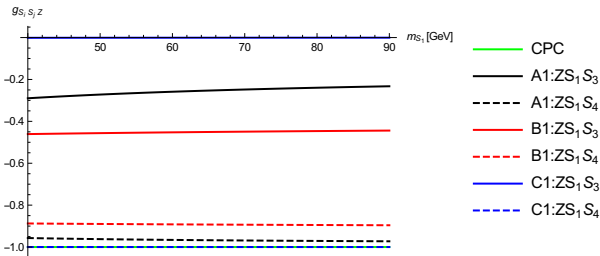
**C1**:  $S_1 S_3, S_2 S_4, S_1 S_4, S_2 S_3$  coann.

depends on  $Z S_i S_j$  couplings

## Z-inert couplings

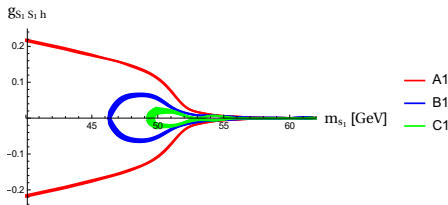
$$\chi_{ZS_1S_3} = \chi_{ZS_2S_4} = \frac{\alpha+\beta}{\sqrt{\alpha^2+1}\sqrt{\beta^2+1}}, \quad \chi_{ZS_1S_4} = \chi_{ZS_2S_3} = \frac{\alpha\beta-1}{\sqrt{\alpha^2+1}\sqrt{\beta^2+1}},$$

$$\chi_{ZS_1S_3}^2 + \chi_{ZS_1S_4}^2 = 1, \quad \chi_{ZS_2S_3}^2 + \chi_{ZS_2S_4}^2 = 1$$



- **mass order:**  $m_{S_1} < m_{S_3} < m_{S_4} < m_{S_2}$
- CPC value  $\chi_{ZS_1S_3} = -1, \chi_{ZS_1S_4} = 0$
- $ZS_1S_3$  – **reduced**; 20 – 50% for **A1**, **B1**,  $\sim 0$  for **C1**
- $ZS_1S_4$  – close to the CPC value  $\rightarrow$  dominant channel for **C1**

## Low DM mass



**A1:** mainly Higgs annihilation, large  $g_{S_1 S_1 h}$

**B1:** Higgs annihilation (smaller  $g_{S_1 S_1 h}$ )

+  $ZS_1 S_3$  coannihilation (reduced with respect to the CPc case)

**C1:** mainly  $ZS_1 S_4$  coannihilation ( $\chi_{ZS_1 S_4} \approx -1$ )

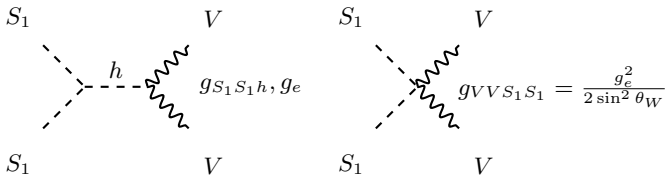
+ Higgs annihilation

Tools used in calculation: LanHEP, arXiv:1412.5016 [physics.comp-ph]; CalcHEP 3.4, Comput. Phys.

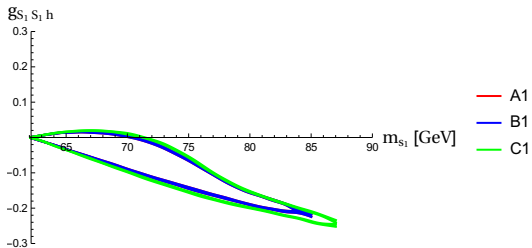
Commun. **184** (2013) 1729; micrOMEGAs 4.2 arXiv:1407.6129 [hep-ph]

## Medium DM mass

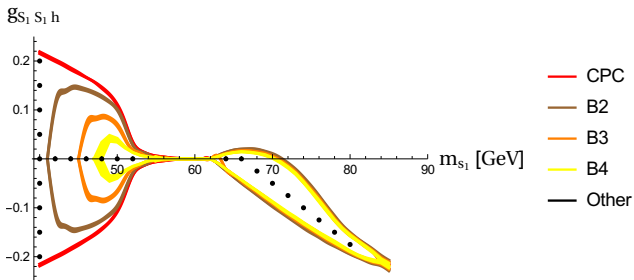
Main annihilation channels into gauge bosons  $V = W^\pm, Z$ :



no dependence on the benchmarks



## Filling the plot



$m_{S_1} < m_h/2$ : many new solutions:

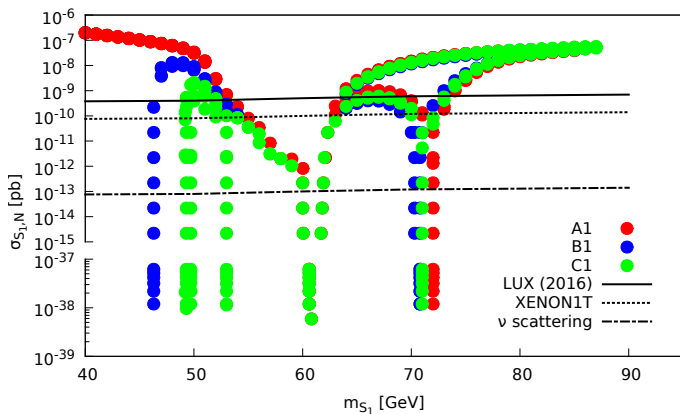
different mass splittings +  $Z S_i S_j$  interaction strength

$m_{S_1} > m_h/2$ : less freedom but still new solutions:

Higgs mediated coannihilation + sign of  $h S_3 S_3$  coupling

## Direct detection

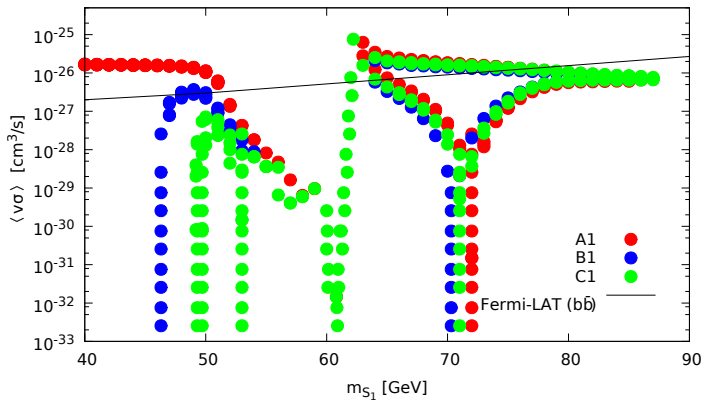
$$\sigma_{S_1, N} \propto \frac{g_{S_1 S_1 h}^2}{(m_{S_1} + m_N)^2}$$



Case **A1**: mostly excluded (large  $g_{S_1 S_1 h}$ )

Cases **B1** and **C1**: mostly within the limits

## Indirect detection



Most of the parameters space in agreement with Fermi-LAT



## LHC constraints

### Higgs invisible decays

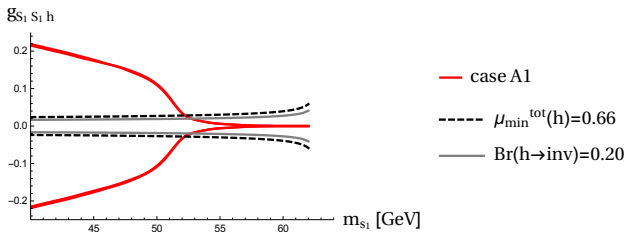
$$\Gamma(h \rightarrow S_1 S_1) = \frac{g_{S_1 S_1 h}^2}{32\pi m_h} \left(1 - \frac{4m_{S_1}^2}{m_h^2}\right)^{1/2}, \quad \text{Br}(h \rightarrow \text{inv}) = \frac{\Gamma(h \rightarrow S_1 S_1)}{\Gamma_h^{\text{SM}} + \Gamma(h \rightarrow S_1 S_1)}$$

### Higgs total decay

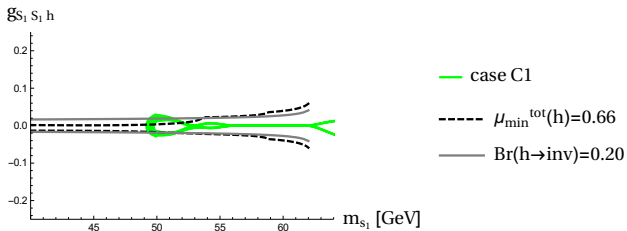
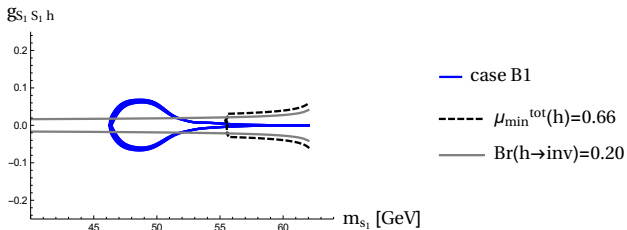
$$\mu_{\text{tot}} = \frac{\text{BR}(h \rightarrow XX)}{\text{BR}(h_{\text{SM}} \rightarrow XX)} = \frac{\Gamma_{\text{tot}}^{\text{SM}}(h)}{\Gamma_{\text{tot}}^{\text{SM}}(h) + \Gamma^{\text{inert}}(h)}$$

### LHC limits

$$\mu_{\text{tot}} = 1.17 \pm 0.17 \quad \text{and} \quad \text{Br}(h \rightarrow \text{inv}) < 0.2$$



## LHC constraints



In general: scenarios of type **C** are the least constrained.

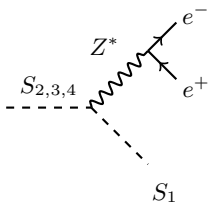
## Conclusions and Outlook

- **3HDM** with  $Z_2$  symmetry: I(2+1)HDM
- viable DM candidate
- large dark sector: **important coannihilation effects in  $\Omega_{DM}h^2$** 
  - **varying strength of gauge-inert couplings**
  - **new regions in agreement with Planck**
- agreement with direct and indirect detection limits:  
 $45 \text{ GeV} \lesssim m_{S_1} \lesssim 62.5 \text{ GeV}$ ,  $64 \text{ GeV} \lesssim m_{S_1} \lesssim 74 \text{ GeV}$ ,  $m_{S_1} \gtrsim 400 \text{ GeV}$ 
  - **as long as DM is practically invisible**
  - other detections prospects?

## Inert cascade decays at the LHC

*work in progress*

$$pp \rightarrow Z \rightarrow S_{2,3,4} S_1 \rightarrow S_1 S_1 Z^* \rightarrow S_1 S_1 e^+ e^-$$

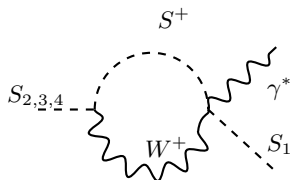
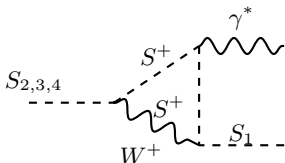


- signature: **missing  $E_T$  and dilepton pair**
- dominant if there is a **large mass splitting** between DM and other inert particles
- process present in the IDM (through  $HAZ$  vertex)
- note: **possible differences with respect to the IDM due to varying strength of  $S_1 S_j Z$  vertex**

## Inert cascade decays at the LHC

*work in progress*

$$pp \rightarrow h \rightarrow S_{2,3,4} S_1 \rightarrow S_1 S_1 \gamma^* \rightarrow S_1 S_1 e^+ e^-$$

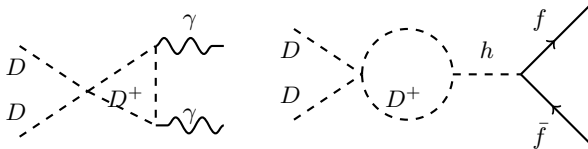


- signature: **missing  $E_T$  and dilepton pair**
- important if there is a **small mass splitting** between DM and other inert particles  $\rightarrow$  scenario C is preferred anyway
- process absent in the IDM (no  $A \rightarrow H\gamma^*$  loop)
- promising preliminary results  $\sigma \sim 10^{-5}$  pb

## DM self-couplings

Dark self-couplings – no impact on standard DM and LHC phenomenology

- In the I(1+1)HDM –  $\lambda_2$
- In the I(2+1)HDM –  $\lambda_1, \lambda_{11,22,12}, \lambda'_{12}$
- Possible relevant corrections to loop processes, e.g.:



$D$  – DM matter,  $D^\pm$  – charged dark scalar

- Astrophysical DM detection experiments:  
DM very weakly coupled to the visible sector  
 $\Rightarrow$  **loop corrections can be important!**

## Final Summary

Scalar DM models:

- interesting phenomenology
- strong constraints on  $g_{DMh}$  from DM experiments
- need to move away from Higgs-portal
- solution: rich particle spectrum & coannihilation effects
- interesting LHC signatures – a tool for testing DM models?
- loop processes & role of self-couplings – can be important
- **Further work needed!**

## References

- Higgs-portal DM models

[B. Patt and F. Wilczek, hep-ph/0605188, X. Chu, T. Hambye, and M. H. Tytgat, JCAP 1205 (2012) 034, A. Djouadi, O. Lebedev, Y. Mambrini, and J. Quevillon, Phys.Lett. B709 (2012) 65–69]

- 3HDM

[V. Keus, S. King, S. Moretti JHEP 1401 (2014) 052, arXiv:1408.0796; V. Keus, S. F. King, S. Moretti and D. Sokolowska, JHEP 1411 (2014) 016, JHEP 1511, 003 (2015)]

- Experimental constraints

[ATLAS and CMS collaborations, JHEP 08 (2016) 045; [http://lux.brown.edu/LUX\\_dark\\_matter/Talks\\_files/LUX\\_NewDarkMatterSearchResult\\_332LiveDays\\_IDM2016\\_160721.pdf](http://lux.brown.edu/LUX_dark_matter/Talks_files/LUX_NewDarkMatterSearchResult_332LiveDays_IDM2016_160721.pdf) (“Dark-matter results from 332 new live days of LUX data, Identification of Dark Matter, The University of Sheffield, Sheffield, UK, 21 July, 2016”), M. Ackermann *et al.* [Fermi-LAT Collaboration], Phys. Rev. Lett. 115 (2015) 23, 231301, XENON1T Collaboration, Springer Proc. Phys. 148 (2013) 93 ]

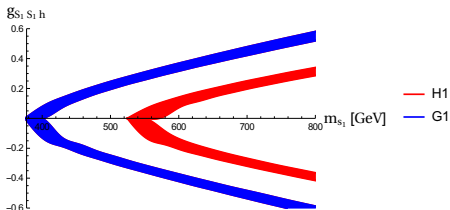
- Numerical Tools

[LanHEP, arXiv:1412.5016 [physics.comp-ph]; CalcHEP 3.4, Comput. Phys. Commun. 184 (2013) 1729; micrOMEGAs 4.2 arXiv:1407.6129 [hep-ph]]



# BACKUP SLIDES

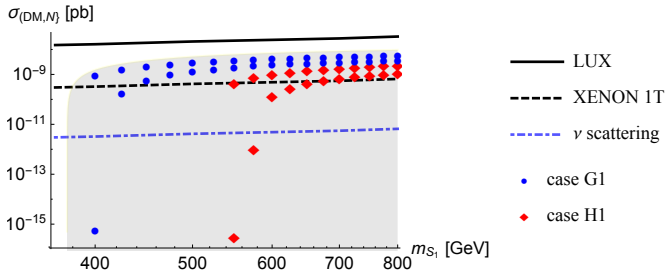
## Heavy DM



- $g_{S_1 S_1 h}$  in Case G  $>$   $g_{S_1 S_1 h}$  Case H
- The same behaviour in both cases
- Lower  $m_{S_1}$  for Case G
- Not really different from the CPc case

## Direct Detection

$$\sigma_{DM,N} \propto \frac{g_{S_1 S_1 h}^2}{(M_{S_1} + M_N)^2}$$



- in agreement with LUX
- within the reach of XENON-1T
- case G (bigger couplings) easier to see/exclude than case H (smaller couplings)

## Mass formulas

$$m_{S_1^\pm}^2 = (-\mu_2^2 - |\mu_{12}^2|) + \frac{1}{2}\lambda_{23}v^2, \quad m_{S_2^\pm}^2 = (-\mu_2^2 + |\mu_{12}^2|) + \frac{1}{2}\lambda_{23}v^2.$$

$$m_{S_1}^2 = \frac{v^2}{2}(\lambda'_{23} + \lambda_{23}) - \Lambda - \mu_2^2,$$

$$m_{S_2}^2 = \frac{v^2}{2}(\lambda'_{23} + \lambda_{23}) + \Lambda - \mu_2^2,$$

$$m_{S_3}^2 = \frac{v^2}{2}(\lambda'_{23} + \lambda_{23}) - \Lambda' - \mu_2^2,$$

$$m_{S_4}^2 = \frac{v^2}{2}(\lambda'_{23} + \lambda_{23}) + \Lambda' - \mu_2^2,$$

$$\Lambda = \sqrt{v^4|\lambda_2|^2 + |\mu_{12}^2|^2 - 2v^2|\lambda_2||\mu_{12}^2|\cos(\theta_{12} + \theta_2)},$$

$$\Lambda' = \sqrt{v^4|\lambda_2|^2 + |\mu_{12}^2|^2 + 2v^2|\lambda_2||\mu_{12}^2|\cos(\theta_{12} + \theta_2)}.$$

$$\alpha = \frac{-|\mu_{12}^2|\cos\theta_{12} + v^2|\lambda_2|\cos\theta_2 - \Lambda}{|\mu_{12}^2|\sin\theta_{12} + v^2|\lambda_2|\sin\theta_2}, \quad \beta = \frac{|\mu_{12}^2|\cos\theta_{12} + v^2|\lambda_2|\cos\theta_2 - \Lambda'}{|\mu_{12}^2|\sin\theta_{12} - v^2|\lambda_2|\sin\theta_2}.$$

## Physical Basis

$$|\mu_{12}^2| = \frac{1}{2}(m_{S_2^\pm}^2 - m_{S_1^\pm}^2),$$

$$\lambda_{23} = \frac{2\mu_2^2}{v^2} + \frac{m_{S_2^\pm}^2 + m_{S_1^\pm}^2}{v^2},$$

$$\lambda'_{23} = \frac{1}{v^2}(m_{S_2}^2 + m_{S_1}^2 - m_{S_2^\pm}^2 - m_{S_1^\pm}^2),$$

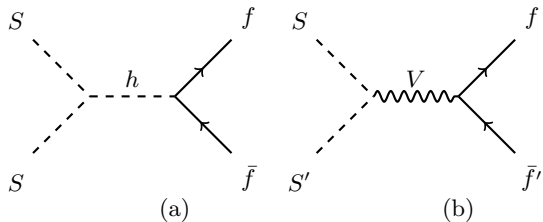
$$\mu_2^2 = \frac{v^2}{2} g_{S_1 S_1 h} - \frac{v^2 |\lambda_2|}{2(1 + \alpha^2)} \left( 4\alpha \sin \theta_2 + 2(\alpha^2 - 1) \cos \theta_2 \right) - \frac{m_{S_2}^2 + m_{S_1}^2}{2},$$

$$|\lambda_2| = \frac{1}{v^2} [|\mu_{12}^2| \cos(\theta_2 + \theta_{12}) +$$

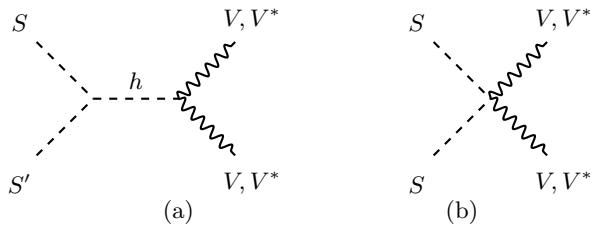
$$\sqrt{|\mu_{12}^2|^2 \cos^2(\theta_2 + \theta_{12}) + \left( \frac{m_{S_2}^2 - m_{S_1}^2}{2} \right)^2 - |\mu_{12}^2|^2}].$$

## DM annihilation diagrams

Light DM annihilation:

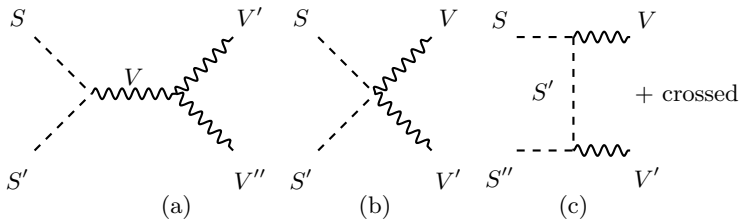


Virtual gauge bosons:



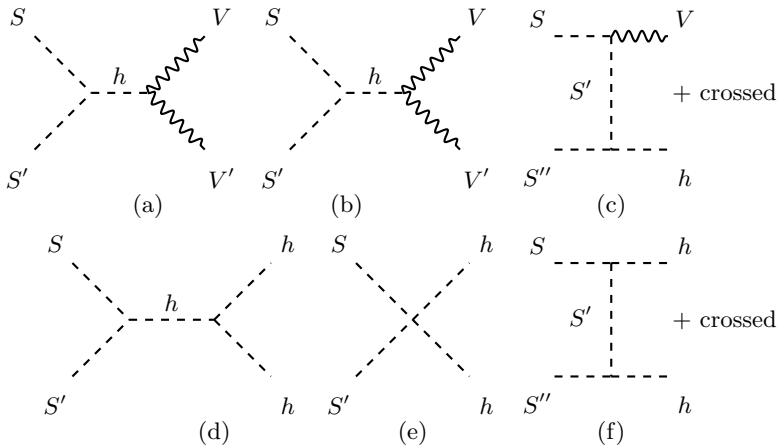
## DM annihilation diagrams - gauge limit

Heavy DM (co)annihilation diagrams with pure gauge boson final states:



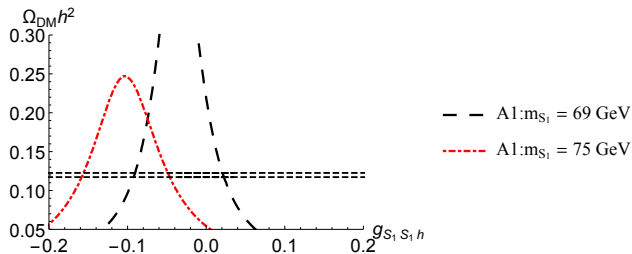
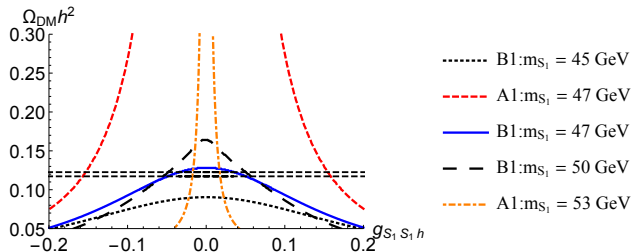
## DM annihilation diagrams

Heavy DM (co)annihilation channels involving the SM-like Higgs boson:





# Relic density



# Higgs-inert couplings

