

NNLO QCD predictions for single jet inclusive production at the LHC

João Pires (MPI Munich)

Based on work in [Phys. Rev. Lett. 118, 072002 (2017)] and new results

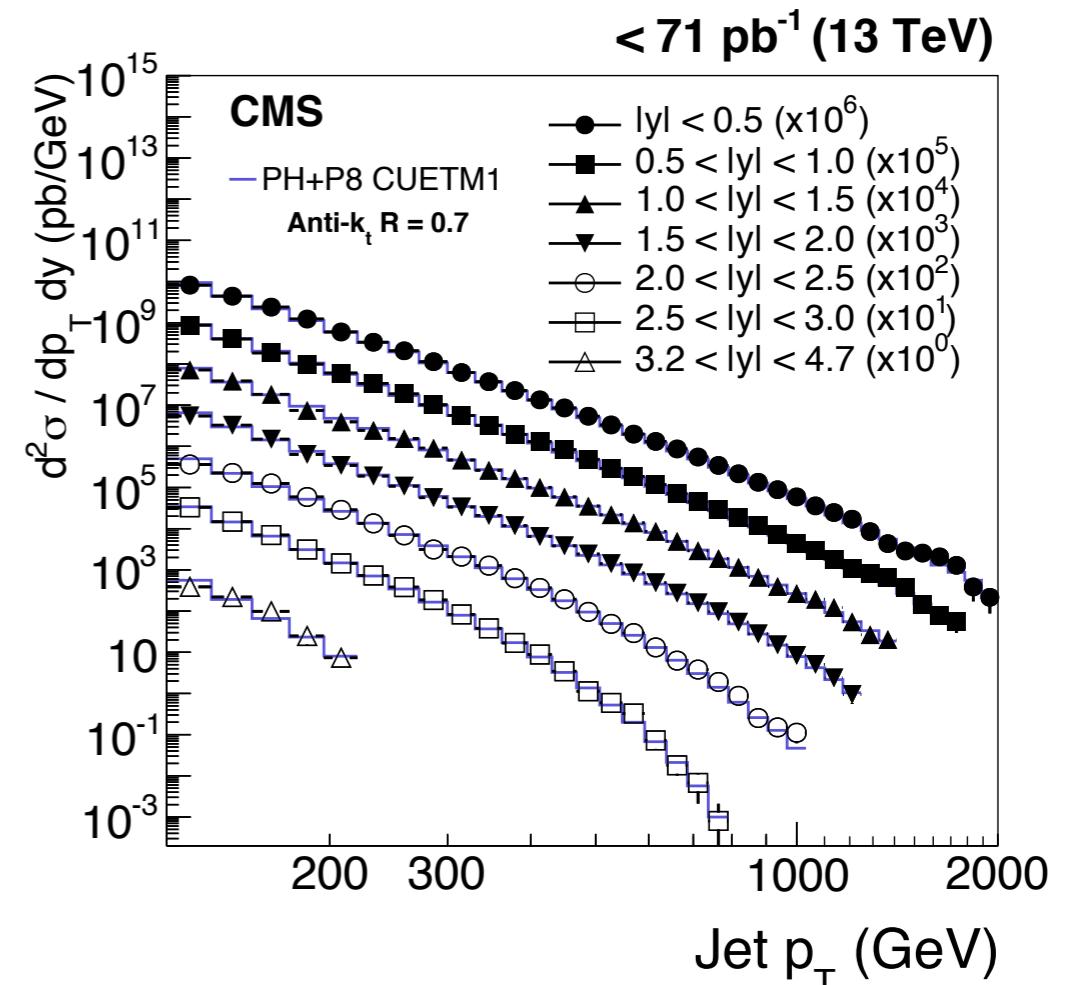
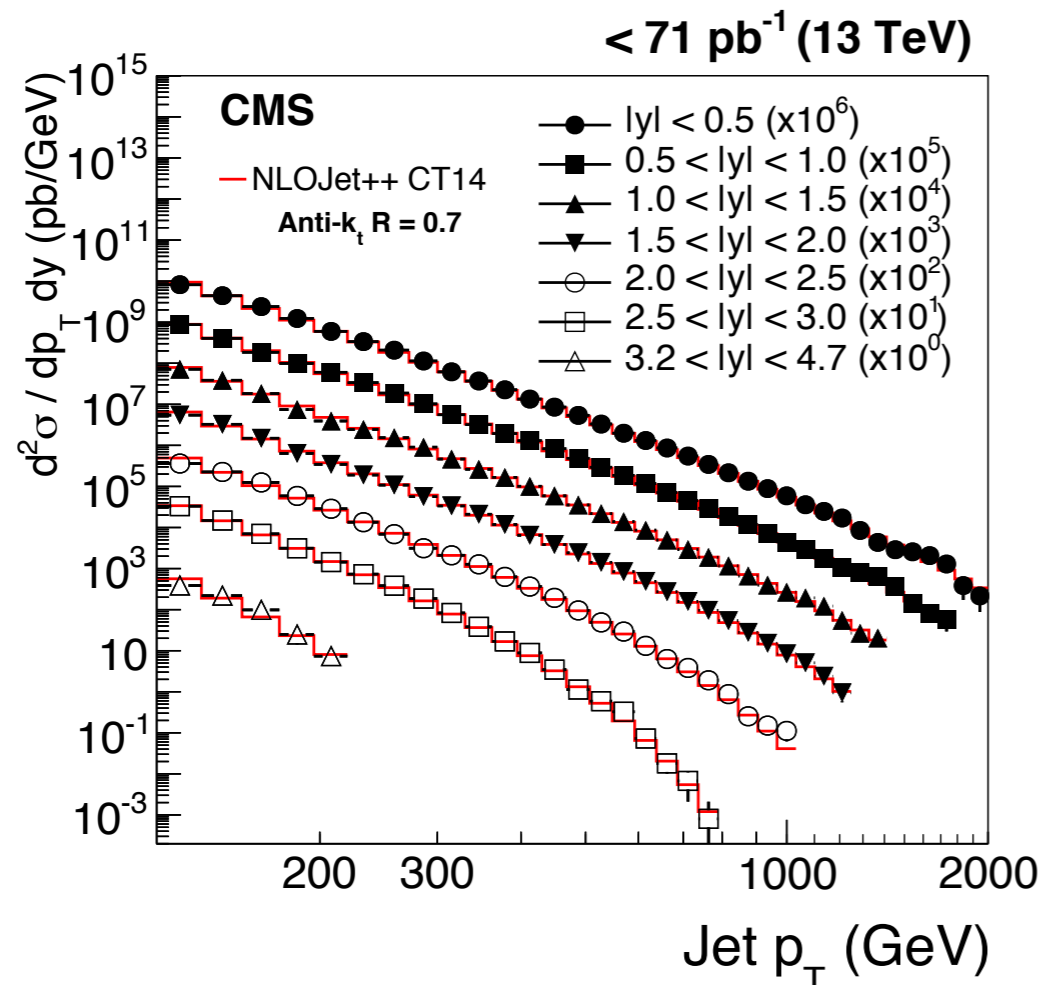
Rencontres de Moriond QCD Session March 30, 2017



Jets at the LHC

- look at production of jets of hadrons with large transverse energy
- for sufficiently high transverse momentum $p_T > 20$ GeV high rates and clean and simple cross section definition

$$\frac{d\sigma}{dp_T dy} = \frac{1}{\mathcal{L}} \frac{N_{jets}}{\Delta p_T \Delta y}$$



Jets at the LHC

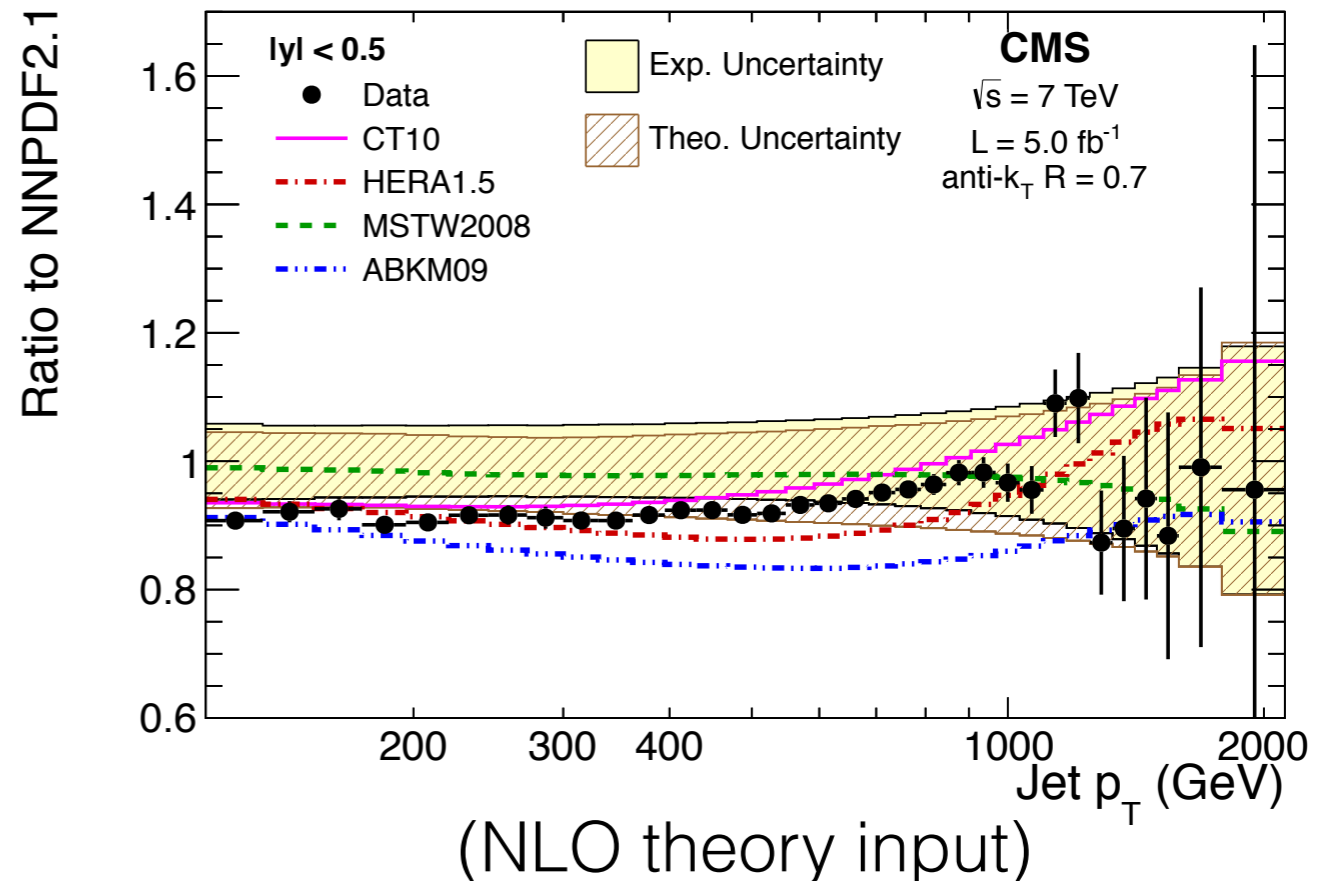
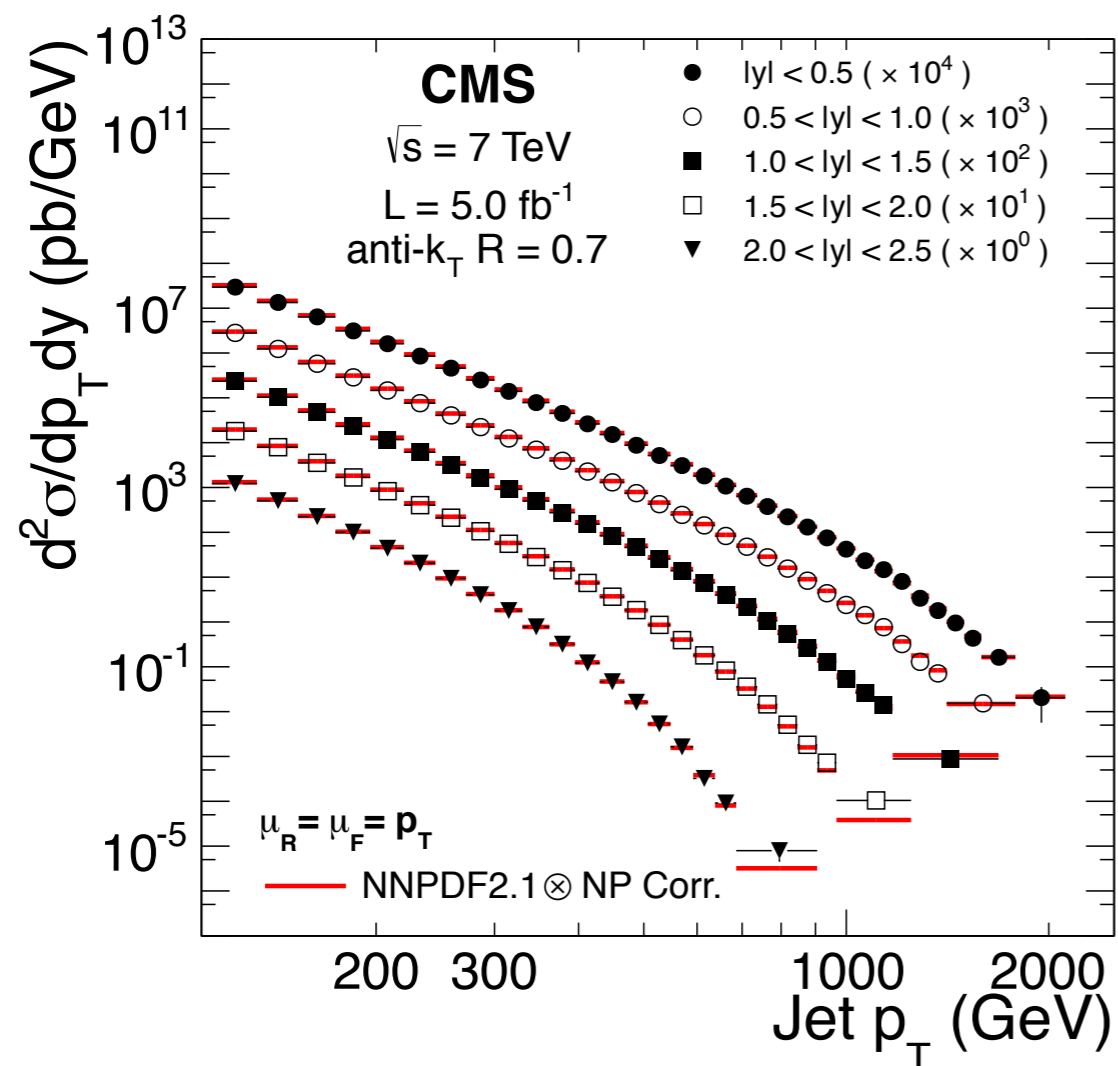
$$\sigma(P_1, P_2) = \sum_{i,j} \int dx_1 dx_2 f_i(x_1, \mu_F^2) f_j(x_2, \mu_F^2) \hat{\sigma}_{ij}(p_1, p_2, \alpha_s(\mu^2), s/\mu^2, s/\mu_F^2)$$

- Inputs from physical parameters
 - PDF's ; strong coupling
- Unphysical inputs
 - renormalization and factorization scales
- Partonic cross section calculable in perturbation theory

$$\hat{\sigma}_{ij} = \hat{\sigma}_{ij}^{LO} + \left(\frac{\alpha_s}{2\pi}\right) \hat{\sigma}_{ij}^{NLO} + \left(\frac{\alpha_s}{2\pi}\right)^2 \hat{\sigma}_{ij}^{NNLO} + \mathcal{O}(\alpha_s^3)$$

Jets at the LHC

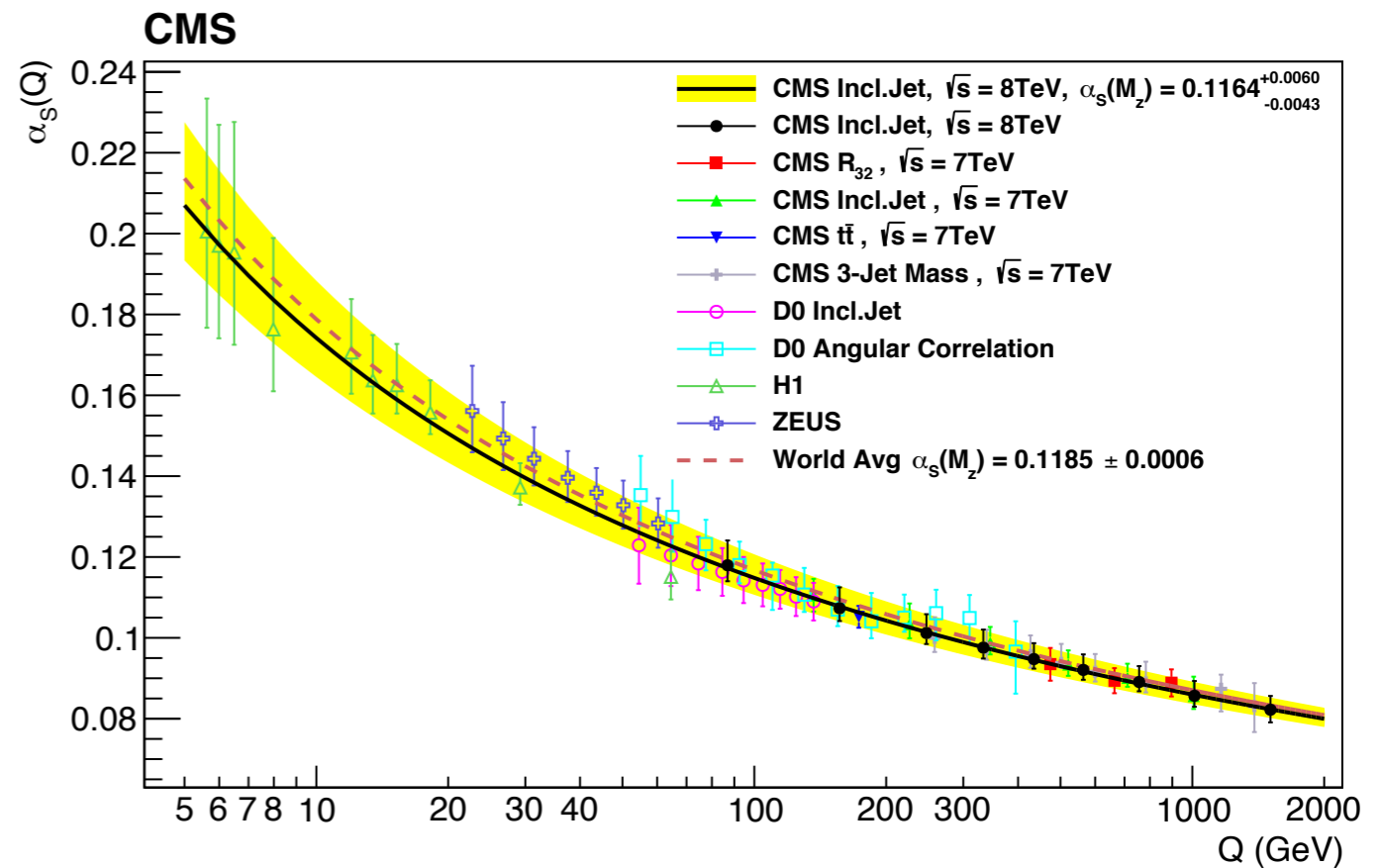
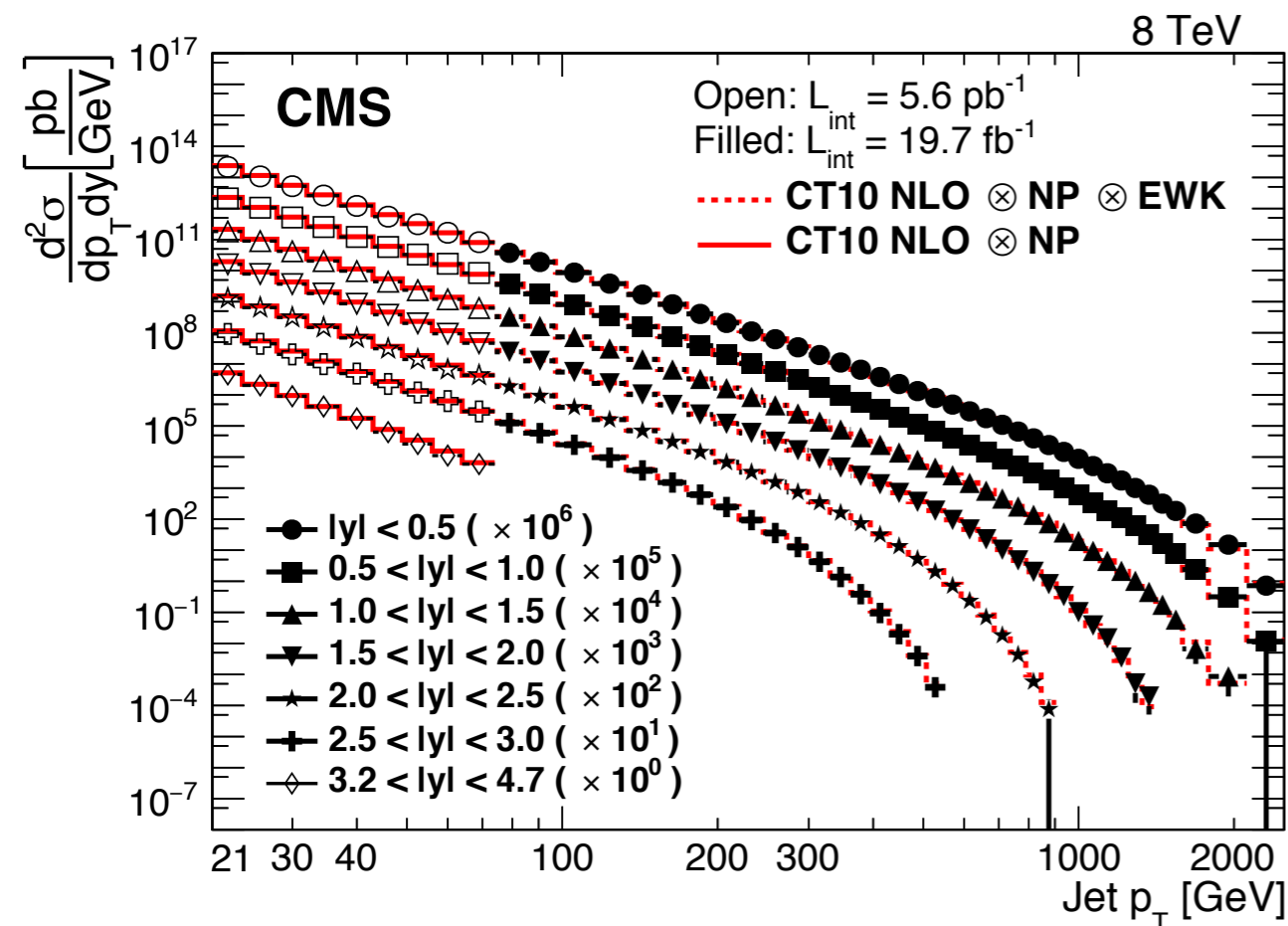
- Percent level experimental accuracy in jet production is a reality at the LHC
 - Rigorous tests of pQCD dynamics across a huge range of kinematics
 - Constrain PDF's



Jets at the LHC

- Percent level experimental accuracy in jet production is a reality at the LHC
 - Rigorous tests of pQCD dynamics across a huge range of kinematics
 - Constrain PDF's
 - Determine $\alpha_s(M_Z)$ and running coupling from a single experiment

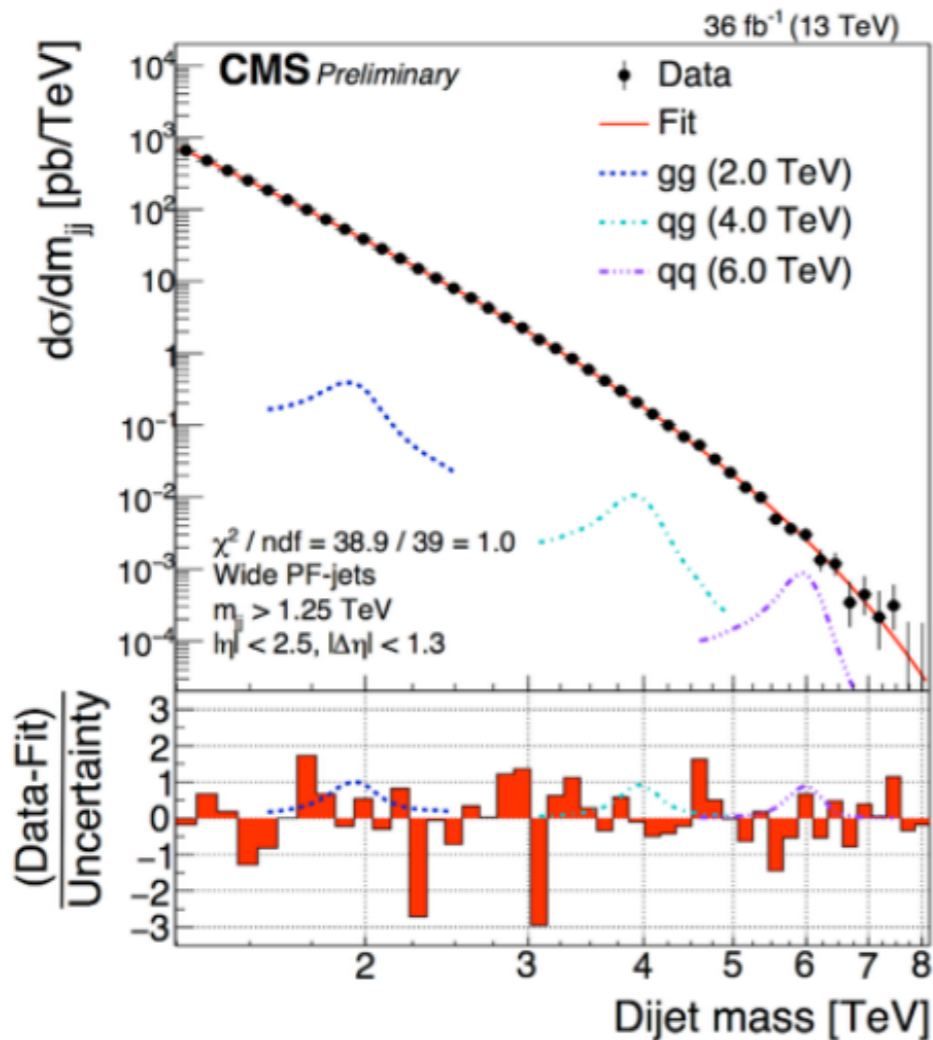
CMS-SMP-14-001



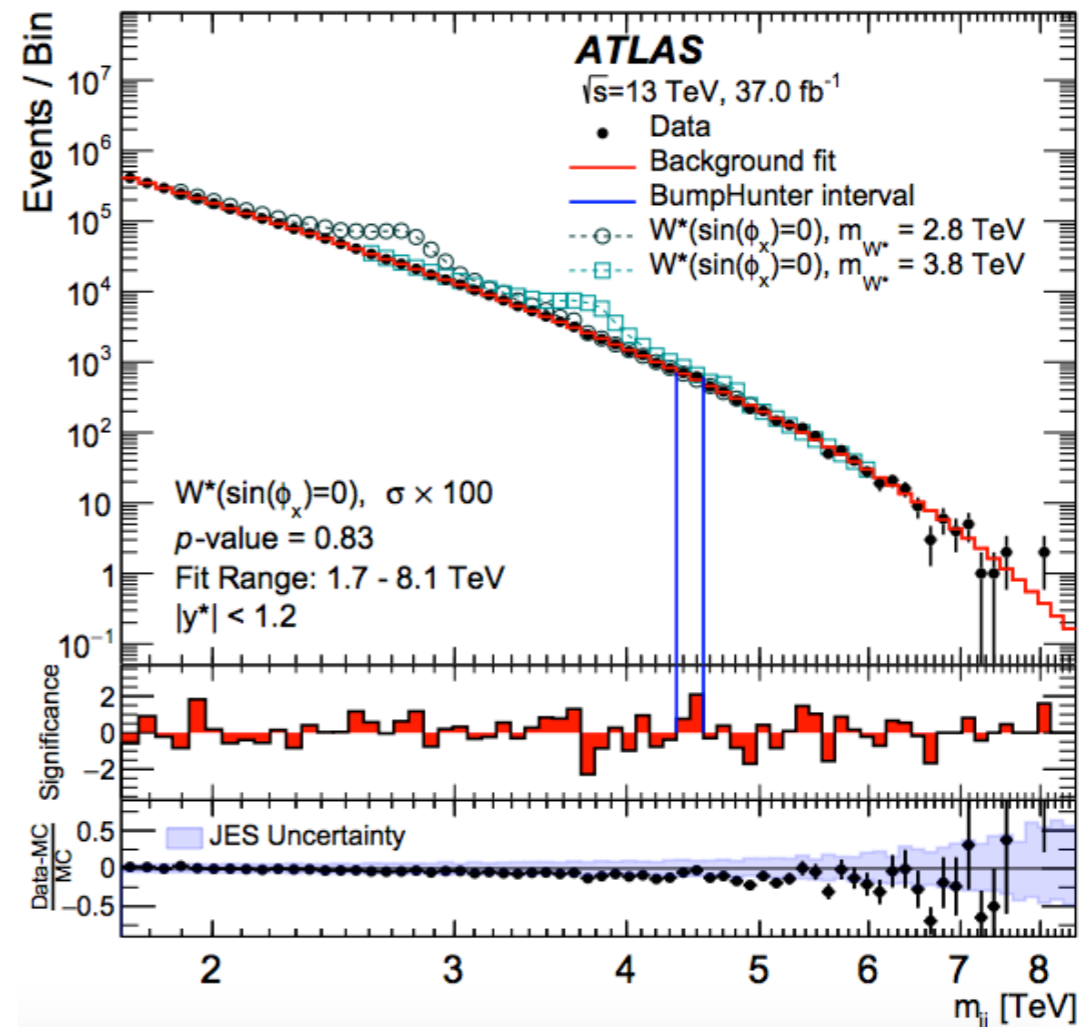
(NLO theory input)

Jets at the LHC

- Percent level experimental accuracy in jet production is a reality at the LHC
 - Rigorous tests of pQCD dynamics across a huge range of kinematics
 - Constrain PDF's
 - Background to high mass resonances decaying to dijet final states



arXiv:1703.09127



Theory state of the art

- NLO QCD [Ellis, Kunszt, Soper '92] [Giele, Glover, Kosower '94] [Nagy 02]
- NLO QCD + PS (POWHEG) [Alioli, Hamilton, Nason, Oleari, Re '11]
- NLO EW [Dittmaier, Huss, Speckner '13]
[Frederix, Frixione, Hirschi, Pagani, Shao, Zaro '16]
- NNLO QCD [Gehrmann-De Ridder, Gehrmann, Glover, JP '13]
[Currie, Gehrmann-De Ridder, Glover, JP '13]
[Currie, Glover, JP '16]

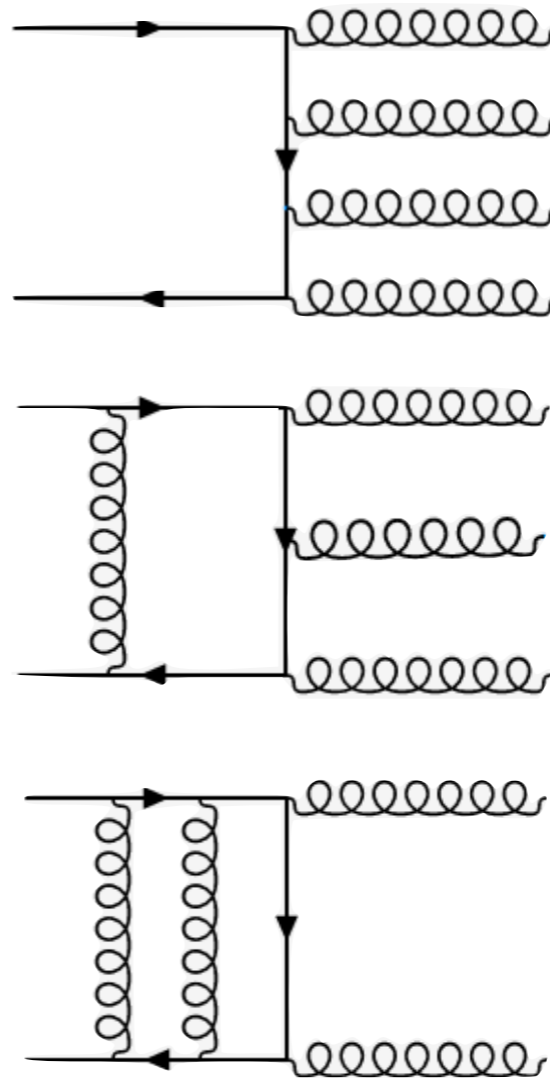
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New approach to look at jet data at the LHC (this talk)

Anatomy of an NNLO calculation

$$\begin{aligned}
 d\hat{\sigma}_{NNLO} &= \int_{d\Phi_4} d\hat{\sigma}_{NNLO}^{RR} \\
 &+ \int_{d\Phi_3} d\hat{\sigma}_{NNLO}^{RV} \\
 &+ \int_{d\Phi_2} d\hat{\sigma}_{NNLO}^{VV}
 \end{aligned}$$



- double-unresolved
- single-unresolved
- single-unresolved
- $1/\epsilon^2$; $1/\epsilon$
- $1/\epsilon^4$; $1/\epsilon^3$;
 $1/\epsilon^2$; $1/\epsilon$;

- 6 parton tree level [Berends, Giele '87] [Mangano, Parke, Xu '87] [Britto, Cachazo, Feng '06]
- 5 parton one-loop [Bern, Dixon, Kosower '93]
- 4 parton two-loop [Anastasiou, Glover, Oleari, Tejeda-Yeomans '01][Bern, De Freitas Dixon '02]
- non-trivial cancellation of infrared singularities at NNLO

NNLO antenna subtraction

$$\begin{aligned}
 d\hat{\sigma}_{NNLO} &= \int_{d\Phi_4} \left(d\hat{\sigma}_{NNLO}^{RR} - d\hat{\sigma}_{NNLO}^S \right) \\
 &+ \int_{d\Phi_3} \left(d\hat{\sigma}_{NNLO}^{RV} - d\hat{\sigma}_{NNLO}^T \right) \\
 &+ \int_{d\Phi_2} \left(d\hat{\sigma}_{NNLO}^{VV} - d\hat{\sigma}_{NNLO}^U \right)
 \end{aligned}$$

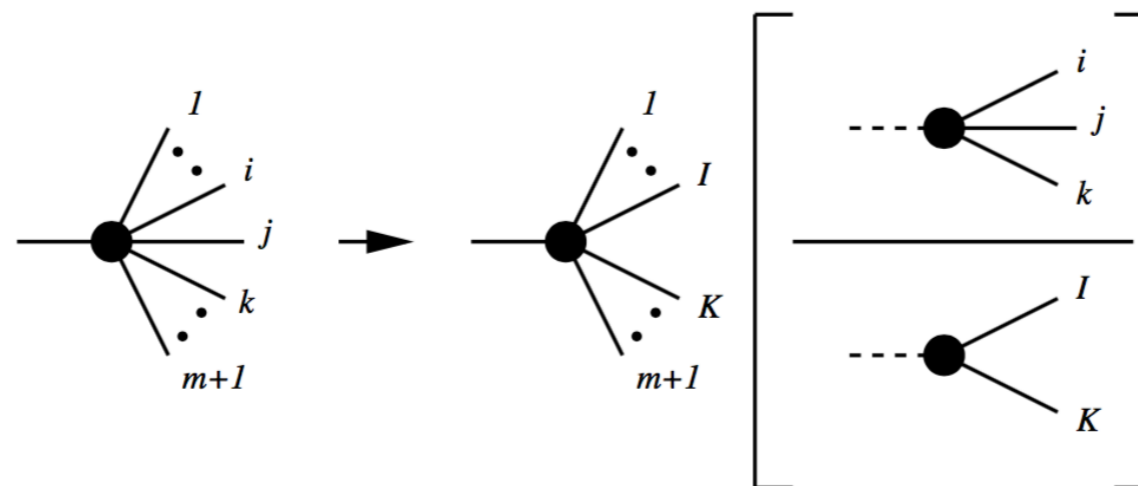
$$d\hat{\sigma}_{NNLO}^S \quad d\hat{\sigma}_{NNLO}^T$$

- mimic RR,RV in unresolved limits

$$d\hat{\sigma}_{NNLO}^T \quad d\hat{\sigma}_{NNLO}^U$$

- analytically cancel the poles in RV and VV matrix elements

- universal factorization properties in IR limits



- phase space factorization

$$d\Phi_{m+1}(p_1, \dots, p_{m+1}; q) = d\Phi_m(p_1, \dots, \tilde{p}_I, \tilde{p}_K, \dots, p_{m+1}; q) \cdot d\Phi_{X_{ijk}}(p_i, p_j, p_k; \tilde{p}_I + \tilde{p}_K)$$

NNLO antenna subtraction

$$\begin{aligned}
 d\hat{\sigma}_{NNLO} &= \int_{d\Phi_4} \left(d\hat{\sigma}_{NNLO}^{RR} - d\hat{\sigma}_{NNLO}^S \right) \\
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$$d\hat{\sigma}_{NNLO}^S \quad d\hat{\sigma}_{NNLO}^T$$

- mimic RR,RV in unresolved limits

$$d\hat{\sigma}_{NNLO}^T \quad d\hat{\sigma}_{NNLO}^U$$

- analytically cancel the poles in RV and VV matrix elements

For pp->jet +X

- all subprocesses included through leading colour $\alpha_s^2 N^2, \alpha_s^2 N N_F, \alpha_s^2 N_F^2$

$$\{gg, qq, q\bar{q}, qq', q\bar{q}'\}$$

- subleading colour contributions to the NNLO coefficient neglected

NNLOJET parton level generator

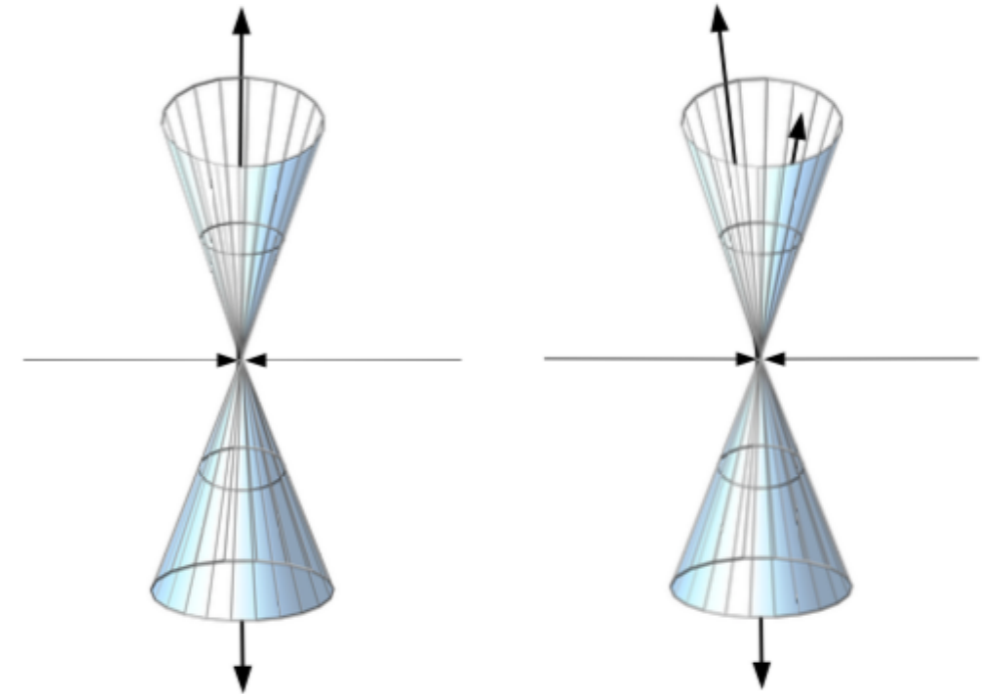
[X. Chen, J. Cruz-Martinez, J. Currie, A. Gehrmann-De Ridder, T. Gehrmann, N. Glover, A. Huss, M. Jaquier, T. Morgan, J. Niehues, JP]

- parton level generator based on antenna subtraction to compute fully differential cross sections at NNLO in QCD
 - $pp \rightarrow H + \text{jet}$
 - $pp \rightarrow Z + \text{jet}$
 - $pp \rightarrow 2 \text{ jets}$
 - $pp \rightarrow H, W, Z$
 - $ep \rightarrow 2 \text{ jets}$ \longrightarrow Jan Niehues talk
 - interface to applfast nnlo tables \longrightarrow Daniel Britzger's talk

Single jet inclusive scale choice

two widely used scale choices:

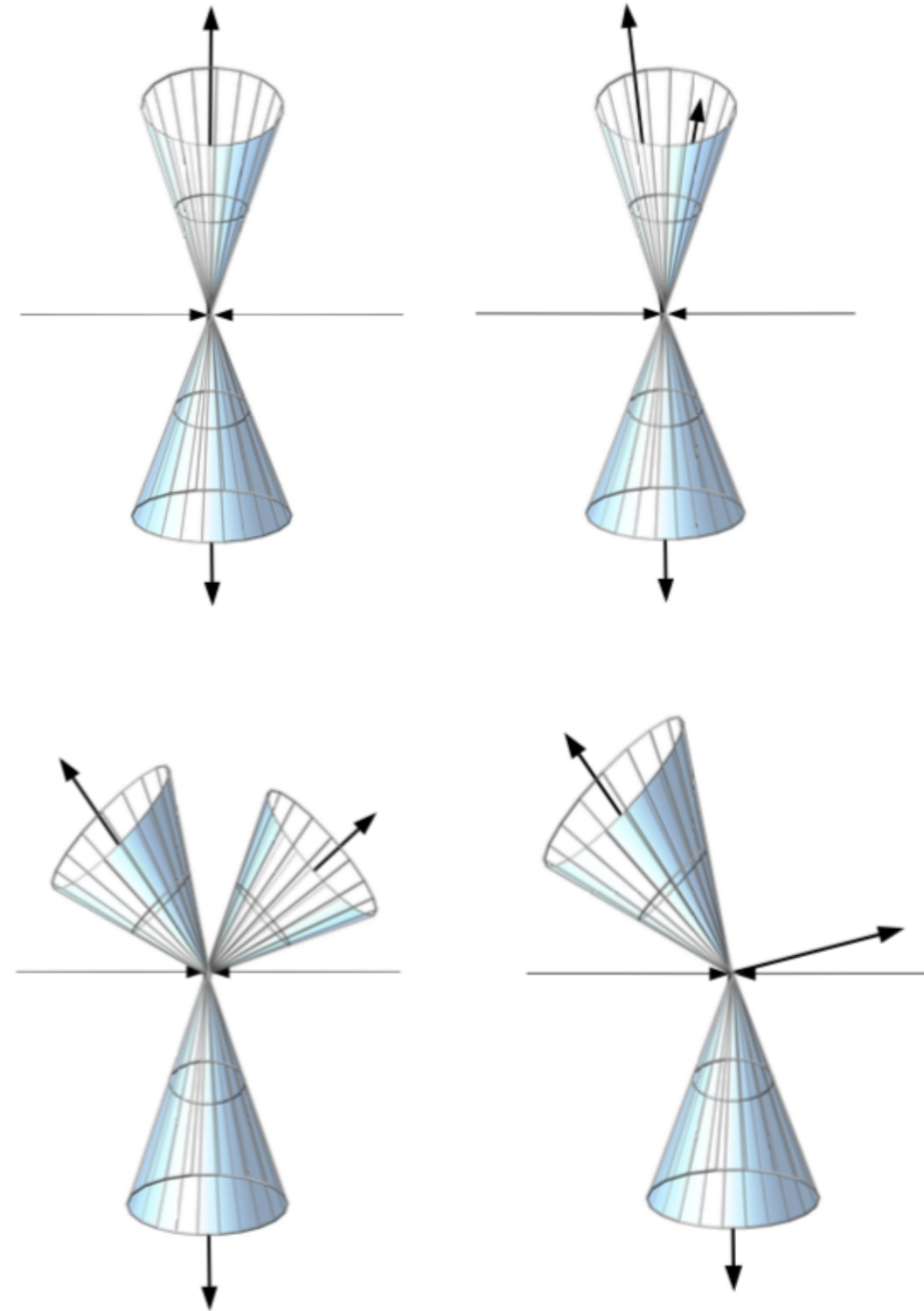
- $\mu_R = \mu_F = \{p_{T1}, p_T\}$
 - leading jet p_T in the event p_{T1}
 - individual jet p_T
- high p_T jets are back to back $\Rightarrow p_T \rightarrow p_{T1}$



Single jet inclusive scale choice

two widely used scale choices:

- $\mu_R = \mu_F = \{p_{T1}, p_T\}$
 - leading jet p_T in the event p_{T1}
 - individual jet p_T
 - high p_T jets are back to back $\Rightarrow p_T \rightarrow p_{T1}$
 - $p_T \neq p_{T1}$ for:
 - 3jet events
 - 3rd jet outside fiducial jet cuts
- \Rightarrow with p_T choice the real emission event with different R gives rise to a different scale \Rightarrow larger $R \Rightarrow$ harder scale $\Rightarrow p_T \rightarrow p_{T1}$
- at NLO the p_{T1} scale choice generates the same hard scale for the event independent of the value of R



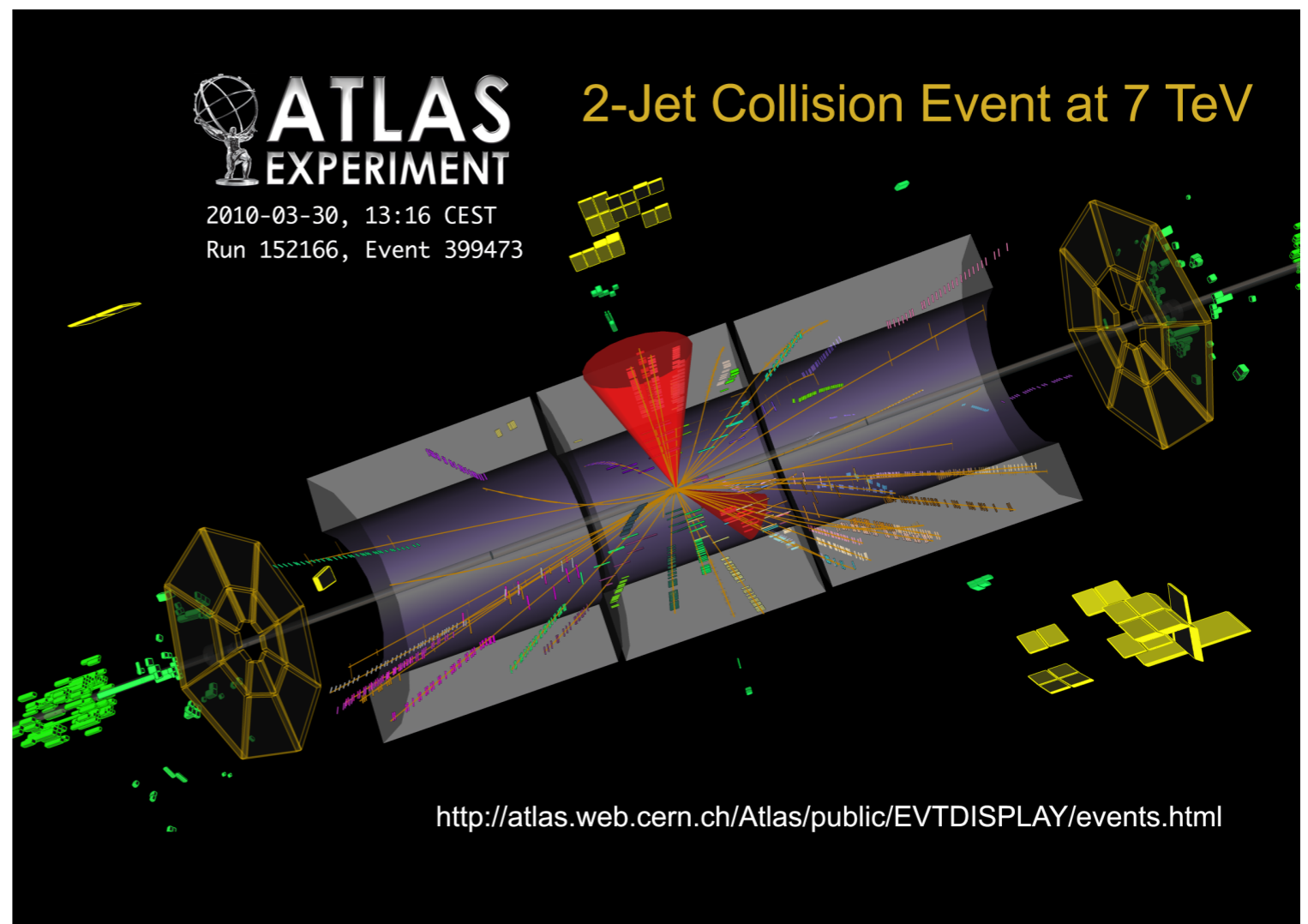
ATLAS jets

Theory setup

- NNPDF3.0_nnlo
- anti- k_T jet algorithm
- $\mu_R = \mu_F = \{p_{T1}, p_T\}$
- vary scales by factors of 2 and 1/2

Comparison to data

- ATLAS 7 TeV 4.5 fb⁻¹
- R=0.4

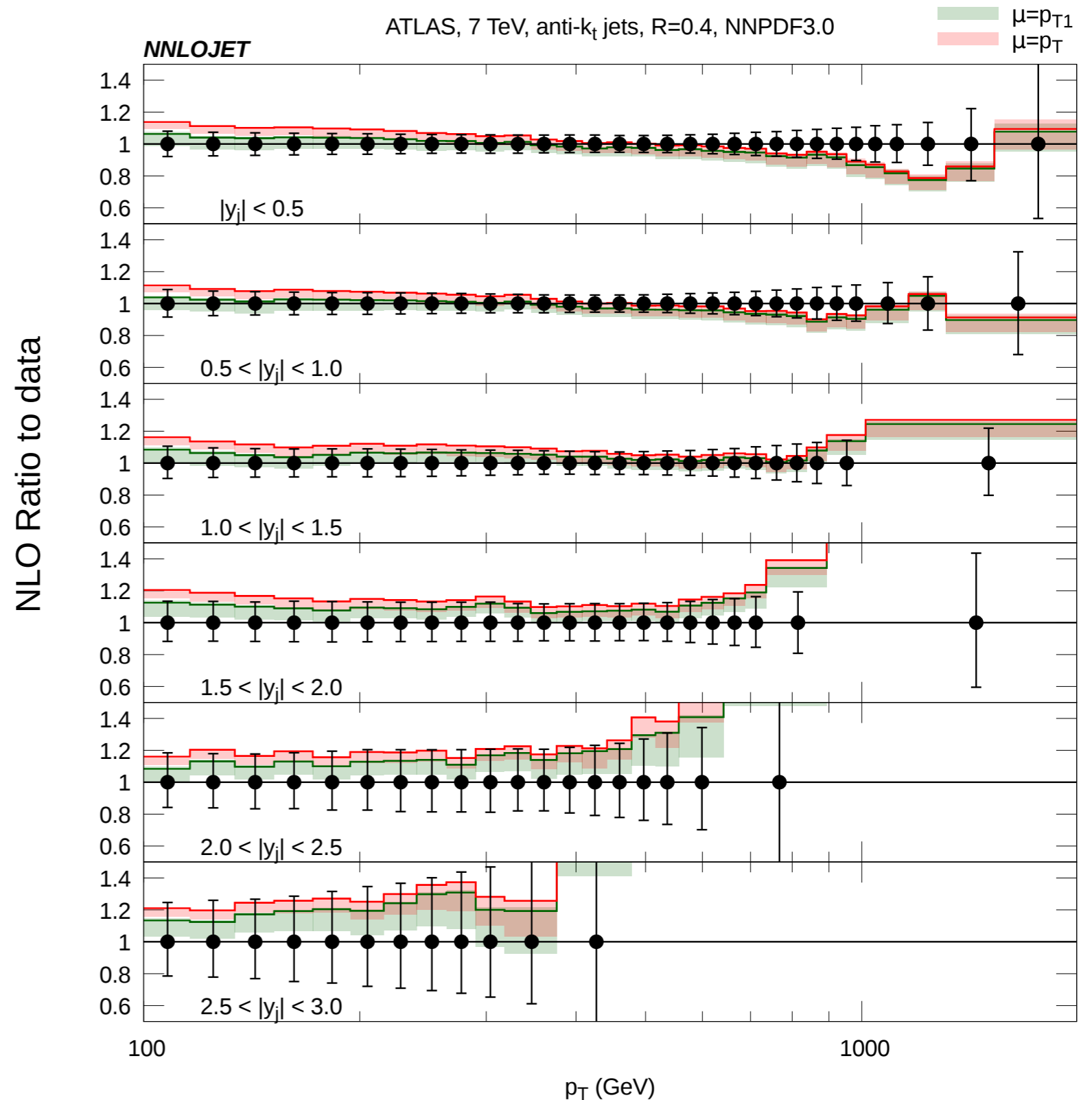


Ratio to NLO

- asymmetric scale band variation
- underestimated at small p_T due to turn over of the NLO coefficient
- 20% uncertainty for central high p_T jets rising to 40% for forward jets

Comparison to data

- non perturbative effects $< 2\%$ effect [JHEP 1509, 141 (2015)]
- data favours the p_{T1} scale choice at NLO

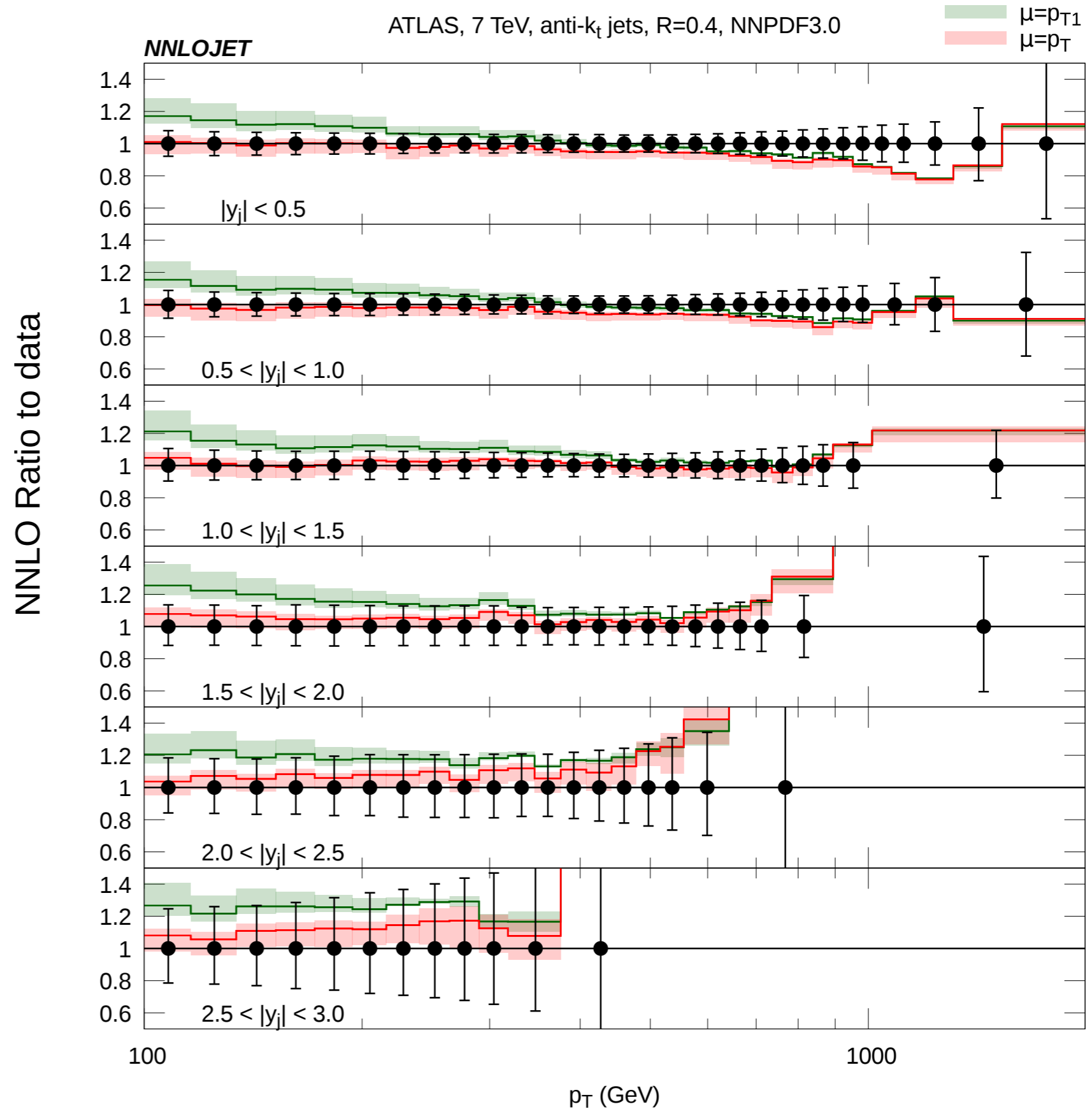


Ratio to NNLO

- symmetric scale band variation
- $p_{T1} \neq p_T$ effects enlarged at NNLO
- 10% scale uncertainty at low p_T and percent level scale uncertainty at high p_T

Comparison to data

- data favours the p_T scale choice at NNLO
- NLO EW effects around 15% for central high p_T jets



K-factor plot

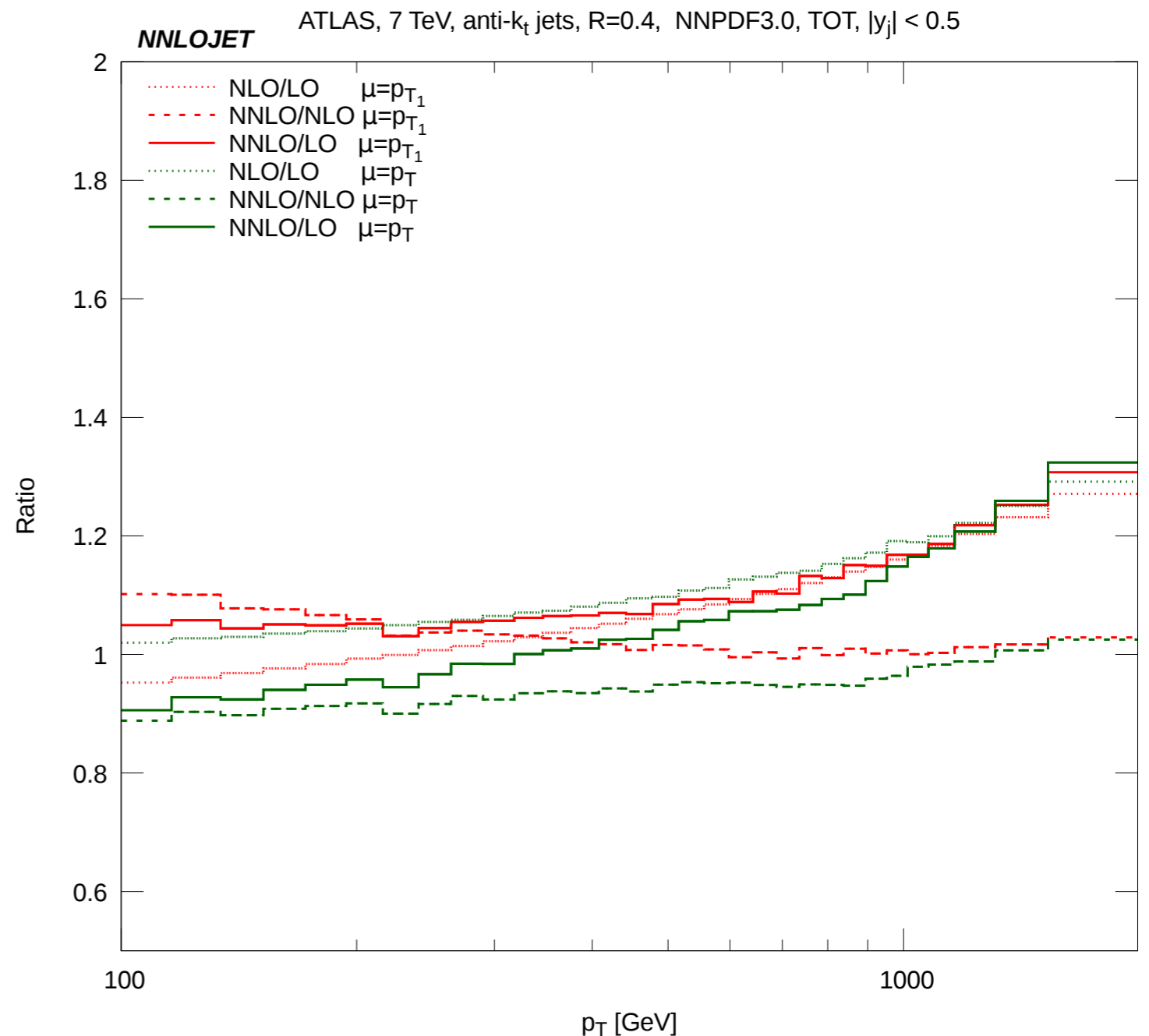
- $p_{T1} \neq p_T$ effects enlarged at NNLO at low p_T
 - decrease for larger R values

Sensible criteria for scale choice for single jet inclusive production

- perturbative stability
- data driven scale choice

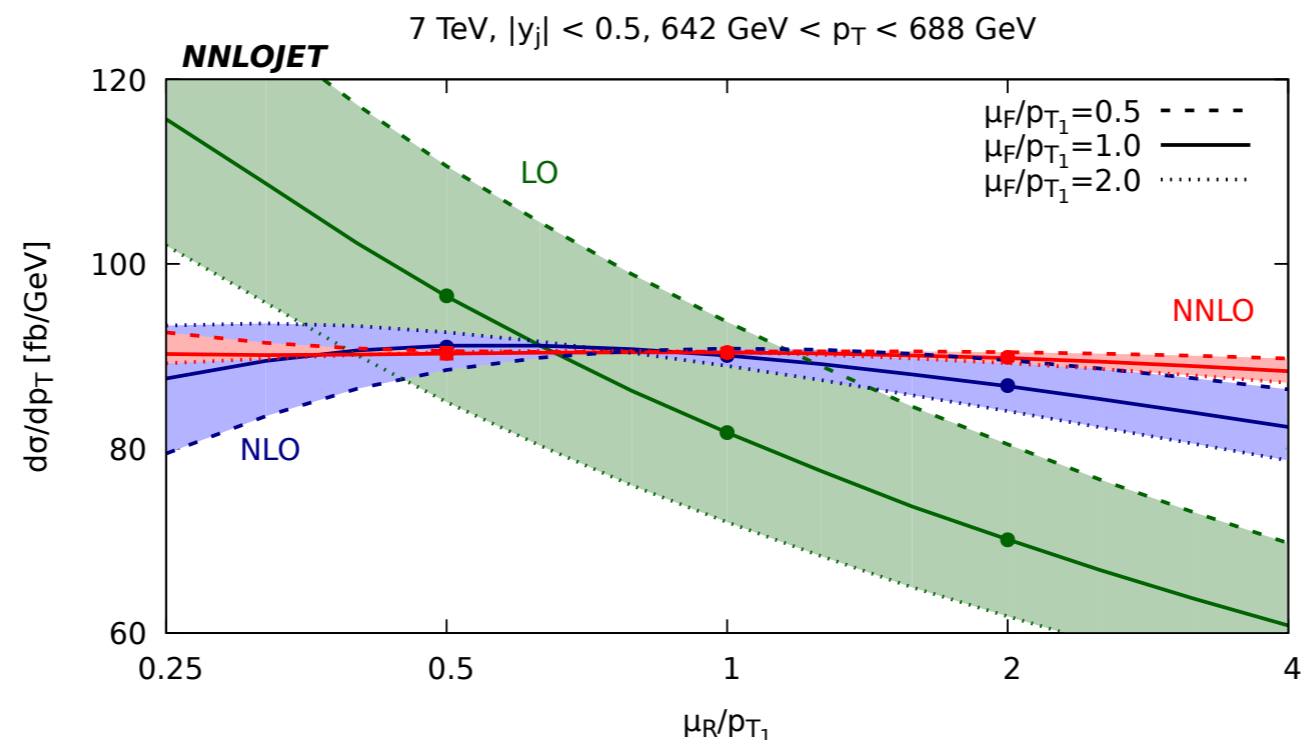
Future steps

- compare with CMS jet data; change R value ; change \sqrt{s}
- Obtain consistent description of jet data at NNLO for all jet data sets at low and high p_T in the central and forward regions for multiple R values



Summary

- Presented a new approach to look at jet data at the LHC with NNLO QCD using antenna subtraction implemented in a new parton level generator NNLOJET
- Percent level experimental accuracy in jet production is a reality at the LHC



- Percent level theory scale uncertainty for high p_T jet production at NNLO below PDF and $\alpha_s(M_Z)$ uncertainties on the cross section
- Observed significant ambiguity at small p_T due to scale choice (underestimated at NLO) which requires further phenomenological studies
- In future: 8 TeV, 13 TeV, CDF data, different PDFs etc... gateway to NNLO jet phenomenology