Top quark mass from rare two-body decays $t \to B_{(s)}^0 + jet$ in pp collisions at the LHC and FCC-hh

Morioncl-QCD 2021

(Virtual), 28th March - 3th April 2021

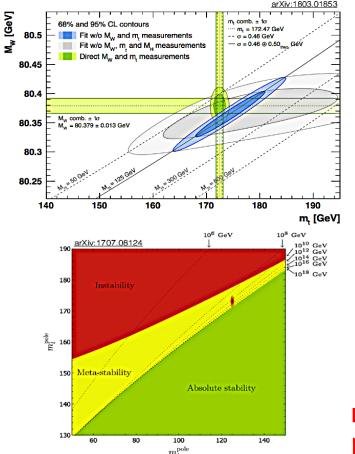
David d'Enterria (CERN) Huasheng Shao (LPTHE, Paris)

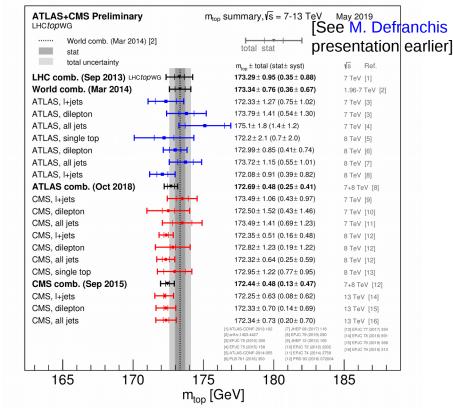
[Details in: DdE & HSS, JHEP 07 (2020) 127 2005.08102 [hep-ph]

Backgd pic: Ski down Mont Valezan (2 years ago today at Moriond-QCD'19)

Top quark mass = fundamental SM parameter

- The top quark, the heaviest SM elementary particle, plays a key role for SM & BSM physics:
 - Largest coupling to Higgs (hierarchy problem).
 - Essential for testing overall SM self-consistency (indirect BSM searches).
 - Key role on the stability of the EW vacuum in the universe.





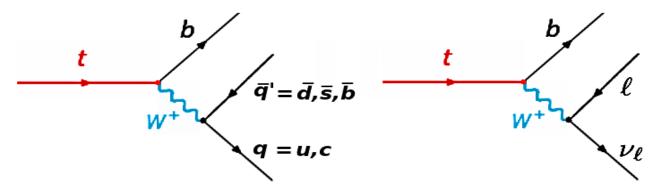
- $m_r = 172.9 \pm 0.4$ (direct), 173.1 ± 0.9 (x-sect.) GeV
- New alternative direct m_t measurements?

Moriond-QCD 2021, March 2021

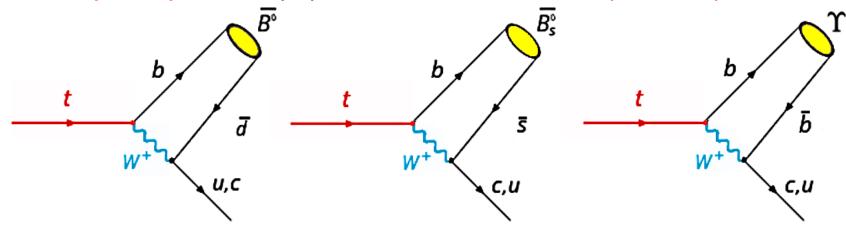
David d'Enterria (CERN)

Two-body decays of the top quark?

■ Standard 3-body top decay into a b-quark plus $W \rightarrow 2$ -quarks, 2-leptons ($V_{tb} \sim 1$):



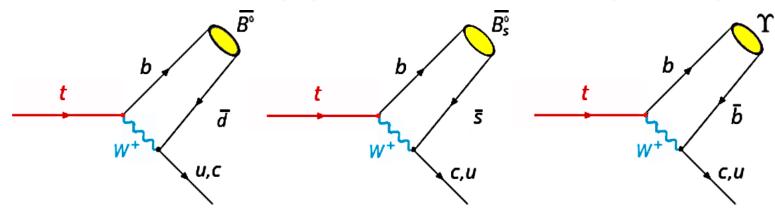
■ Two-body decay of the top quark into a bottom-meson plus u,c-quark?



- ▶ One quark from the W-decay recombines with the b-quark to form a B-meson bound state: (b, \overline{d}) , (b, \overline{s}) , or (b, \overline{b}) . Such decays have never been studied before.
- What are expected branching ratios? Visible at LHC/FCC? Top mass from them?

Two-body decays of the top quark

Two-body decay of the top quark into a bottom-meson plus u,c-quark:



- ▶ 4-momentum conservation between B, b, and q implies that the top-decay W boson is quite offshell: $m_w \approx m_t \sqrt{(m_g/m_b)} = 45$ GeV.
- ► Away-side u,c quark carries $p_T \approx m_t/2 = 85$ GeV, i.e. it is more boosted than the W decay quarks in standard $t \rightarrow bW \rightarrow bqq'$ decays.
- Recombination of b+q quark to form B-meson well computable in the Non-Relativistic QCD (NRQCD) approach [*] typically used for quarkonia. Partial width of (bq), meson determined via:

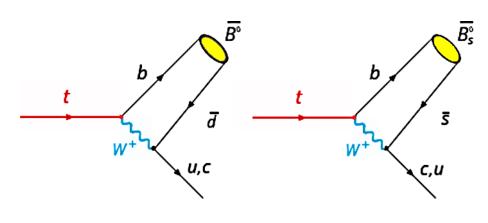
$$d\Gamma[a \to \overline{B} + X] = \sum d\hat{\Gamma}[a \to (b\overline{q})_n + X] \rho[(b\overline{q})_n \to \overline{B}]$$

Sum over (color&angular) Fock states: $n={}^{2s+1}L_J^{[c]}$ non-pQCD probability

non-pQCD probability
(NRQCD velocity power counting)

*] E. Braaten, Y. Jia, and T. Mehen, PRD66 (2002) 034003.

NRQCD calculation of $t \rightarrow B_{(s)}^0 + q$ decay



Only S-wave Fock states (higher states suppressed by $m_b/\Lambda_{OCD} \approx 20$)

$$d\Gamma[a \to \overline{B} + X] = \sum_{n} \frac{d\hat{\Gamma}[a \to (b\overline{q})_{n} + X]}{\rho[(b\overline{q})_{n} \to \overline{B}]}$$
$$n = {}^{1}S_{0}^{[1]}, S_{1}^{[1]}, S_{0}^{[8]}, S_{1}^{[8]}$$

■ HQ spin symmetry: Only 2 indep. non-pQCD probs, written as function of LDMEs:

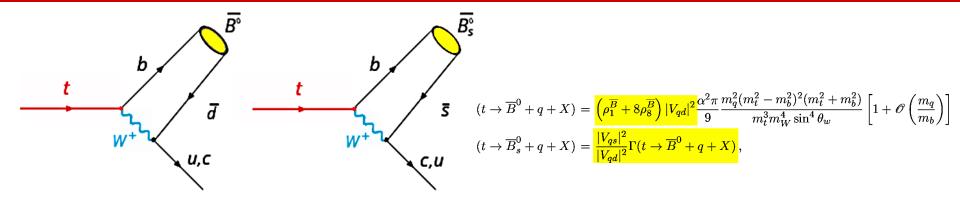
$$\begin{split} \rho_{1}^{\overline{B}} & \equiv \rho[(b\overline{q})_{1S_{0}^{[1]}} \to \overline{B}], & \langle \mathscr{O}^{\overline{B}}(^{1}S_{0}^{[1]}) \rangle = 2N_{c}\frac{4m_{b}m_{\overline{q}}^{2}}{3}\rho_{1,8}^{\overline{B}} \\ \rho_{8}^{\overline{B}} & \equiv \rho[(b\overline{q})_{1S_{0}^{[8]}} \to \overline{B}], & \langle \mathscr{O}^{\overline{B}}(^{1}S_{0}^{[1]}) \rangle = (N_{c}^{2} - 1)\frac{4m_{b}m_{\overline{q}}^{2}}{3}\rho_{1,8}^{\overline{B}} \\ 3\rho_{1,8}^{\overline{B}} & = \rho[(b\overline{q})_{3S_{1}^{[1,8]}} \to \overline{B}] & \langle \mathscr{O}^{\overline{B}}(^{3}S_{1}^{[1,8]}) \rangle = 3\langle \mathscr{O}^{\overline{B}}(^{1}S_{0}^{[1,8]}) \rangle \end{split}$$

■ Final partial widths computed by relating the $b \rightarrow B^0$ to known $c \rightarrow D^*$ -meson FFs:

$$\Gamma(t \to \overline{B}^{0} + q + X) = \frac{\left(\rho_{1}^{\overline{B}} + 8\rho_{8}^{\overline{B}}\right) |V_{qd}|^{2}}{9} \frac{\alpha^{2}\pi}{m_{t}^{3} m_{W}^{4} \sin^{4}\theta_{w}} \left[1 + \mathcal{O}\left(\frac{m_{q}}{m_{b}}\right)\right]$$

$$\Gamma(t \to \overline{B}^{0}_{s} + q + X) = \frac{|V_{qs}|^{2}}{|V_{qd}|^{2}} \Gamma(t \to \overline{B}^{0} + q + X), \quad \rho_{1,8}^{\overline{B}} = \frac{m_{c}}{m_{b}} \rho_{1,8}^{D} : \quad \rho_{1}^{\overline{B}} = 0.4 \text{ and } \rho_{8}^{\overline{B}} = 0.8$$

NRQCD calculation of $t \rightarrow B_{(s)}^0 + q$ decay

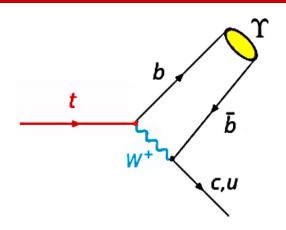


■ Final $t \rightarrow B^0(s)+q$ partial widths & branching ratios (for $\Gamma_t = 1.35$ GeV):

Decay mode	Partial decay width (GeV)	Branching fraction
$t \to \overline{B}^0 + u$	$5.4\cdot 10^{-5}$	$4.0\cdot10^{-5}$
$t o \overline{B}^0 + c$	$2.9\cdot 10^{-6}$	$2.1\cdot 10^{-6}$
$t \to \overline{B}_s^0 + c$	$5.4\cdot 10^{-5}$	$4.0\cdot 10^{-5}$
$t \to \overline{B}_s^0 + u$	$2.9\cdot 10^{-6}$	$2.1\cdot 10^{-6}$
$t \to \overline{B}_{(s)}^0 + \text{jet (total)}$	$1.14\cdot 10^{-4}$	$8.4\cdot 10^{-5}$

- ► The top-quark has a small but non-negligible (~10⁻⁴) probability to decay into B+q
- ► Particular decays depend on V_{qq} , values: $t \rightarrow B^0 + u$, $B_s^0 + c$ dominate (~50%-50%)
- [TH uncertainties large (~100%, knowledge of B-meson LDMEs) but result holds].

NRQCD calculation of $t \rightarrow Y + q$ decay



- Exemplary NRQCD use case: Veloc. scaling $v \approx \sqrt{0.1}$ valid
- S- and P-wave Fock states contribute.
- ▶ BUT, CKM-suppressed wrt. $t \rightarrow B+q$: $|V_{cb,ub}|^2/|V_{ud,cs}|^2 \approx 10^{-3}-10^{-5}$

$$d\Gamma[a \to \overline{B} + X] = \sum_{n} d\hat{\Gamma}[a \to (b\overline{q})_{n} + X] \frac{\rho[(b\overline{q})_{n} \to \overline{B}]}{\rho[(b\overline{q})_{n} \to \overline{B}]}$$

Non-pQCD probabilities as function of well-known bottomonia LDMEs:

For S-states:

$$\Gamma(t \to (H_{b\overline{b}})_{3S_{1}^{[1]}} + q) = \frac{\langle \mathscr{O}^{H_{b\overline{b}}}(^{3}S_{1}^{[1]}) \rangle}{m_{b}^{3}} \frac{\Gamma_{0}(t \to (b\overline{b})_{S} + q)}{216}$$

$$\Gamma(t \to (H_{b\overline{b}})_{3S_{1}^{[8]}} + q) = \frac{\langle \mathscr{O}^{H_{b\overline{b}}}(^{3}S_{1}^{[8]}) \rangle}{m_{b}^{3}} \frac{\Gamma_{0}(t \to (b\overline{b})_{S} + q)}{36}$$

$$\Gamma(t \to (H_{b\overline{b}})_{1S_{0}^{[8]}} + q) = \frac{\langle \mathscr{O}^{H_{b\overline{b}}}(^{1}S_{0}^{[8]}) \rangle}{m_{b}^{3}} \frac{\Gamma_{0}(t \to (b\overline{b})_{S} + q)}{12}$$

...with the prefactor in the widths:

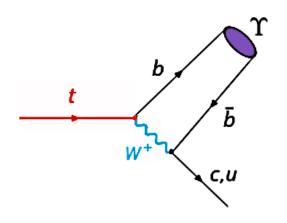
$$\Gamma_{0}(t \to (b\bar{b})_{S} + q) \equiv \frac{|V_{qb}|^{2}}{(m_{t}^{2} - 2m_{W}^{2})^{2} \sin^{4}\theta_{w}} \frac{(b\bar{b})_{S}P_{1}^{2} + D}{\Gamma(t \to (H_{b\bar{b}})_{3}P_{2}^{[1]} + q)} = \frac{(\mathcal{O}^{H_{b\bar{b}}}(^{3}P_{1}^{[2]}))}{m_{b}^{5}} \frac{\Gamma_{0}(t \to (b\bar{b})_{P} + q)}{540},$$

$$\Gamma_{0}(t \to (b\bar{b})_{P} + q) \equiv \frac{(V_{qb})^{2}}{(m_{t}^{2} - 2m_{W}^{2})^{4} \sin^{4}\theta_{w}} \frac{(b\bar{b})_{S}P_{1}^{[1]} + D}{(m_{b\bar{b}}^{2})_{S}P_{1}^{[1]} + Q} = \frac{(\mathcal{O}^{H_{b\bar{b}}}(^{3}P_{1}^{[2]}))}{m_{b}^{5}} \frac{\Gamma_{0}(t \to (b\bar{b})_{P} + Q)}{540},$$

For P-states:

$$\Gamma(t \to (H_{b\bar{b}})_{3S_{1}^{[1]}} + q) = \frac{\langle \mathscr{O}^{H_{b\bar{b}}}(3S_{1}^{[1]}) \rangle}{m_{b}^{3}} \frac{\Gamma_{0}(t \to (b\bar{b})_{S} + q)}{216} \qquad \Gamma(t \to (H_{b\bar{b}})_{3P_{0}^{[8]}} + q) = \frac{\langle \mathscr{O}^{H_{b\bar{b}}}(3P_{0}^{[8]}) \rangle}{m_{b}^{5}} \frac{\Gamma_{0}(t \to (b\bar{b})_{S} + q)}{36} \qquad \Gamma(t \to (H_{b\bar{b}})_{3P_{0}^{[8]}} + q) = \frac{\langle \mathscr{O}^{H_{b\bar{b}}}(3P_{0}^{[8]}) \rangle}{m_{b}^{5}} \frac{\Gamma_{0}(t \to (b\bar{b})_{S} + q)}{36} \qquad \Gamma(t \to (H_{b\bar{b}})_{3P_{0}^{[8]}} + q) = \frac{\langle \mathscr{O}^{H_{b\bar{b}}}(3P_{0}^{[8]}) \rangle}{m_{b}^{5}} \frac{\Gamma_{0}(t \to (b\bar{b})_{S} + q)}{36} \qquad \Gamma(t \to (H_{b\bar{b}})_{3P_{0}^{[8]}} + q) = \frac{\langle \mathscr{O}^{H_{b\bar{b}}}(3P_{0}^{[8]}) \rangle}{m_{b}^{5}} \frac{\Gamma_{0}(t \to (b\bar{b})_{S} + q)}{m_{b}^{5}} \qquad \Gamma(t \to (H_{b\bar{b}})_{3P_{0}^{[8]}} + q) = \frac{\langle \mathscr{O}^{H_{b\bar{b}}}(3P_{0}^{[8]}) \rangle}{m_{b}^{5}} \frac{\Gamma_{0}(t \to (b\bar{b})_{S} + q)}{m_{b}^{5}} \qquad \Gamma(t \to (H_{b\bar{b}})_{3P_{0}^{[8]}} + q) = \frac{\langle \mathscr{O}^{H_{b\bar{b}}}(3P_{0}^{[8]}) \rangle}{m_{b}^{5}} \frac{\Gamma_{0}(t \to (b\bar{b})_{S} + q)}{m_{b}^{5}} \qquad \Gamma(t \to (H_{b\bar{b}})_{3P_{0}^{[1]}} + q) = \frac{\langle \mathscr{O}^{H_{b\bar{b}}}(3P_{0}^{[8]}) \rangle}{m_{b}^{5}} \frac{\Gamma_{0}(t \to (b\bar{b})_{S} + q)}{m_{b}^{5}} \qquad \Gamma(t \to (H_{b\bar{b}})_{3P_{0}^{[1]}} + q) = \frac{\langle \mathscr{O}^{H_{b\bar{b}}}(3P_{0}^{[8]}) \rangle}{m_{b}^{5}} \frac{\Gamma_{0}(t \to (b\bar{b})_{S} + q)}{m_{b}^{5}} \qquad \Gamma(t \to (H_{b\bar{b}})_{3P_{0}^{[1]}} + q) = \frac{\langle \mathscr{O}^{H_{b\bar{b}}}(3P_{0}^{[8]}) \rangle}{m_{b}^{5}} \frac{\Gamma_{0}(t \to (b\bar{b})_{S} + q)}{108} \qquad \Gamma(t \to (H_{b\bar{b}})_{3P_{0}^{[1]}} + q) = \frac{\langle \mathscr{O}^{H_{b\bar{b}}}(3P_{0}^{[8]}) \rangle}{m_{b}^{5}} \frac{\Gamma_{0}(t \to (b\bar{b})_{S} + q)}{108} \qquad \Gamma(t \to (H_{b\bar{b}})_{3P_{0}^{[1]}} + q) = \frac{\langle \mathscr{O}^{H_{b\bar{b}}}(3P_{0}^{[1]}) \rangle}{m_{b}^{5}} \frac{\Gamma_{0}(t \to (b\bar{b})_{S} + q)}{108} \qquad \Gamma(t \to (H_{b\bar{b}})_{3P_{0}^{[1]}} + q) = \frac{\langle \mathscr{O}^{H_{b\bar{b}}}(3P_{0}^{[1]}) \rangle}{m_{b}^{5}} \frac{\Gamma_{0}(t \to (b\bar{b})_{S} + q)}{108} \qquad \Gamma(t \to (H_{b\bar{b}})_{3P_{0}^{[1]}} + q) = \frac{\langle \mathscr{O}^{H_{b\bar{b}}}(3P_{0}^{[1]}) \rangle}{m_{b}^{5}} \frac{\Gamma_{0}(t \to (b\bar{b})_{S} + q)}{108} \qquad \Gamma(t \to (H_{b\bar{b}})_{3P_{0}^{[1]}} + q) = \frac{\langle \mathscr{O}^{H_{b\bar{b}}}(3P_{0}^{[1]}) \rangle}{m_{b}^{5}} \frac{\Gamma_{0}(t \to (b\bar{b})_{S} + q)}{108} \qquad \Gamma(t \to (H_{b\bar{b}})_{3P_{0}^{[1]}} + q) = \frac{\langle \mathscr{O}^{H_{b\bar{b}}}(3P_{0}^{[1]}) \rangle}{m_{b}^{5}} \frac{\Gamma_{0}(t \to (h_{b\bar{b}})_{S} + q)}{108}$$

NRQCD calculation of $t \rightarrow \Upsilon + q$ decay



- **Exemplary NRQCD use case:** Veloc. scaling $v \approx \sqrt{0.1}$ valid
- S- and P-wave Fock states contribute.
- ▶ BUT, CKM-suppressed wrt. $t \rightarrow B+q$: $|V_{cb,ub}|^2/|V_{ud,cs}|^2 \approx 10^{-3}-10^{-5}$

$$d\Gamma[a \to \overline{B} + X] = \sum_{n} \underline{d\hat{\Gamma}[a \to (b\overline{q})_{n} + X]} \rho[(b\overline{q})_{n} \to \overline{B}]$$

t → Y+q partial widths for all bottomonia states & four sets of LDMEs:

top decay (bottomonium state)	Partia Set I	-	width (10 Set III	o ⁻⁹ GeV) Set IV
$t \to \Upsilon(1S) + q$	1.62	2.01	1.61	1.36
$t \to \Upsilon(2S) + q$	0.71	0.23	0.67	0.57
$t \to \Upsilon(3S) + q$	0.51	0.47	0.42	0.37
$t \to \Upsilon(nS) + q$	2.84	2.71	2.70	2.30

■ Final $t \rightarrow Y+q$ partial widths & branching ratios (for $\Gamma_t = 1.35$ GeV):

Decay mode	Partial decay width (GeV)	Branching fraction
$t \to \Upsilon(nS) + c$	$2.5\cdot 10^{-9}$	$1.9 \cdot 10^{-9}$
$t \to \Upsilon(nS) + u$	$2.5\cdot 10^{-11}$	$1.9\cdot 10^{-11}$

- ▶ The top-quark has a very small (~10-9) probability to decay into Υ +c.
- [Small TH uncertainties ~10%, thanks to good knowledge of b-bbar LDMEs].

Experimental observation of $t \rightarrow B_{(s)}^0 + q$?

■ How many rare $\mathcal{O}(10^{-4})$ 2-body decays at HL-LHC (3 ab⁻¹), FCC-hh (20 ab⁻¹)?

\sqrt{s}	$\sigma_{ m NNLO+NNLL}^{tar{t}, m incl.}$ (top++)
14 TeV	$980 \pm 17 (PDF) ^{+24}_{-35} (scale) pb$
$100~{\rm TeV}$	$34.80 \pm 1.20 (PDF) ^{+1.00}_{-1.65} (scale) \mathrm{nb}$

We expect 6·10⁹, 1.5·10¹² top-quark pair events at the HL-LHC, FCC

(\times 250 more at FCC thanks to \times 35 σ and \times 7 \mathscr{L})

- Number of $t \rightarrow B$ +jet decays in triggered top-pair events at HL-LHC, FCC-hh:
 - Fully hadronic: $t\bar{t} \to W^+ b \ W^- \bar{b} \to q\bar{q}' b \ q''\bar{q}'''\bar{b} \ (\mathscr{B} = 45.7\%)$
 - Lepton+jets: $t\bar{t} \to W^+ b \ W^- \bar{b} \to q\bar{q}' b \ \ell^- \nu_\ell \bar{b} \ (\mathscr{B} = 43.8\%)$
 - Dileptons: $t\bar{t} \to W^+ b \ W^- \bar{b} \to \ell^+ \nu_\ell b \ \ell^- \nu_\ell \bar{b} \ (\mathscr{B} = 10.5\%)$

Run MadGraph@NLO (NNPDF3.0) for pp → ttbar → 6j, ℓ+j at LHC/FCC

- ► Typical acceptance×efficiency for ttbar multi-jets & lepton+jets final states: ~20%
 - At least 4 jets (reconstructed with the anti- $k_{\rm T}$ algorithm R = 0.5) with $p_{\rm T} > 25$ GeV and $|\eta| < 3.0$ (5.0 at the FCC);
 - A total scalar sum of transverse energy in the reconstructed jets above roughly twice the top mass, $H_T > 350 \text{ GeV}$;
 - At least 1 b jet tagged with a typical 75% efficiency;

- One isolated (cone radius $R_{\rm isol} = 0.3$) charged lepton ℓ with $p_{\rm T} > 30$ GeV and $|\eta| < 2.5$ ($|\eta| < 5.0$ at the FCC);
- At least 2 jets (reconstructed with the anti- $k_{\rm T}$ algorithm R=0.5) with $p_{\rm T}\geq 25$ GeV and $|\eta|<3.0$ (5.0 at the FCC), and separated from the lepton by $\Delta R(\ell,j)>0.4$;
- 1 b-jet tagged with a typical 75% efficiency;

<i>[</i> 2	top-quark pair <mark>%× accep</mark>	$otance \times efficiency losses$
\sqrt{s}	fully hadronic	ℓ + jets
14 TeV	$0.457 \times 0.57 \times 0.75 = 0.20$	$0.35 \times 0.64 \times 0.75 = 0.17$
$100 \mathrm{TeV}$	$0.457 \times 0.76 \times 0.75 = 0.26$	$0.35 \times 0.71 \times 0.75 = 0.19$

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■ How many rare $\mathcal{O}(10^{-4})$ 2-body decays at HL-LHC (3 ab⁻¹), FCC-hh (20 ab⁻¹)?

$$\sqrt{s}$$
 $\sigma_{\text{NNLO+NNLL}}^{t\bar{t},\text{incl.}}$ (top++)

14 TeV $980 \pm 17 \, (\text{PDF})_{-35}^{+24} (\text{scale}) \, \text{pb}$

100 TeV $34.80 \pm 1.20 \, (\text{PDF})_{-1.65}^{+1.00} (\text{scale}) \, \text{nb}$

We expect 6·10⁹, 1.5·10¹² top-quark pair events at the HL-LHC, FCC

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 - A total scalar sum of transverse energy in the reconstructed jets above roughly twice the top mass, $H_T > 350$ GeV;
 - At least 1 b jet tagged with a typical 75% efficiency;

LHC:
$$N_{\mathrm{t} \to \overline{\mathrm{B}}_{(\mathrm{s})}^{0} + \mathrm{jet}}^{t\overline{t}} \approx 1.5 \cdot 10^{5}$$

FCC: $N_{\mathrm{t} \to \overline{\mathrm{B}}_{(\mathrm{s})}^{0} + \mathrm{jet}}^{t\overline{\mathrm{f}} \, \mathrm{full \ had.}} pprox 4 \cdot 10^{7}$

- One isolated (cone radius $R_{\rm isol}=0.3$) charged lepton ℓ with $p_{\scriptscriptstyle \rm T}>30$ GeV and $|\eta|<2.5$ ($|\eta|<5.0$ at the FCC);
- At least 2 jets (reconstructed with the anti- $k_{\rm T}$ algorithm R=0.5) with $p_{\rm T}\geq 25$ GeV and $|\eta|<3.0$ (5.0 at the FCC), and separated from the lepton by $\Delta R(\ell,j)>0.4$;
- 1 b-jet tagged with a typical 75% efficiency;

LHC:
$$N_{\rm t \to \overline{B}_{(s)}^0 + \rm jet}^{tt \, \ell + \rm jet} \approx 62\,000$$

FCC:
$$N_{t \to \overline{\mathrm{B}}_{(\mathrm{s})}^{0} + \mathrm{jet}}^{t\bar{t}} \approx 1.6 \cdot 10^{7}$$

Search for $t \rightarrow B_{(s)}^0$ +jet via $B_{(s)}^0 \rightarrow J/\phi$ +hh final states

■ 6 exclusive $B_{(s)}^0 \rightarrow J/\phi$ +hh decay channels with total branching ratios 10⁻⁵–10⁻⁴:

$\overline{B}_{(s)}^0$ meson	Total branching fraction		$_{ m tance}$
decay mode	(product of individual \mathcal{B} 's)		FCC
$\overline{B}^0 \to J/\psi(\mu^+\mu^-) K^{*0}(K^{\pm}\pi^{\mp})$	$1.27 \cdot 10^{-3} \times 0.0596 \times 1 \approx 7.6 \cdot 10^{-5}$	50%	54%
$\overline B^0 o J/\psi(\mu^+\mu^-) K^+\pi^-$	$1.15 \cdot 10^{-3} \times 0.0596 \approx 6.9 \cdot 10^{-5}$	62%	67%
$\overline{B}^0 \to J/\psi(\mu^+\mu^-) K_s^0(\pi^+\pi^-)$	$0.87 \cdot 10^{-3} \times 0.0596 \times 0.69 \approx 3.6 \cdot 10^{-5}$	52%	58%
$\overline{B}_s^0 \to J/\psi(\mu^+\mu^+)\phi(K^+K^-)$	$1.08 \cdot 10^{-3} \times 0.0596 \times 0.49 \approx 3.2 \cdot 10^{-5}$	61%	64%
$\overline B^0_s o J/\psi(\mu^+\mu^+)K^+K^-$	$0.79 \cdot 10^{-3} \times 0.0596 \approx 4.7 \cdot 10^{-5}$	63%	67%
$\overline{B}_s^0 \to J/\psi(\mu^+\mu^+) \pi^+\pi^-$	$0.21 \cdot 10^{-3} \times 0.0596 \approx 1.25 \cdot 10^{-5}$	66%	68%

Acceptances determined via PYTHIA8 simulations for all 6 final states in $tt \rightarrow B^0_{(s)} + X \rightarrow J/\phi + hh + X$ events in pp at $\sqrt{s} = 14$, 100 TeV. Reco effic. assumed at ~0.5.

Top pair	$\overline{B}_{(s)}^0$ meson	Events after all selection cuts		
final state	decay mode	LHC (3 ab^{-1})	$FCC (20 \text{ ab}^{-1})$	
$t\overline{t} o \overline{B}^0 + u, c$				
o fully hadronic	$\overline B^0 o J/\psi(\mu^+\mu^-)K^+\pi^-$	3.1	1050	
∘ ℓ+ jets	$B \rightarrow J/\psi(\mu \cdot \mu^-) K \cdot \pi$	1.4	380	
o fully hadronic	$\overline{B}^0 \to J/\psi(\mu^+\mu^-) K^{*0}(K^{\pm}\pi^{\mp})$	2.7	940	
∘ ℓ+ jets	$B \rightarrow J/\psi(\mu^+\mu^-)K^-(K^-\pi^+)$	1.2	340	
o fully hadronic	$\overline B^0 o J/\psi(\mu^+\mu^-)K^0_s(\pi^+\pi^-)$	1.3	470	
∘ ℓ+ jets	$B \rightarrow J/\psi(\mu^+\mu^-)K_s(\pi^+\pi^-)$	0.6	170	
Sum of all channels		10.2	3 300	
$t\overline{t} o \overline{B}_s^0 + c, u$				
o fully hadronic	$\overline{B}_s^0 \to J/\psi(\mu^+\mu^+) K^+K^-$	2.1	710	
∘ ℓ+ jets	$B_s \to J/\psi(\mu \cdot \mu \cdot) K \cdot K$	0.9	250	
o fully hadronic	$\overline{B}^0_s \to J/\psi(\mu^+\mu^+)\phi(K^+K^-)$	1.4	460	
∘ ℓ+ jets	$B_s \to J/\psi(\mu \cdot \mu \cdot) \psi(\mathbf{K} \cdot \mathbf{K})$	0.6	170	
o fully hadronic	$\overline B{}^0_c o J/\psi(\mu^+\mu^+)\pi^+\pi^-$	0.6	200	
∘ ℓ+ jets	$B_s \to J/\psi(\mu^+\mu^+)\pi^+\pi^-$	0.3	70	
Sum of all channels		5.7	1 900	
TOTAL		16	5 200	

5-thousand (~15) 2-body top-quark decays reconstructed at FCC-hh (LHC).

Search for $t \rightarrow B_{(s)}^0$ +jet via $B_{(s)}^0 \rightarrow D$ +h(h) final states

■ 6 exclusive $B_{(s)}^0 \to D+h(h)$ decay channels with total branching ratios ~10⁻⁴:

$\overline{B}_{(s)}^0$ meson	Total branching fraction	Acceptance	
decay mode	(product of individual \mathcal{B} 's)	LHC	FCC
$\overline{B}^0 \to D^+(K^-\pi^+\pi^+) \pi^-$	$2.52 \cdot 10^{-3} \times 0.094 \approx 2.37 \cdot 10^{-4}$	55%	65%
$\overline{B}^0 \to \overline{D}^{*,+}(D^0(K^-\pi^+)\pi^+) \pi^-$	$2.74 \cdot 10^{-3} \times 0.677 \times 0.039 \approx \frac{7.2 \cdot 10^{-5}}{10^{-5}}$	55%	65%
$\overline{B}^0 \to \overline{D}^0(K^-\pi^+) \pi^+\pi^-$	$0.88 \cdot 10^{-3} \times 0.039 \approx 3.43 \cdot 10^{-5}$	75%	80%
$\overline{B}^0 \to D^+(K^-\pi^+\pi^+) K^-$	$0.186 \cdot 10^{-3} \times 0.094 \approx 1.75 \cdot 10^{-5}$	55%	65%
$\overline{B}_{s}^{0} \to D_{s}^{+}(K^{-}K^{+}\pi^{+})\pi^{-}$	$3.0 \cdot 10^{-3} \times 0.055 \approx 1.6 \cdot 10^{-4}$	60%	60%
$\overline{B}_s^0 \to \overline{D}^0(K^-\pi^+) K^-\pi^+$	$1.04 \cdot 10^{-3} \times 0.039 \approx 4 \cdot 10^{-5}$	60%	65%

Acceptances determined via PYTHIA8 simulations (ongoing) for all 6 final states in $tt \rightarrow B^0_{(s)} + X \rightarrow D + h(h) + X$ events in pp at $\sqrt{s} = 14$, 100 TeV. Reco effic. assumed at ~0.5.

ion cuts (20 ab ⁻¹)
<u> </u>
500
500
200
100
380
630
220
250
90
300
200
800
570
200
800
1 000

▶ 10-thousand (~30) 2-body top-quark decays reconstructed at FCC-hh (LHC).

Search for 2-body top decay via t → j_{bottom}+j_{charm}

- 12 exclusive B⁰_(s) → J/φ+hh,D+h(h) decays are very clean, but have small BRs. Try instead to reconstruct the B-meson as a b-jet & search for the 2-body decay via t → j(bottom)+j(charm)? [t → j(bottom)+j(u-quark) swamped by backgd.]
- Analysis of ttbar ℓ +jets triggered events requiring N=3 jets w/ exactly 2 b-tags + 1 c-tag (p_T>80 GeV) within acceptance with (mis)tagging performances:
 - Bottom jets: b-jet tagging efficiency: 75%; b-jet mistagging probability for a c-quark: 5%, b-jet mistagging probability for udsg (light quarks or gluon): 0.5%.
 - Charm jets: Two different c-jet working points considered: (i) c-jet tagging efficiency: 65%, c-jet mistagging probability for a b quark: 10%, c-jet mistagging probability for udsg: 10% [HcT:]; and (ii) c-jet tagging efficiency: 35%, c-jet mistagging probability for a b-quark: 5%, and c-jet mistagging probability for udsg: 1% [LcT].

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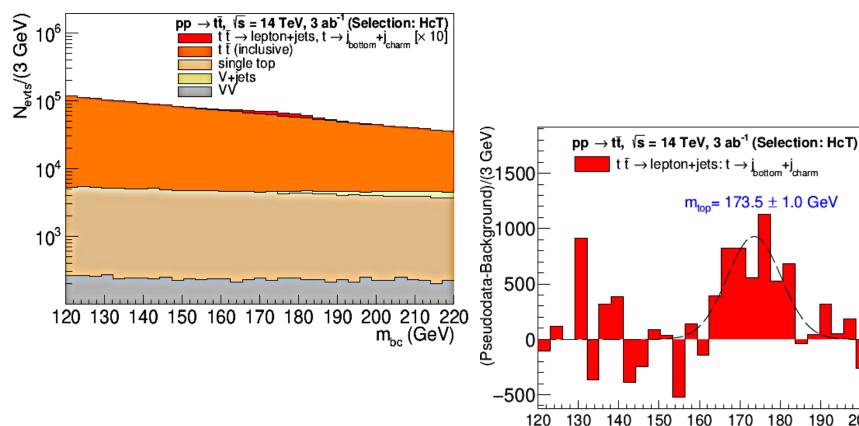
- MadGraph@NLO (NNPDF3.0)
 - + PYTHIA8 PS simulation for all backgrounds:
 - pp \rightarrow ttbar \rightarrow 6j, ℓ +j
 - single-t
 - W+jets, Z+jets
 - **■** VV (V=W,Z)
 - QCD multi-jets

nrogogg	pp at $\sqrt{s} = 14$	$4 {\rm TeV} (3 {\rm ab}^{-1})$	pp at $\sqrt{s} = 10$	$0 { m TeV} (20 { m ab}^{-1})$
process	'HcT'	$^{\prime}\mathrm{LcT}^{\prime}$	'HcT'	$^{\prime}\mathrm{LcT}^{\prime}$
$t\bar{t}$ (signal)	5 300	2500	$1.4 \cdot 10^6$	$7.5 \cdot 10^5$
$t\bar{t}$ (backgd.)	$2.1(1.7) \cdot 10^6$	$9.3(7.4) \cdot 10^5$	$3.0(2.4)\cdot 10^8$	$1.1(0.9) \cdot 10^8$
tW	$7.0(5.8) \cdot 10^4$	$3.1(2.5) \cdot 10^4$	$1.2(1.0) \cdot 10^7$	$4.7(4.4) \cdot 10^6$
$t+\mathrm{jet}$	$7.8(6.6) \cdot 10^4$	$2.7(2.4) \cdot 10^4$	$1.1(1.0) \cdot 10^7$	$4.0(3.7) \cdot 10^6$
$V+\mathrm{jet}$	$1.5(1.3) \cdot 10^5$	$7.6(7.6) \cdot 10^4$	$1.1(0.8) \cdot 10^7$	$1.1(0.8) \cdot 10^7$
VV	$7.7(6.3) \cdot 10^3$	$2.8(2.3) \cdot 10^3$	$9.1(7.3) \cdot 10^5$	$3.4(3.1) \cdot 10^5$
Multijet	$< 4.8 \cdot 10^4$	$<4.8\cdot 10^4$	$< 1.2 \cdot 10^7$	$<1.2\cdot 10^7$
\mathcal{S} (std. dev.)	6.1	4.5	130.	110.

► Few million (thousands) $t \rightarrow b+c$ expected at FCC-hh (LHC).

Search for 2-body top decay via t → j_{bottom}+j_{charm} (LHC)

- Expected dijet signal $t \rightarrow b+c$ ($m_{jj}\sim 10$ GeV resolution) after cuts injected on top of MadGraph@NLO+ PYTHIA8 simulation for dominant pp \rightarrow ttbar \rightarrow 6j, ℓ +j backgds
- Invariant mass dijet distribution for best (mis)tagging selection criterion:

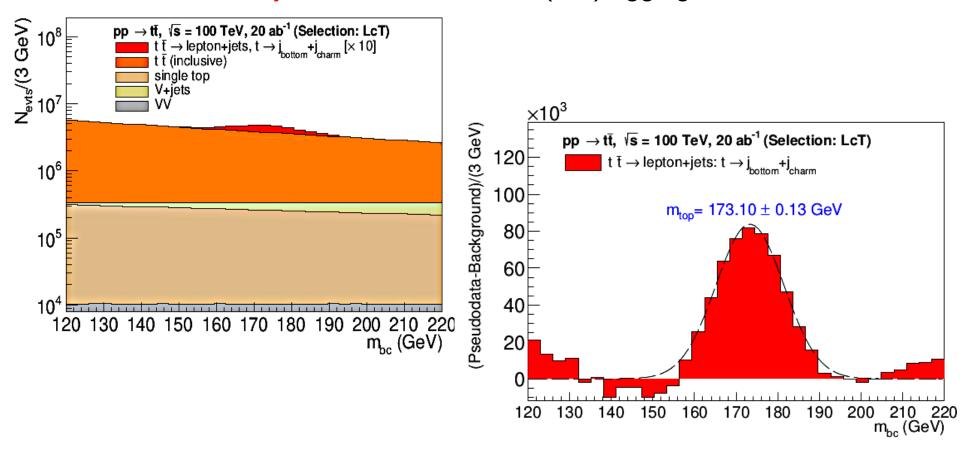


▶ Observation warranted at HL-LHC (significance >6 σ) for t \rightarrow b+c signal.

200 210 220 m_{bc} (GeV)

Search for 2-body top decay via t → j_{bottom}+j_{charm} (FCC)

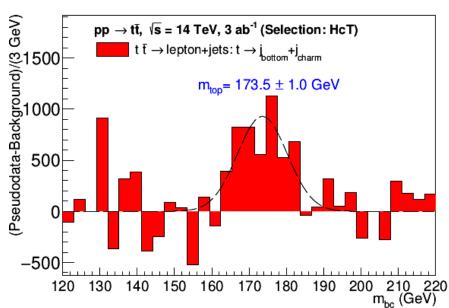
- Expected dijet signal $t \rightarrow b+c$ ($m_{jj}\sim 10$ GeV resolution) after cuts injected on top of MadGraph@NLO+ PYTHIA8 simulation for dominant pp \rightarrow ttbar \rightarrow 6j, ℓ +j backgds
- Invariant mass dijet distribution for best (mis)tagging selection criterion:



► Huge $t \rightarrow b+c$ signal (significance >100 σ) over background at the FCC.

Top mass from $t \rightarrow j_{bottom} + j_{charm}$ invariant mass

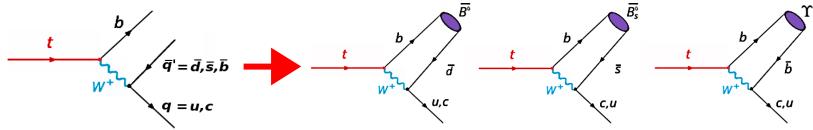
- Two-body top-quark decay provides novel interesting possibilities for simple top mass extractions from 2-body invariant mass analysis of:
- 1) $m_{inv}(B^0+jet)$ from 12 exclusive $B^0_{(s)} \rightarrow J/\phi + hh, D+h(h)$ final states plus jet:
 - Small systematical uncertainties dominated by single-jet energy scale.
 - Example Stat.uncertainties dominated by small data samples (~50, 15000 evts at LHC,FCC)
- 2) m_{inv} (b-jet + c-jet):
 - Small statistical uncertainties (~million evts at FCC).
 - (2) Moderate systematical uncertainties dominated by double-jet energy scale.
- Simple Gaussian fit of the backgd-subtracted m_{bc} distribution at LHC:



- Top mass extraction with ~1 GeV fit uncertainty.
- Dijet energy scale should be fitted concurrently & assoc. systematics propagated.
- Promising novel extraction of m_{top} with different systematics to other methods

Summary

First-ever study of 2-body decays of a top-quark into B-meson + u,c jet:



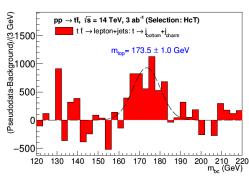
Partial widths & branching ratios computed in NRQCD recombination:

Decay mode	Partial decay width (GeV)	Branching fraction
$t \to \overline{B}^0 + u$	$5.4\cdot10^{-5}$	$4.0\cdot 10^{-5}$
$t \to \overline{B}^0 + c$	$2.9 \cdot 10^{-6}$	$2.1\cdot 10^{-6}$
$t \to \overline{B}_s^0 + c$	$5.4\cdot10^{-5}$	$4.0\cdot 10^{-5}$
$t \to \overline{B}_s^0 + u$	$2.9 \cdot 10^{-6}$	$2.1\cdot 10^{-6}$
$t \to \overline{B}_{(s)}^0 + \text{jet (total)}$	$1.14 \cdot 10^{-4}$	$8.4 \cdot 10^{-5}$

- ► Top-quark has small but visible (~10-4) probability to decay into $t \rightarrow B^0+u$, B^0_s+c .
- ► Top-quark has tiny (~10-9) probability to decay into Y+c.
- Feasibility study in pp collisions at HL-LHC & FCC-hh with NNLO rates & MadGraph@NLO+PY8 simulations: exclusive channels & b+c dijet final state

pp system	Number of produced		Number of 2-body top-quark events		
$\sqrt{s},~\mathcal{L}_{\mathrm{int}}$	events expected		after $\mathcal{B} \times$ acceptance \times efficiency cuts		
	$t\bar{t}$ $t \to \overline{B}_{(s)}^0 + \mathrm{jet}$		$\overline{B}_{(s)}^0 \to J/\psi hh$	$\overline{B}_{(s)}^0 \to D_{(s)}^{0,\pm} h(h)$	$t \to b + c$
$14 \text{TeV}, 3 \text{ab}^{-1}$	$3 \cdot 10^9$	$2.5\cdot 10^5$	16	33	5300
$100 \text{TeV}, 20 \text{ab}^{-1}$	$7\cdot 10^{11}$	$6\cdot 10^7$	5 200	11000	$1.4\cdot 10^6$

Novel precise method to measure top mass via simple invariant-mass analysis of 2-body final states:



Backup slides

t → j_{bottom} + j_{charm} background

Cut-flow numbers for accumulated selection criteria for the dominant ttbar background:

Selection criteria	pp at $\sqrt{s} = 14 \text{TeV} (3 \text{ab}^{-1})$		pp at $\sqrt{s} = 100 \text{TeV} (20 \text{ab}^{-1})$	
	'HcT'	$^{\prime}\mathrm{LcT}^{\prime}$	'HcT'	$^{\prime}\mathrm{LcT}^{\prime}$
(0): none	$1.0 \cdot 10^{9}$	$1.0 \cdot 10^{9}$	$2.4 \cdot 10^{11}$	$2.4 \cdot 10^{11}$
(1): trigger $(\ell+jets)$	$4.8 \cdot 10^{8}$	$4.8\cdot 10^8$	$1.3 \cdot 10^{11}$	$1.3\cdot10^{11}$
(2): (1) + $(N_{j_b} = 2 \&\& N_{j_c} = 1)$	$8.2 \cdot 10^7$	$3.8\cdot 10^7$	$2.3\cdot 10^{10}$	$1.2\cdot10^{10}$
(3): $(1) + (2) + 120 \le m_{b,c} \le 220 \text{GeV}$	$5.3 \cdot 10^7$	$2.6\cdot 10^7$	$1.3 \cdot 10^{10}$	$7.3 \cdot 10^9$

Table 10. Total number of events expected for $t\bar{t} \to N_j + \ell^{\pm} + X$ with $\ell^{\pm} = \mu^{\pm}, e^{\pm}$, in pp collisions at $\sqrt{s} = 14$ and 100 TeV, passing the analysis criteria for the 'HcT' and 'LcT' charm jet reconstruction performances discussed in the text. The "trigger" row lists the number of events that pass the trigger cuts presented in subsection 3.1.2, and $N_{j_b} = 2 \&\& N_{j_c} = 1$ indicates the exact requirement of two b- and one c-quark jets reconstructed in the event with the kinematic criteria discussed in the text.